

Energy Loss in the Hot QCD Brick I

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With many thanks to Brian Cole, Miklos Gyulassy, Yuri Kovchegov, and TECHQM Collab.

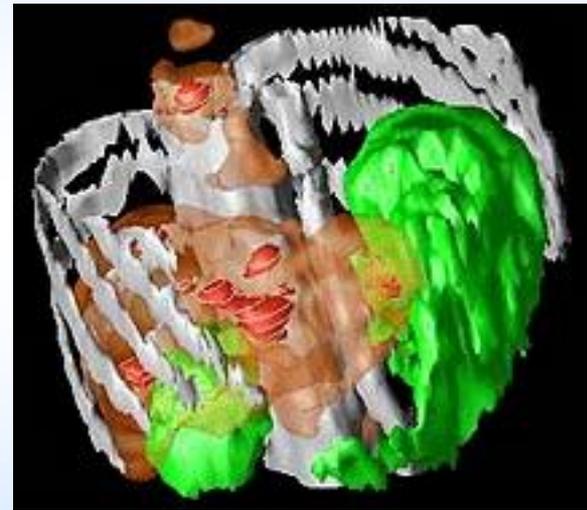
Outline

- High- p_T motivation
- Differences seen in parameters OK
- Collinear approximation leads to large systematic theoretical uncertainty
- Thermal gluon mass effects (surprisingly!) small
- Conclusions

Why Study High- p_T Particles?

- Learn about QGP medium
 - Measure suppression, use theory to invert
 - density profile and evolution
 - transport coefficients

- Tomography in Medicine

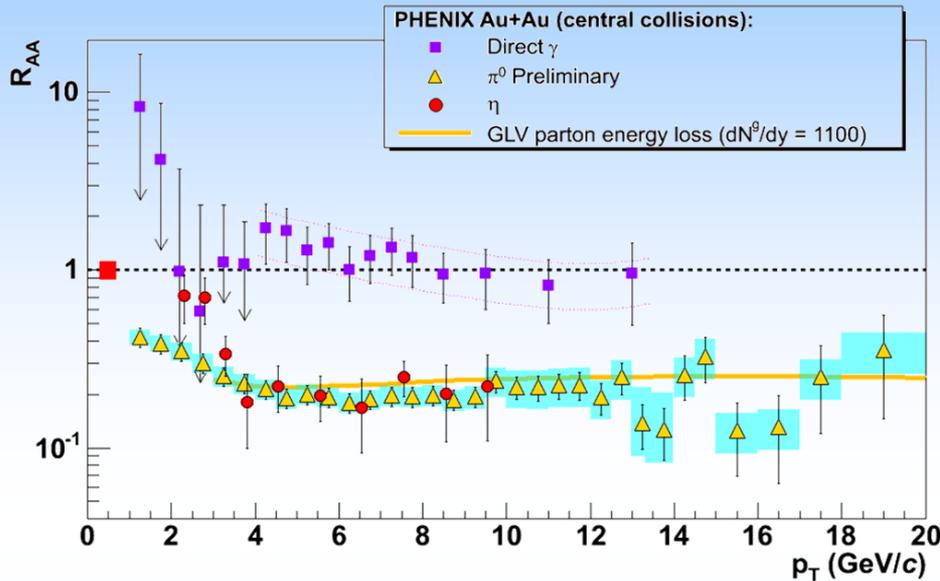


SPECT-CT Scan uses internal γ photons and external X-rays

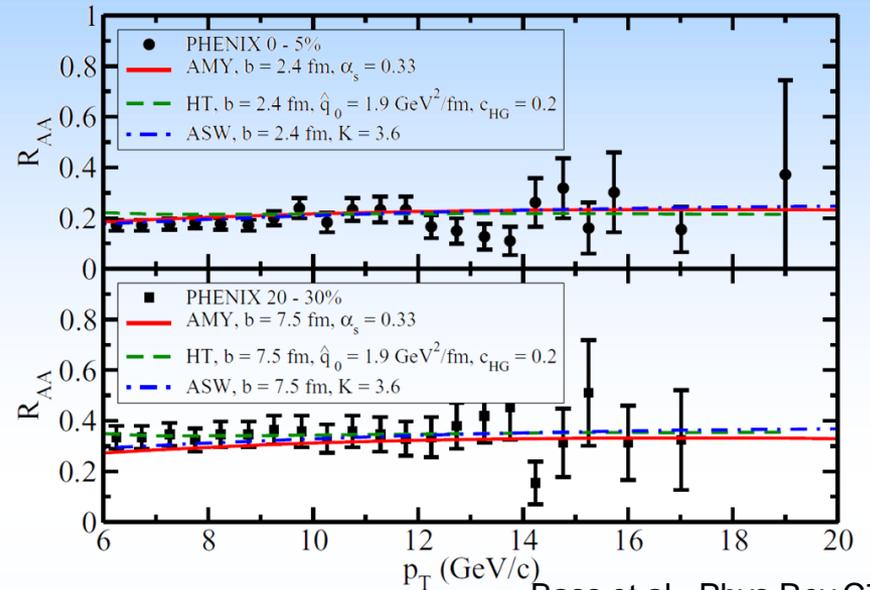
Acronyms, Acronyms

- Four major pQCD formalisms for Rad E-loss
 - Opacity expansion: GLV (DGLV), ASW-SH
 - Multiple soft scattering: BDMPS (ASW-MS)
 - Higher Twist: HT
 - Thermal field theory: AMY

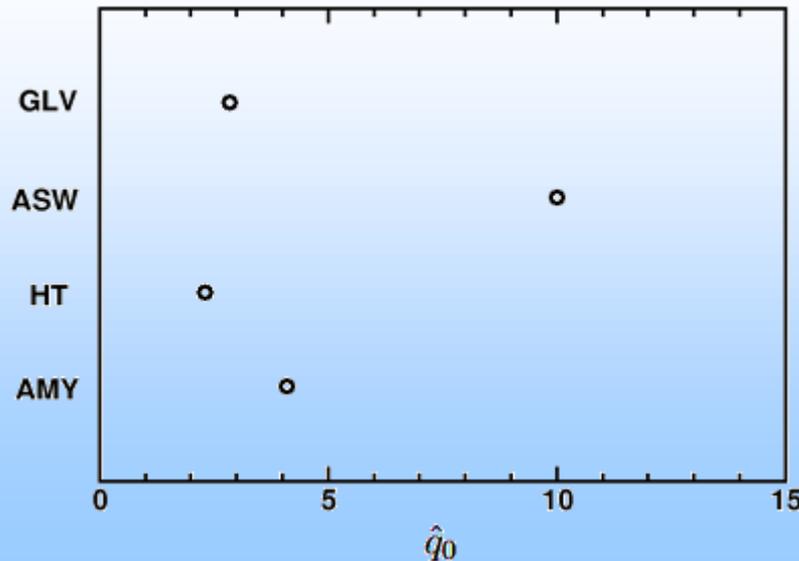
Different Formalisms, Different Results



Y. Akiba for the PHENIX collaboration,
hep-ex/0510008



Bass et al., Phys.Rev.C79:
024901,2009



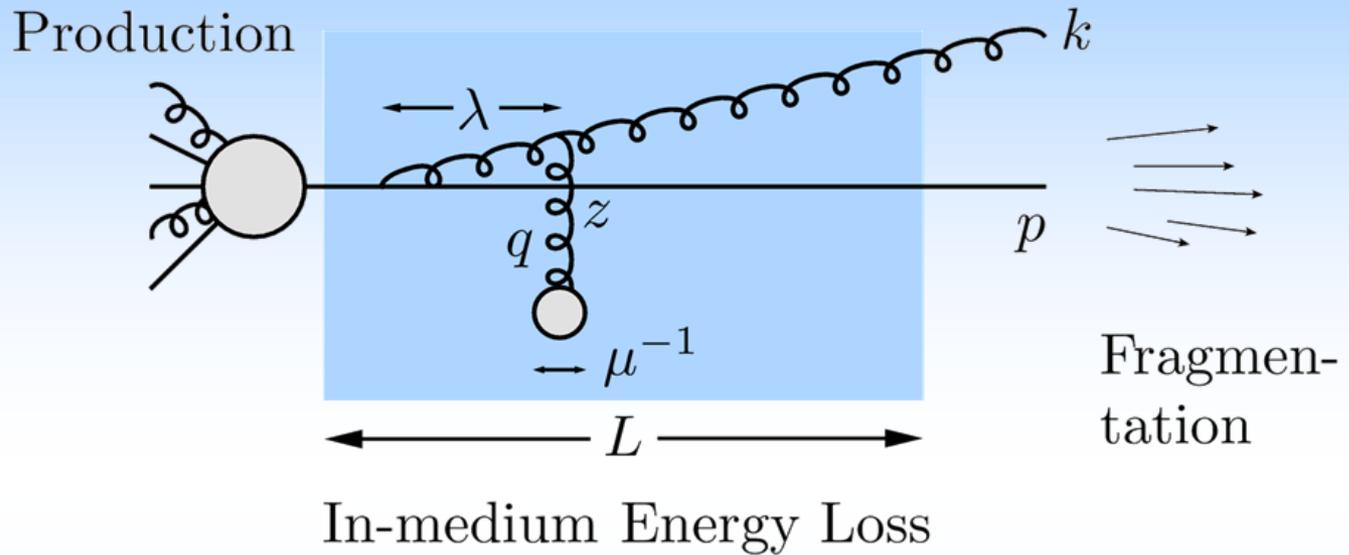
What to Do, What to Do?

- Expectation of theorists: settle on single value of extracted parameter
 - Within pQCD: very diff. assumptions about the relevant physics
 - *Shouldn't* expect a single density, q hat, etc. from different physics
 - Treat all theories as equal?
 - Ex: AdS/CFT
- Need experimental observables to determine the relevant physics
 - Use suppression to learn about QCD, QGP
 - Brick problem: see Marta's talk

Need for Theoretical Uncertainty

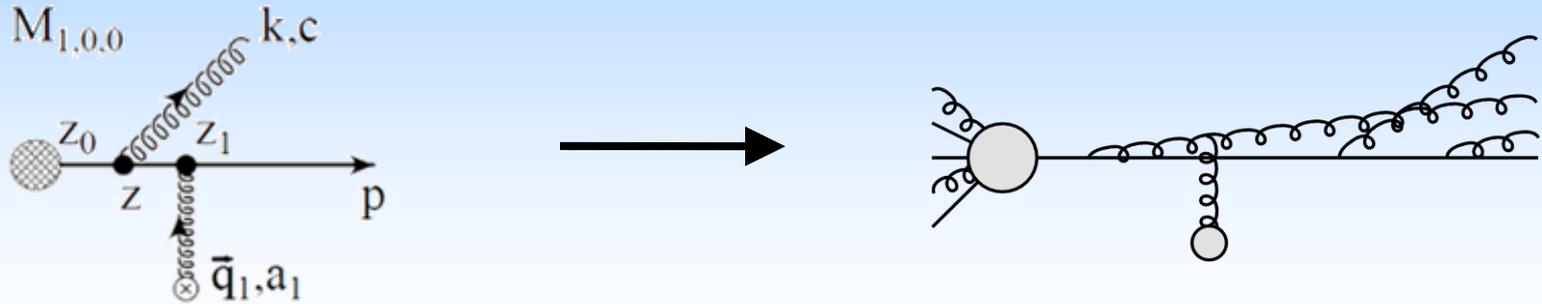
- **Want to rigorously:**
 - falsify theories
 - quantify medium
- **Therefore need:**
 - Precise observables
 - Precise theory
- **Distinguish between systematic uncertainties:**
 - between formalisms
 - Due to diff. physics assumptions
 - within formalisms
 - Due to simplifying approximations
- **Focus specifically on opacity expansion**
 - GLV; ASW-SH (never been compared with data)

Energy Loss



- $R_{AA} \sim \int (1-\epsilon)^n P(\epsilon) d\epsilon$
 - $E_f = (1-\epsilon)E_i$
- Opacity expansions finds single inclusive gluon emission spectrum
 - $dN_g/dxdk_T dq_T$

Poisson Convolution



Gyulassy, Levai, and Vitev NPB594 (2001)

- Find $P(\epsilon)$ by convolving dN_g/dx
 - Approximates probabilistic multiple gluon emission
 - assume independent emissions

Opacity Expansion Calculation

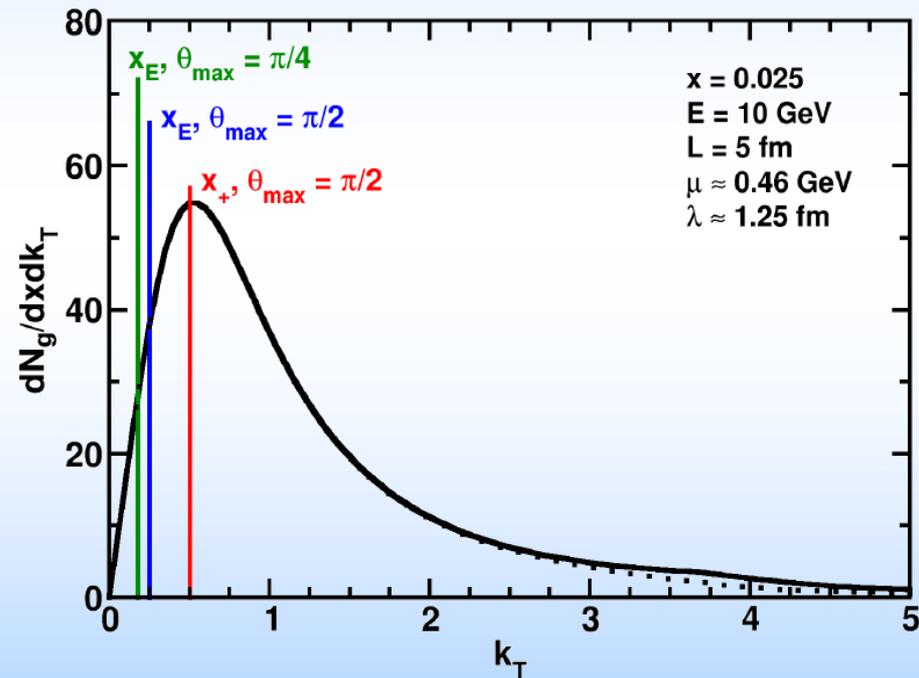
- Want to find dN_g/dx
 - Make approximations to simplify derivation
 - Small angle emission: $k_T \ll xE$
 - Note: ALL current formalisms use collinear approximation
 - Derived $dN_g/dxdk_T$ violates collinear approx

$$x \frac{dN_g}{dx} = \frac{C_R \alpha_s L}{\pi \lambda} \int_0^{q_{\max}^2} \frac{2q^2 \mu^2 dq^2}{(4xE\hbar c/L)^2 + (q^2 + \beta^2)^2} \\ \times \int_0^{k_{\max}^2} \frac{dk^2}{k^2 + \beta^2} \frac{k^2(k^2 - q^2 + \mu^2) - \beta^2(k^2 - q^2 - \mu^2)}{((k - q)^2 + \mu^2)^{3/2} ((k + q)^2 + \mu^2)^{3/2}} + \mathcal{O}\left((k_T/xE)^2\right)$$

- Both IR and UV safe
- Enforce small angle emission through UV cutoff in k_T

Uncertainty from Collinear Approx

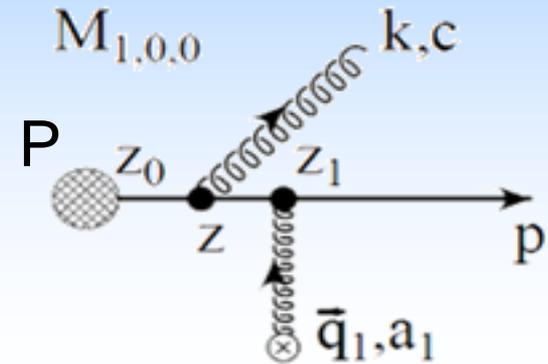
- Derived $dN_g/dxdk_T$ maximally violates collinear approximation
 - dN_g/dx depends sensitively on k_T cutoff
 - Despite UV safety
 - For effect on extracted prop., must understand x
 - Discovered through TECHQM Brick Problem



WAH and B Cole, arXiv:0910.1823

ASW-SH Definition of x

- ASW-SH: x_E
 - Minkowski coords



$$P = (E, E, 0)$$

$$p = ((1 - x_E)E, \sqrt{((1 - x_E)E)^2 - (q_{\perp} - k_{\perp})^2}, q_{\perp} - k_{\perp})$$

$$k = (x_E E, \sqrt{(x_E E)^2 - k_{\perp}^2}, k_{\perp})$$

- Always on-shell

GLV Definition of x

- GLV: x_+

- Light-cone coords

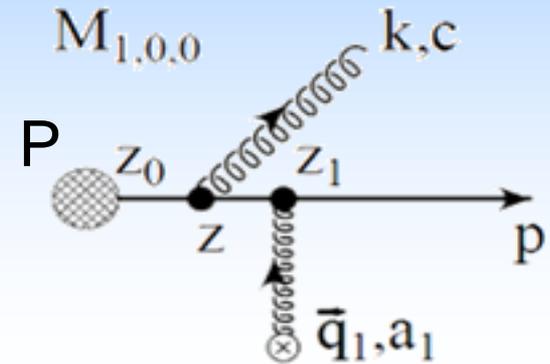
$$P = (E, E, 0, 0) = [E^+, 0, 0]$$

$$k = [x_+ E^+, \frac{k_\perp^2}{x_+ E^+}, \mathbf{k}_\perp]$$

$$p = [(1 - x_+) E^+, \frac{(\mathbf{q}_\perp - \mathbf{k}_\perp)^2}{(1 - x_+) E^+}, \mathbf{q}_\perp - \mathbf{k}_\perp]$$

- Always on-shell

$$k = \left(\frac{1}{2} x_+ E^+ \left[1 + \left(\frac{\mathbf{k}_\perp}{x_+ E^+} \right)^2 \right], \frac{1}{2} x_+ E^+ \left[1 - \left(\frac{\mathbf{k}_\perp}{x_+ E^+} \right)^2 \right], \mathbf{k}_\perp \right)$$



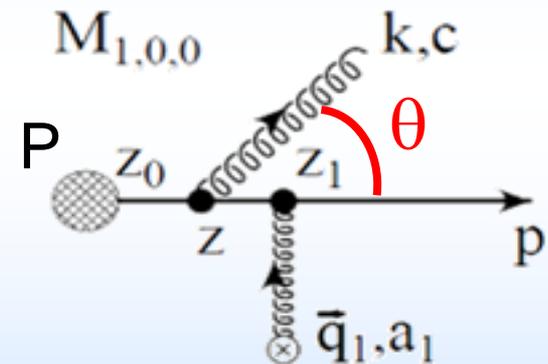
Coordinate Transformations

$$x_+ = \frac{x_E}{2} \left(1 + \sqrt{1 - \left(\frac{\mathbf{k}_\perp}{x_E E} \right)^2} \right) \quad x_E = x_+ \left(1 + \left(\frac{\mathbf{k}_\perp}{x_+ E^+} \right)^2 \right)$$

– Same in the limit $k_\perp/xE \rightarrow 0!$

- UV cutoff given by restricting maximum angle of emission

$$k_{\max} = \begin{cases} x_+ E^+ \tan(\theta_{\max}/2), & x = x_+, \\ x_E E \sin(\theta_{\max}), & x = x_E. \end{cases}$$



- Previous comparisons with data took $\theta_{\max} = \pi/2$
- Vary θ_{\max} to estimate systematic theoretical uncertainty

Jacobians

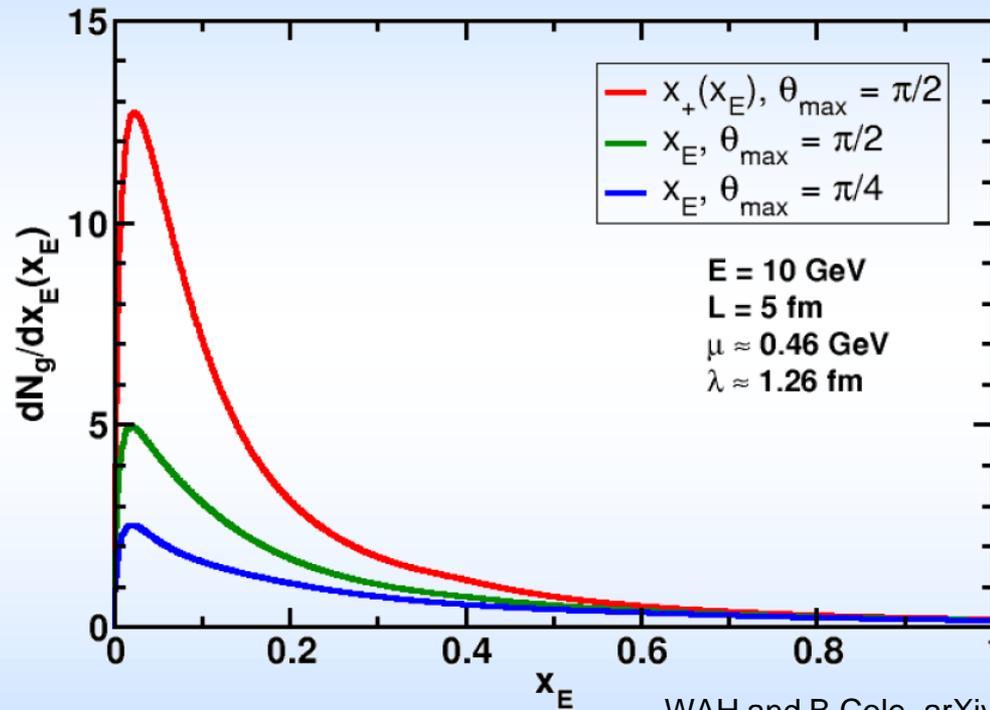
- ϵ is fraction of longitudinal momentum
 - Need dN_g/dx_E to find $P(\epsilon)$
 - A Jacobian is required for $x = x_+$ interpretation

$$\frac{dN_g^J}{dx_E}(x_E) \equiv \int^{x_E E \sin(\theta_{\max})} dk_T \frac{dx_+}{dx_E} \frac{dN_g}{dx_+ dk_T}(x_+(x_E)),$$

$$\frac{dx_+}{dx_E} = \frac{1}{2} \left[1 + \left(1 - \left(\frac{k_T}{x_E E} \right)^2 \right)^{-1/2} \right].$$

Rad. Gluon Kin. Sensitivities

- UV



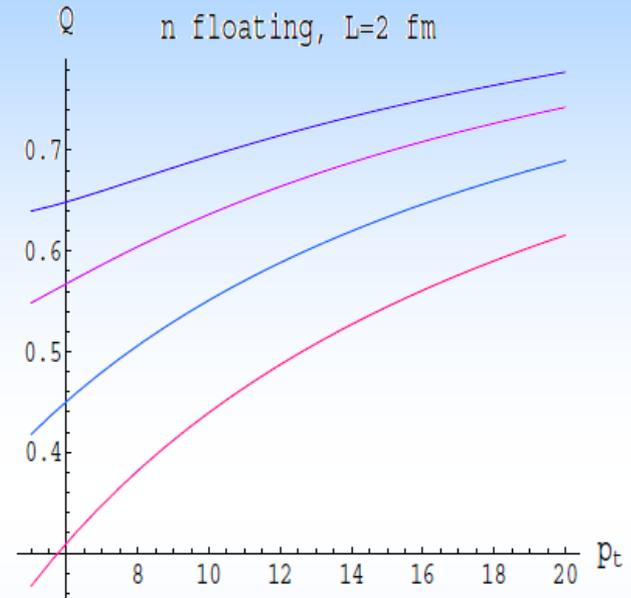
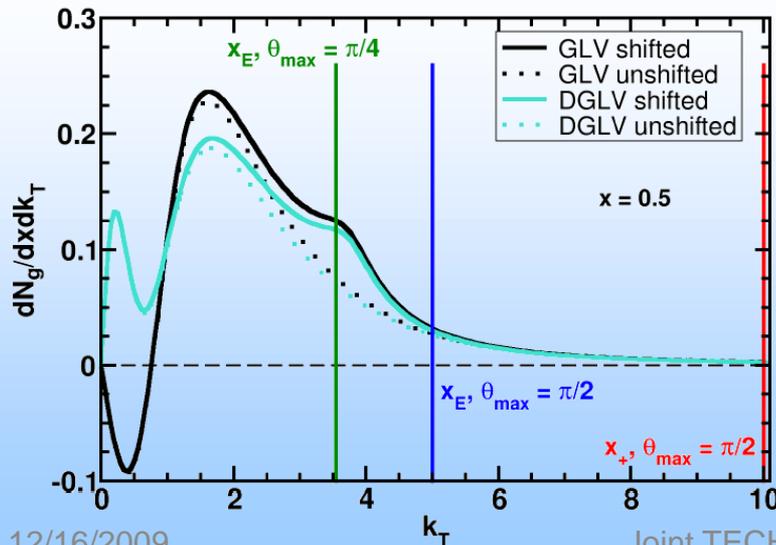
WAH and B Cole, arXiv:0910.1823

- What about IR?

Collinearity and Gluon Mass

- Massless gluons:
 - Large IR cutoff sensitivity
- Gluons with thermal mass

$$\tau_{\text{coh}} = \frac{2xE}{q_T^2 + \beta^2}$$



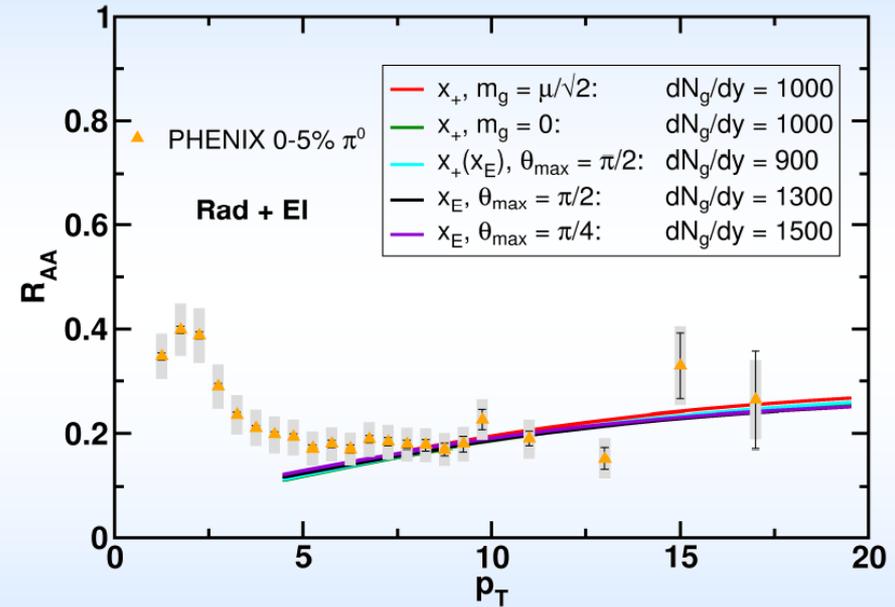
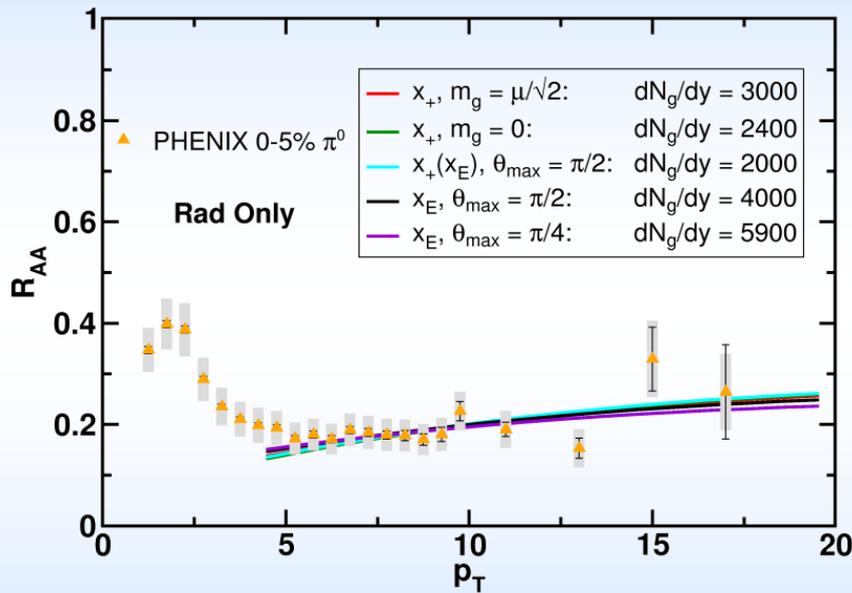
BDMS, JHEP 0109 (2001)

m_g^2

Larger x better respects $k_T \ll xE$

Results

- Quantitatively compare to PHENIX data

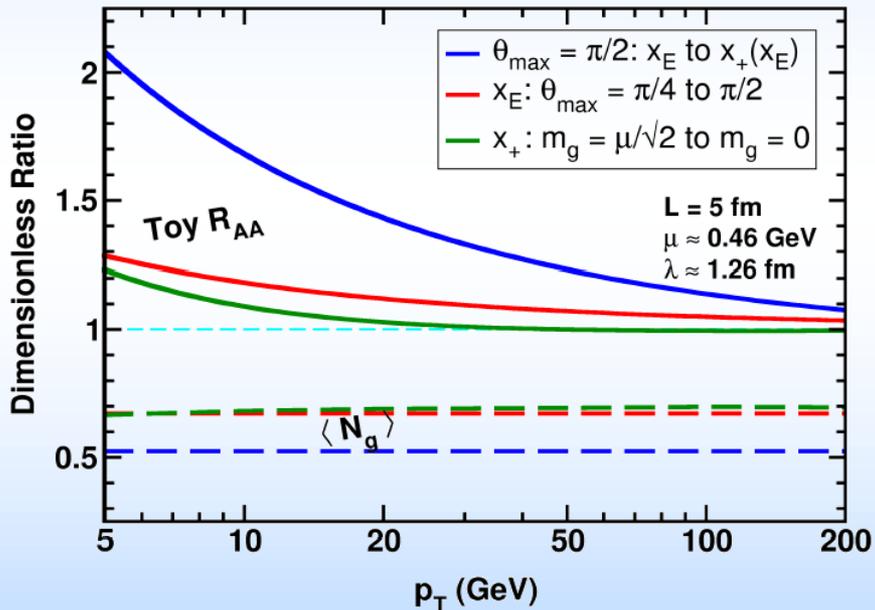


WAH and B Cole, arXiv:0910.1823

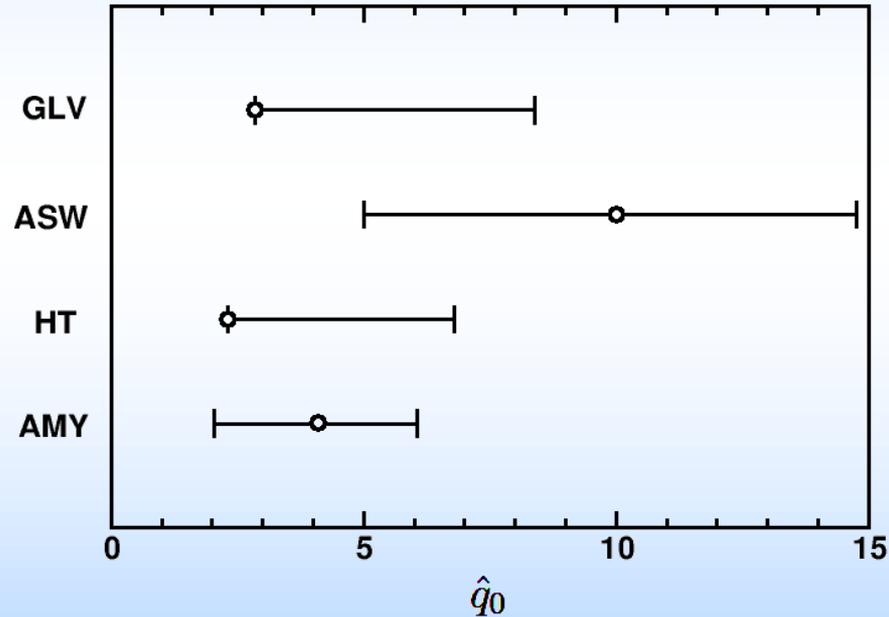
– Assumed infinite Elastic precision

Parton Energy Dependence

- Dependence on parton energy



- Uncertainty on \hat{q}



– Assume all formalisms equally affected

WAH and B Cole, arXiv:0910.1823

Conclusions

- Use different physics, expect different parameters
- $dN_g/dxdk_T$ drastically violates collinear approx.
 - Large angle radiation is important
 - Not under theoretical control
 - Leads to ~200% systematic theoretical uncertainty in extracted medium density at RHIC
 - Extracted parameters from pQCD formalisms likely consistent within uncertainty
- Thermal gluon mass effects small
- LHC
 - R_{AA} insensitive to collinear approx. at ultra-high- p_T
 - $\langle N_g \rangle$ uncertainty E-independent
 - Large cone jets not under theoretical control