

Discovery  
in  
Drell-Yan  
at the  
LHC

Neil Christensen

University of Wisconsin - Madison

in collaboration with

Cheng-Wei Chiang, Gui-Jun Ding and Tao Han

to be published soon!

SM

???

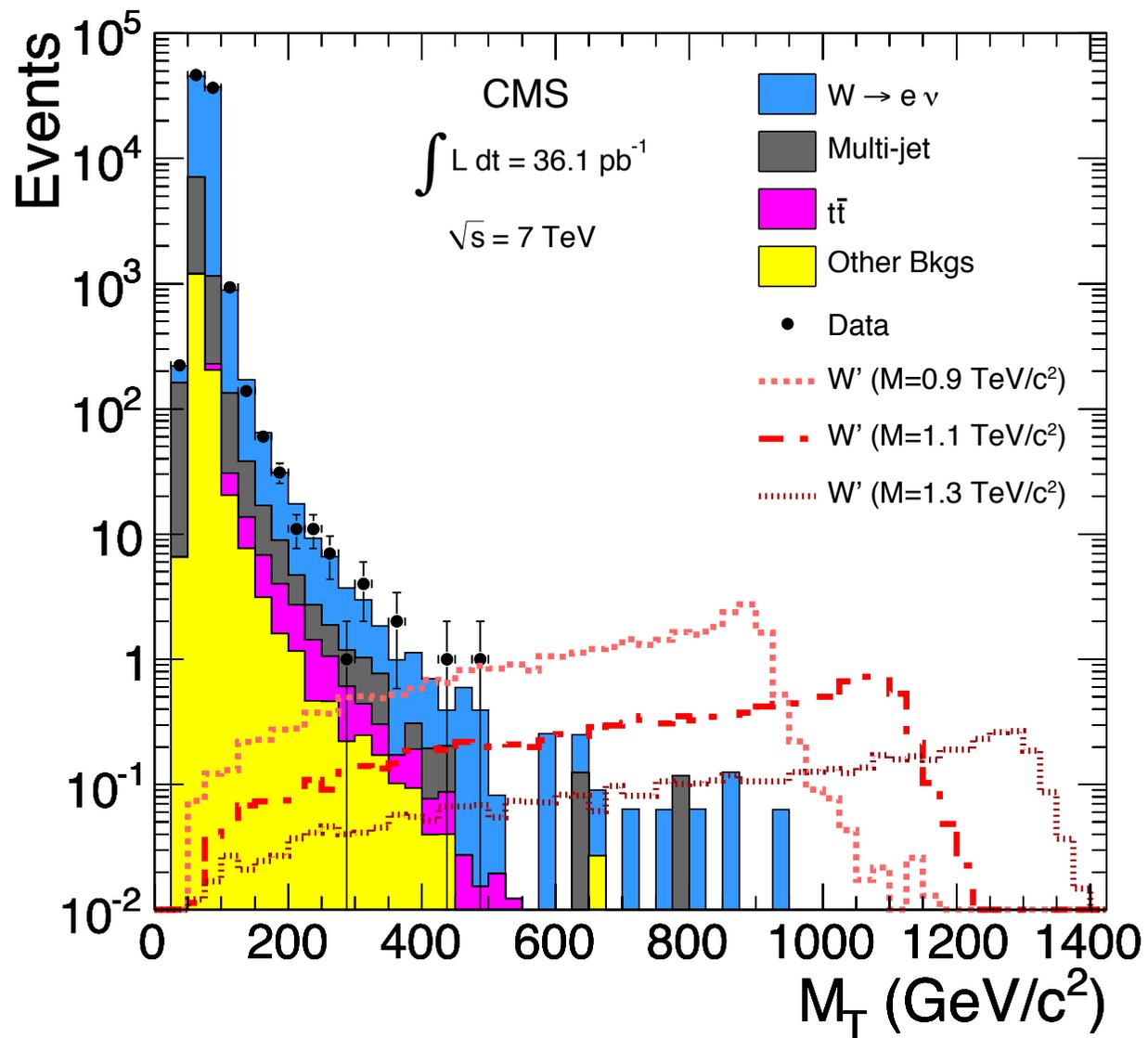
Supersymmetry

Extra  
Dimensions

Little Higgs

Higgsless

New Strong  
Dynamics



SM

???

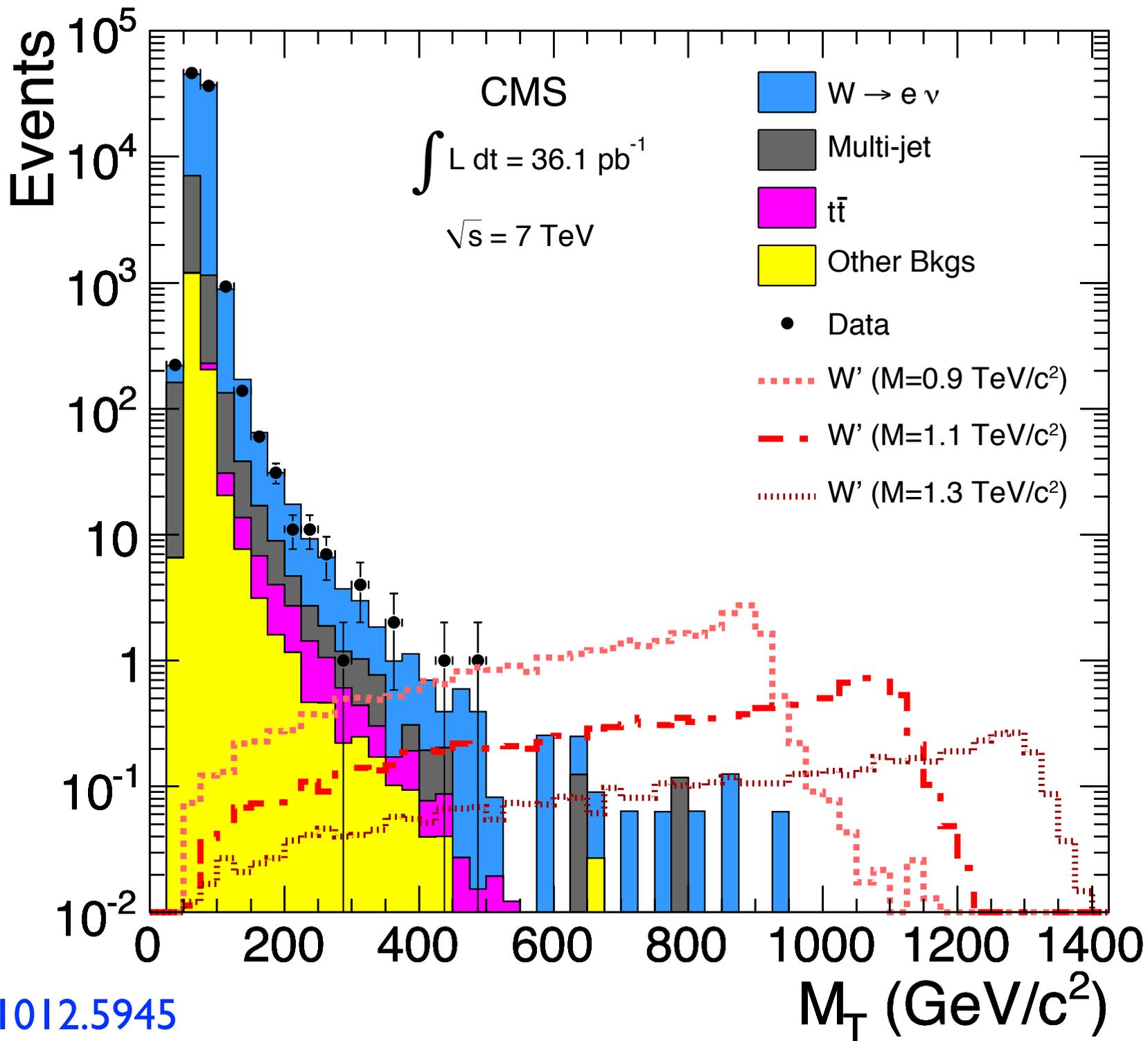
Supersymmetry

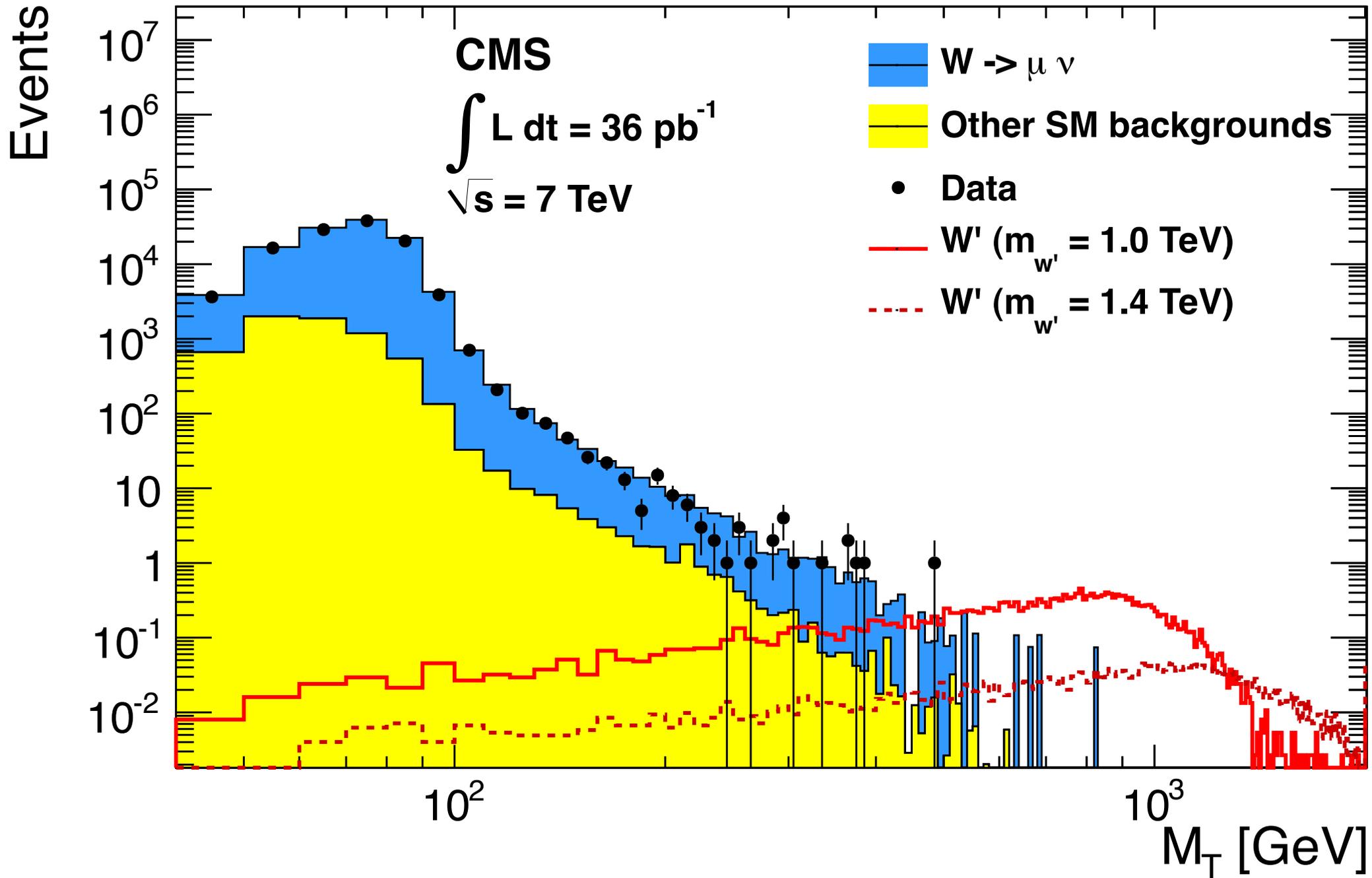
Extra  
Dimensions

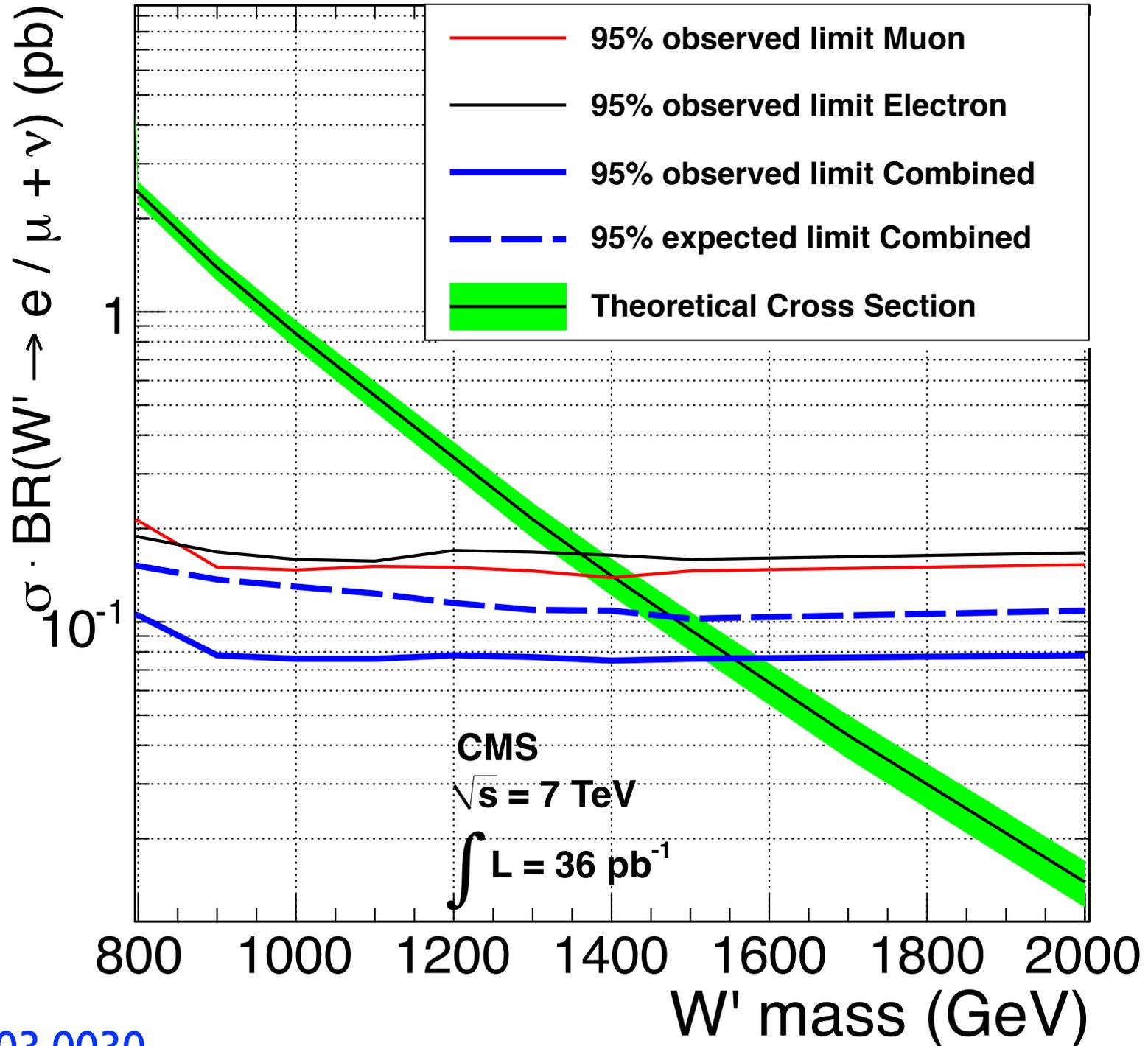
Little Higgs

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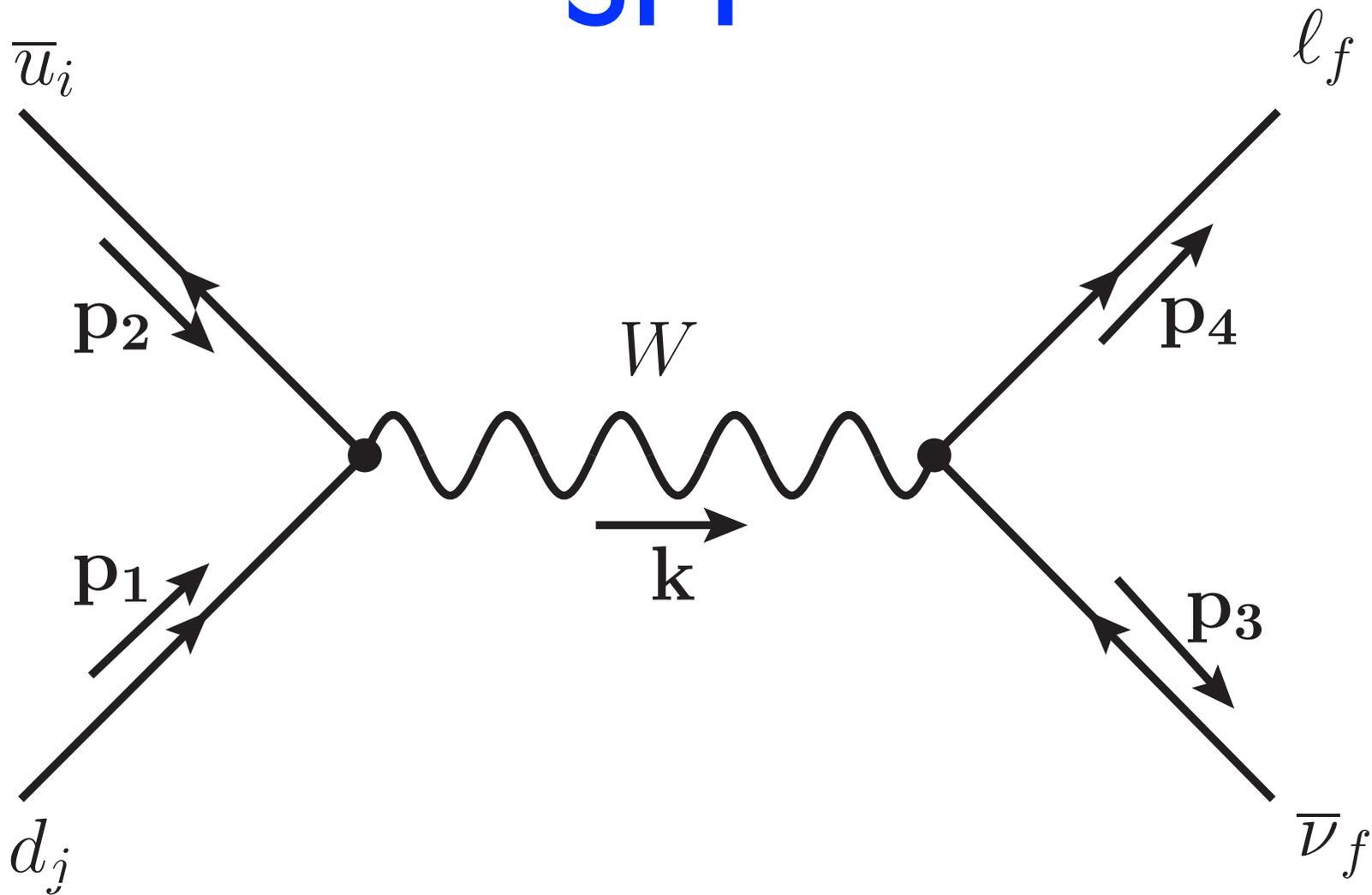




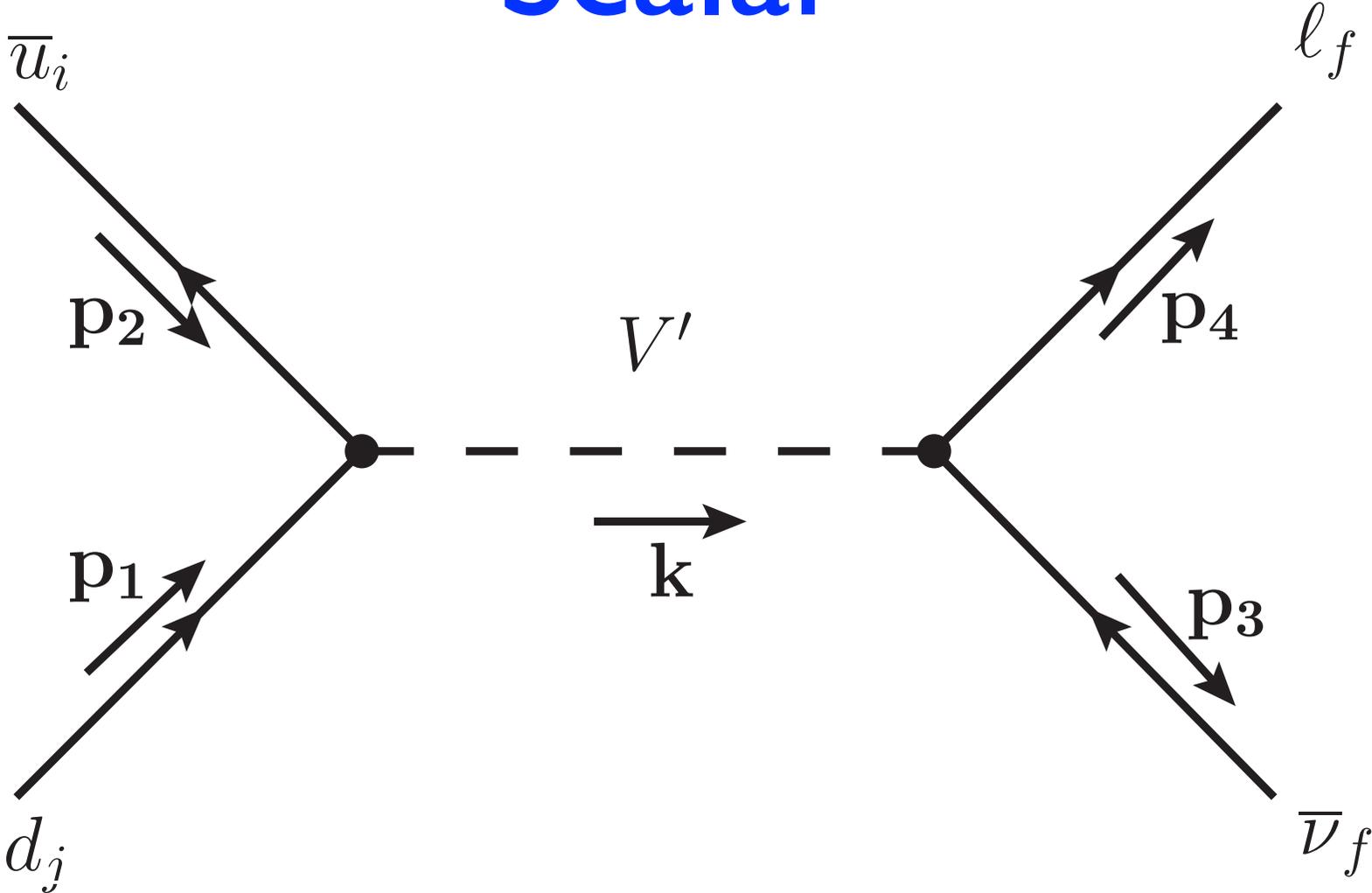
This channel is very simple.

A model independent analysis can be done.

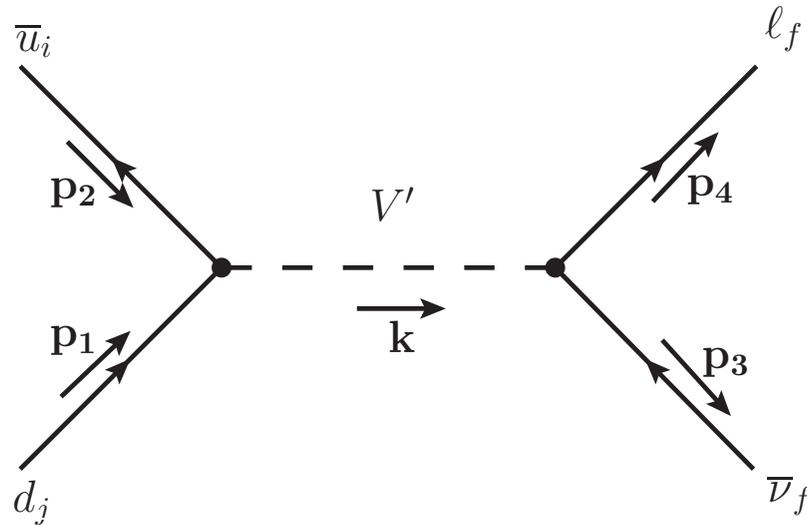
SM



# Scalar

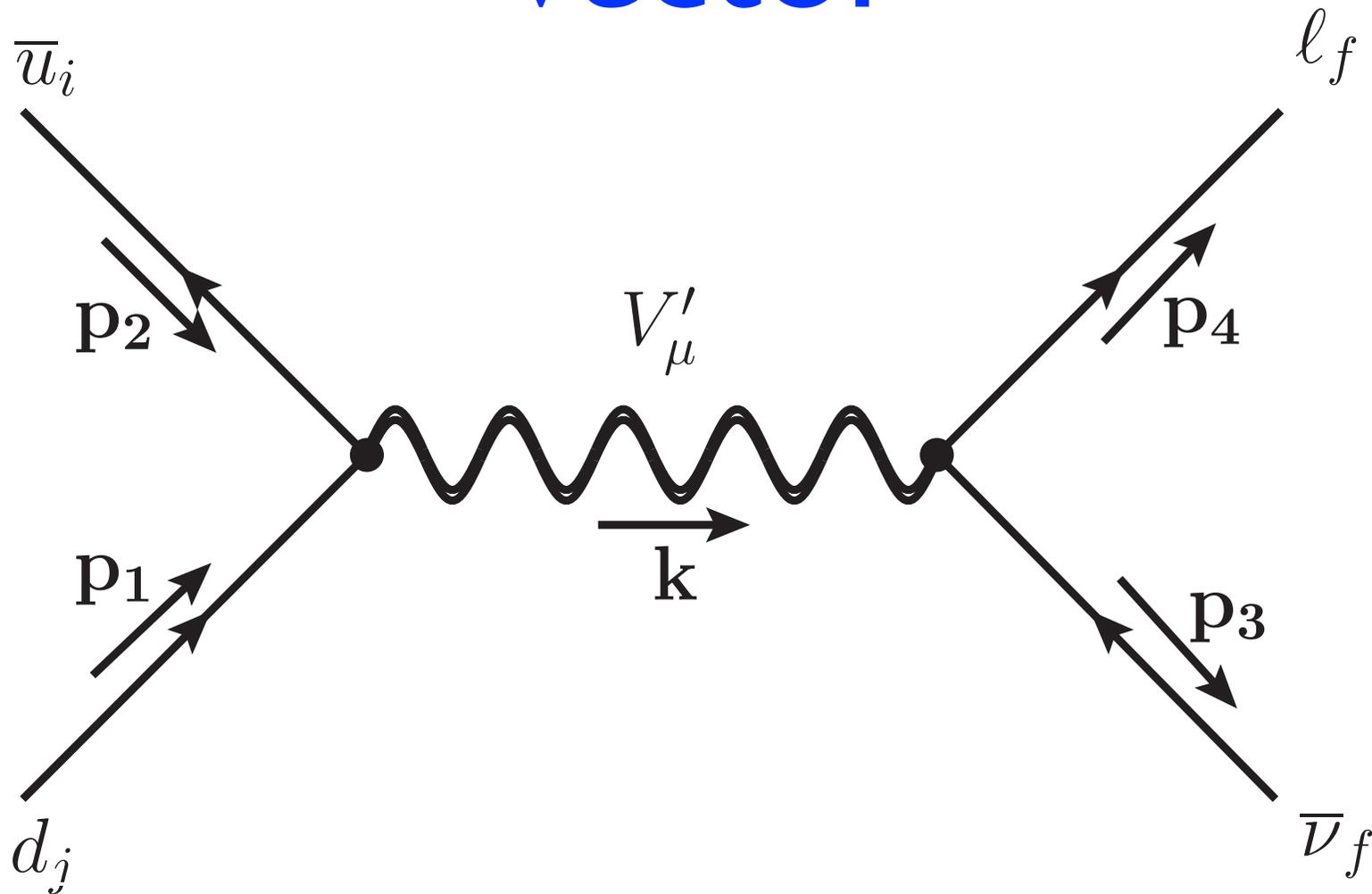


# Scalar

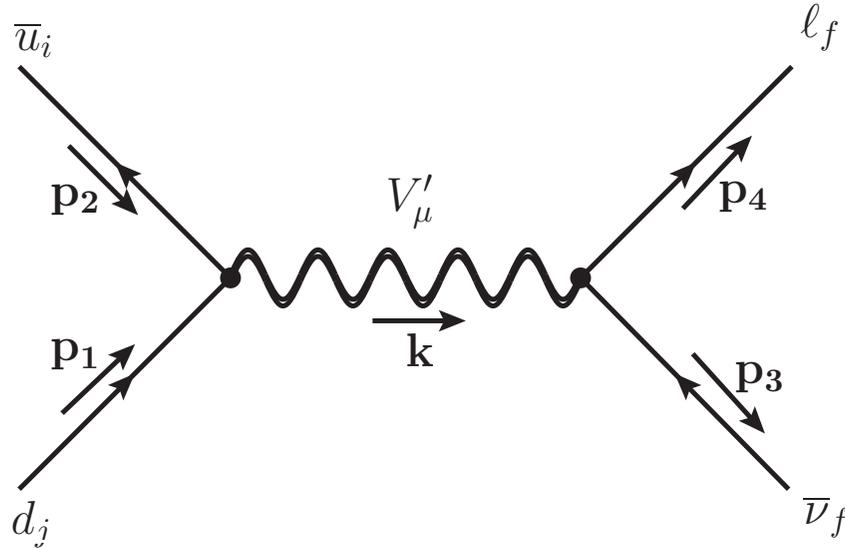


$$\begin{aligned} \mathcal{L}_{V'}^{\text{scalar}} &= \bar{u}_i \left( h_{Sij}^q + ih_{Pij}^q \gamma_5 \right) d_j V' + \text{h.c.} \\ &\quad + \bar{\nu}_i \left( h_{Sij}^\ell + ih_{Pij}^\ell \gamma_5 \right) \ell_j V' + \text{h.c.} \end{aligned}$$

# Vector

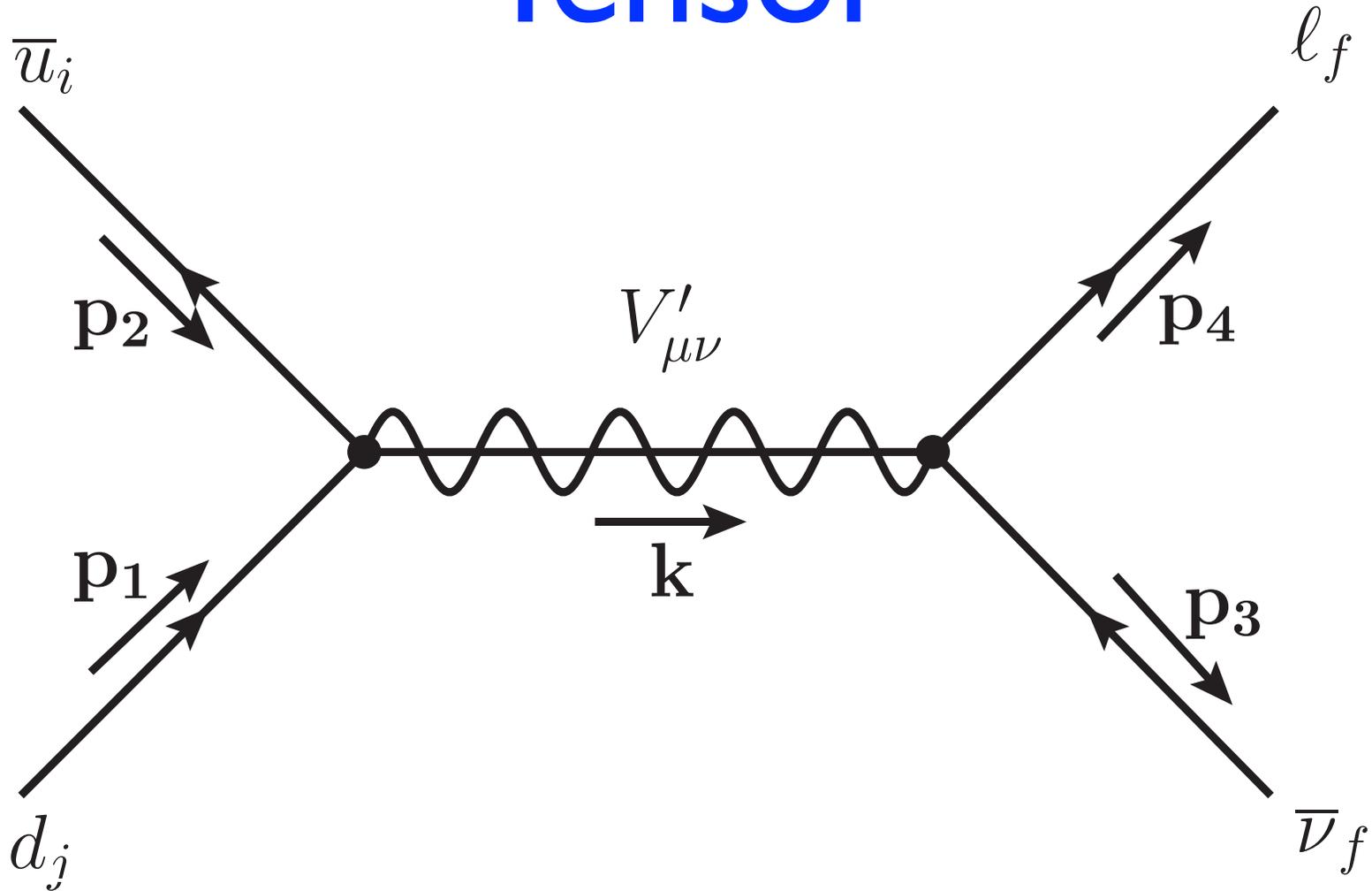


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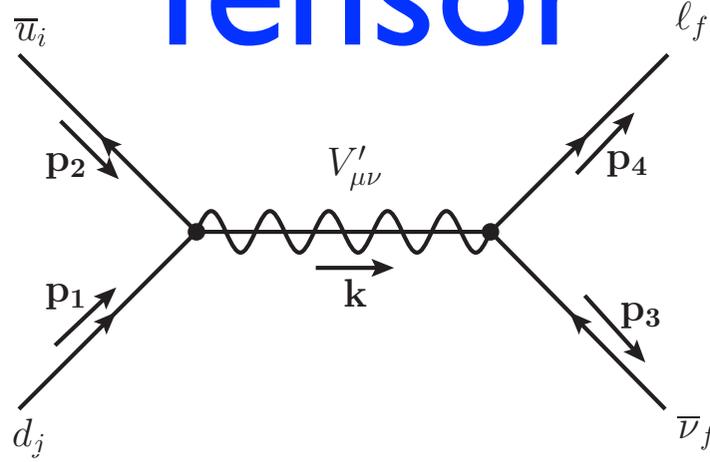


$$\begin{aligned} \mathcal{L}_{V'}^{\text{vector}} &= \bar{u}_i \gamma^\mu \left( h_{Vij}^q + h_{Aij}^q \gamma_5 \right) d_j V'_\mu + \text{h.c.} \\ &+ \bar{v}_i \gamma^\mu \left( h_{Vij}^\ell + h_{Aij}^\ell \gamma_5 \right) \ell_j V'_\mu + \text{h.c.} \end{aligned}$$

# Tensor



# Tensor



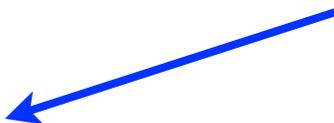
$$\begin{aligned}
 \mathcal{L}_{V'}^{\text{tensor}} = & \frac{i}{\Lambda} \left[ \bar{u}_i \left( h_{Tij}^q - h_{ATij}^q \gamma_5 \right) \left( \gamma^\mu \partial^\nu d_j + \gamma^\nu \partial^\mu d_j \right) V'_{\mu\nu} \right. \\
 & \left. - \left( \partial^\mu \bar{u}_i \gamma^\nu + \partial^\nu \bar{u}_i \gamma^\mu \right) \left( \tilde{h}_{Tij}^q + \tilde{h}_{ATij}^q \gamma_5 \right) d_j V'_{\mu\nu} \right] \\
 & + \frac{i}{\Lambda} \left[ \bar{v}_i \left( h_{Tij}^\ell - h_{ATij}^\ell \gamma_5 \right) \left( \gamma^\mu \partial^\nu \ell_j + \gamma^\nu \partial^\mu \ell_j \right) V'_{\mu\nu} \right. \\
 & \left. - \left( \partial^\mu \bar{v}_i \gamma^\nu + \partial^\nu \bar{v}_i \gamma^\mu \right) \left( \tilde{h}_{Tij}^\ell + \tilde{h}_{ATij}^\ell \gamma_5 \right) \ell_j V'_{\mu\nu} \right] \\
 & + \text{h.c.}
 \end{aligned}$$

Model File

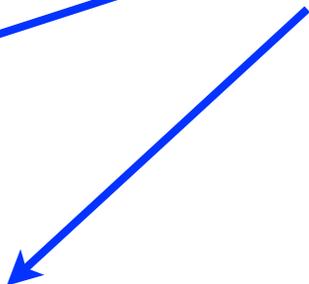


FeynRules

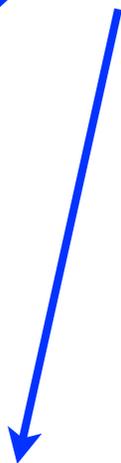
FeynArts



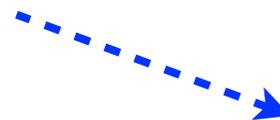
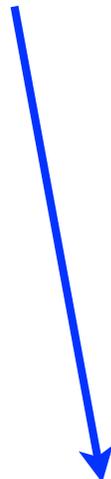
MadGraph  
4/5



CalcHEP

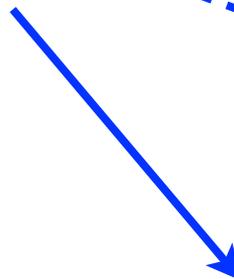


Sherpa



Herwig

Whizard  
1/2

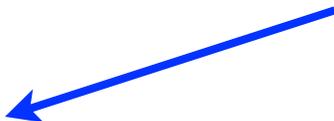


Model File



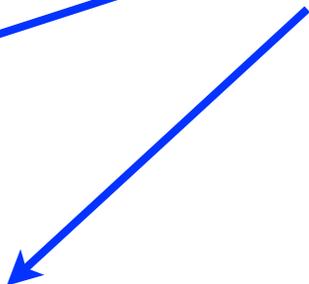
FeynRules

FeynArts

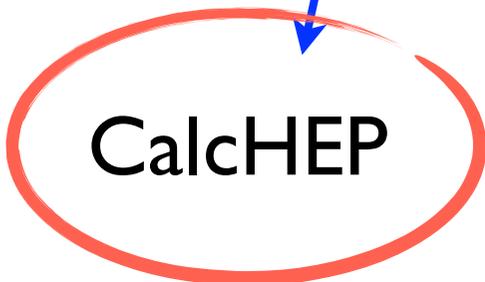


MadGraph

4/5



CalcHEP

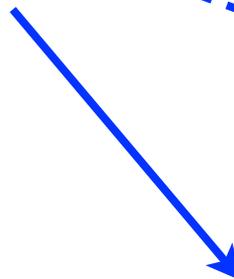


Sherpa

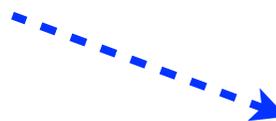


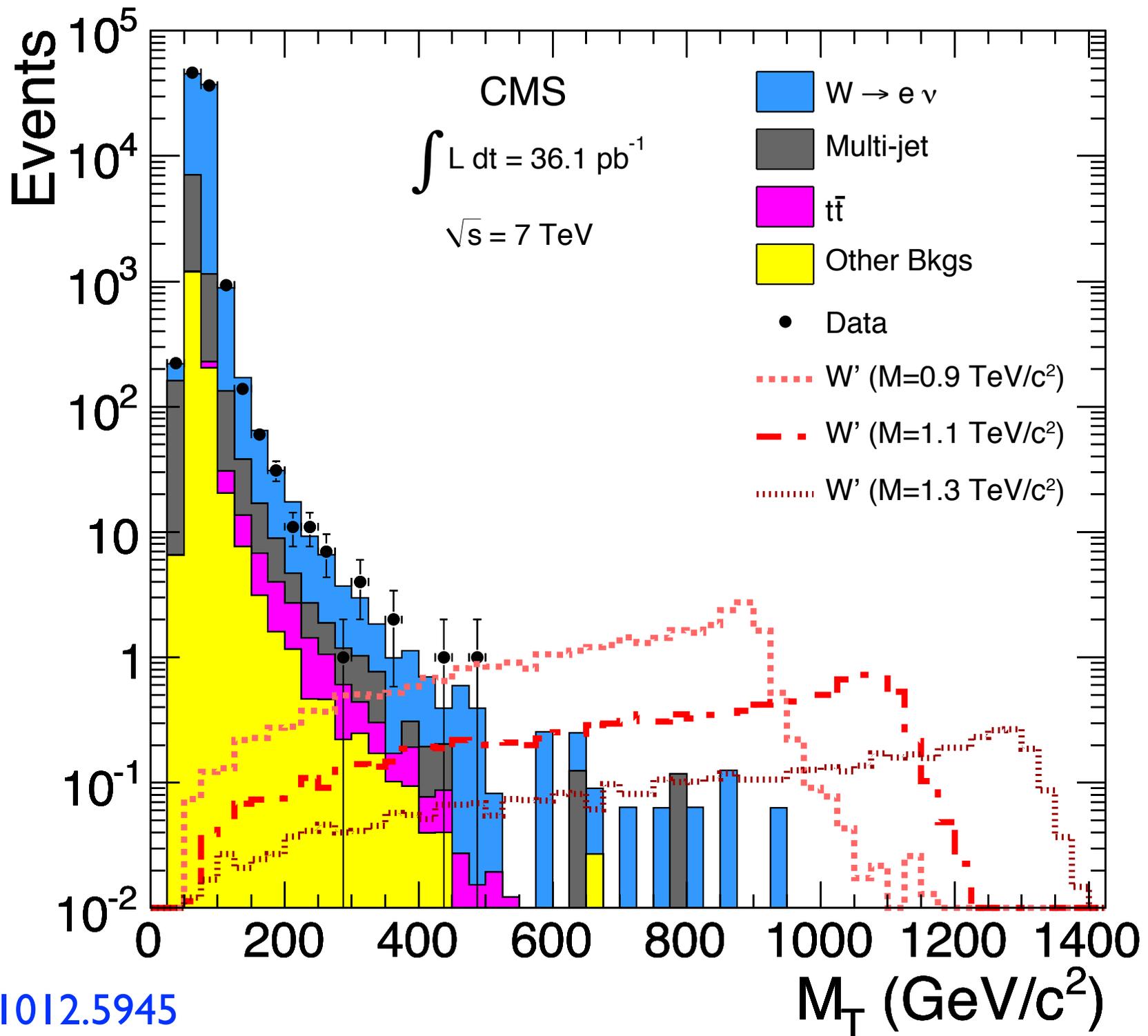
Whizard

1/2

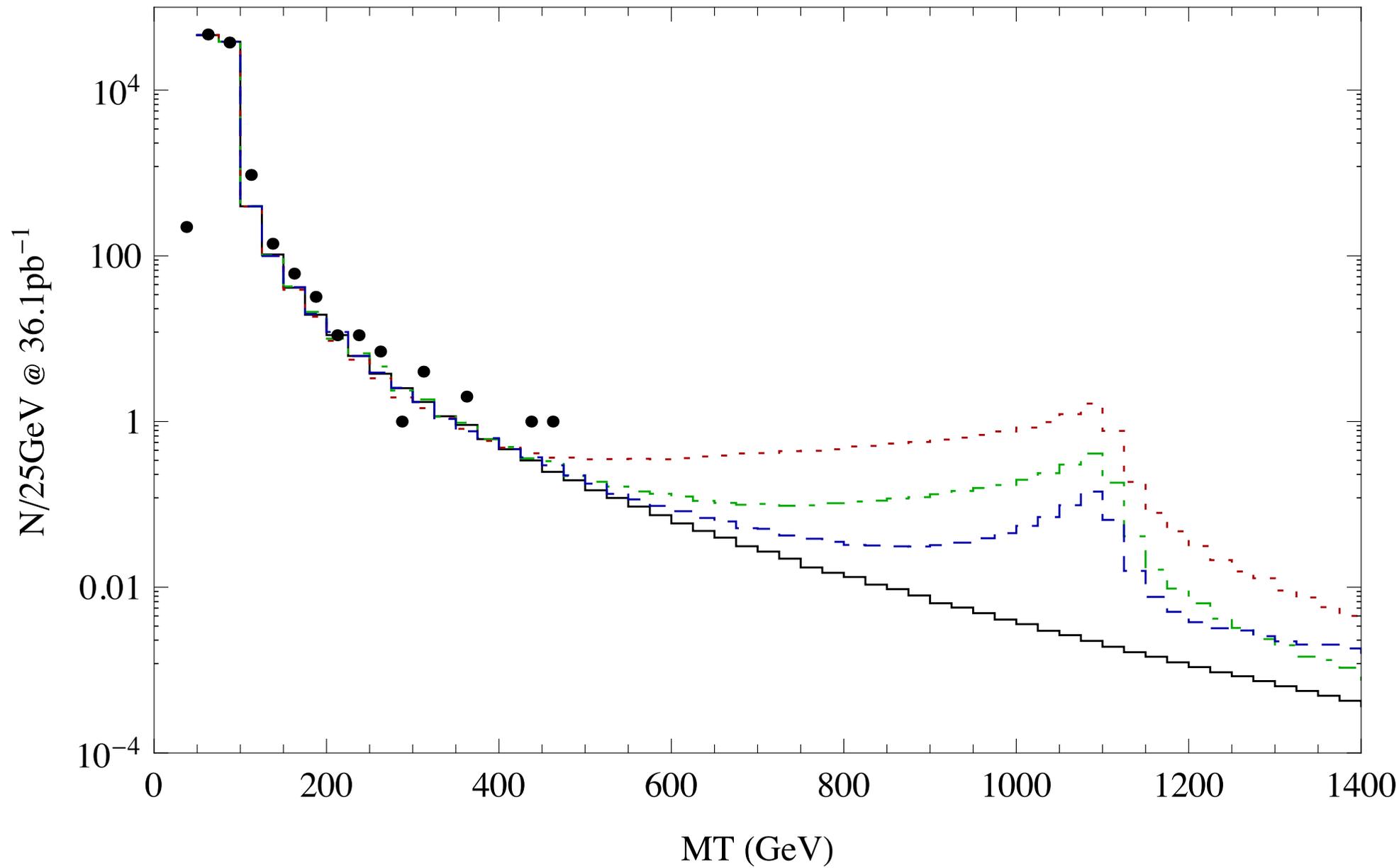


Herwig

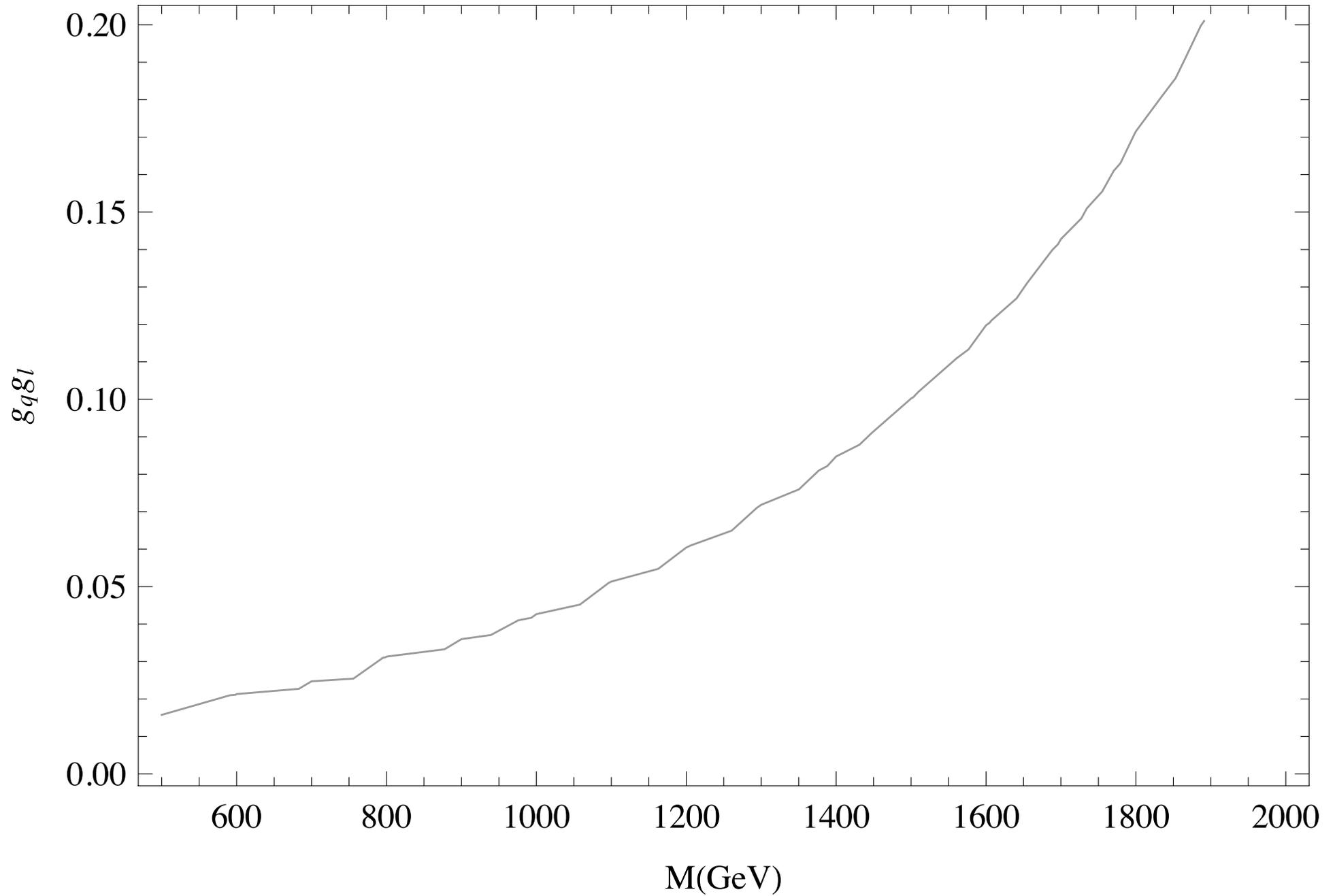


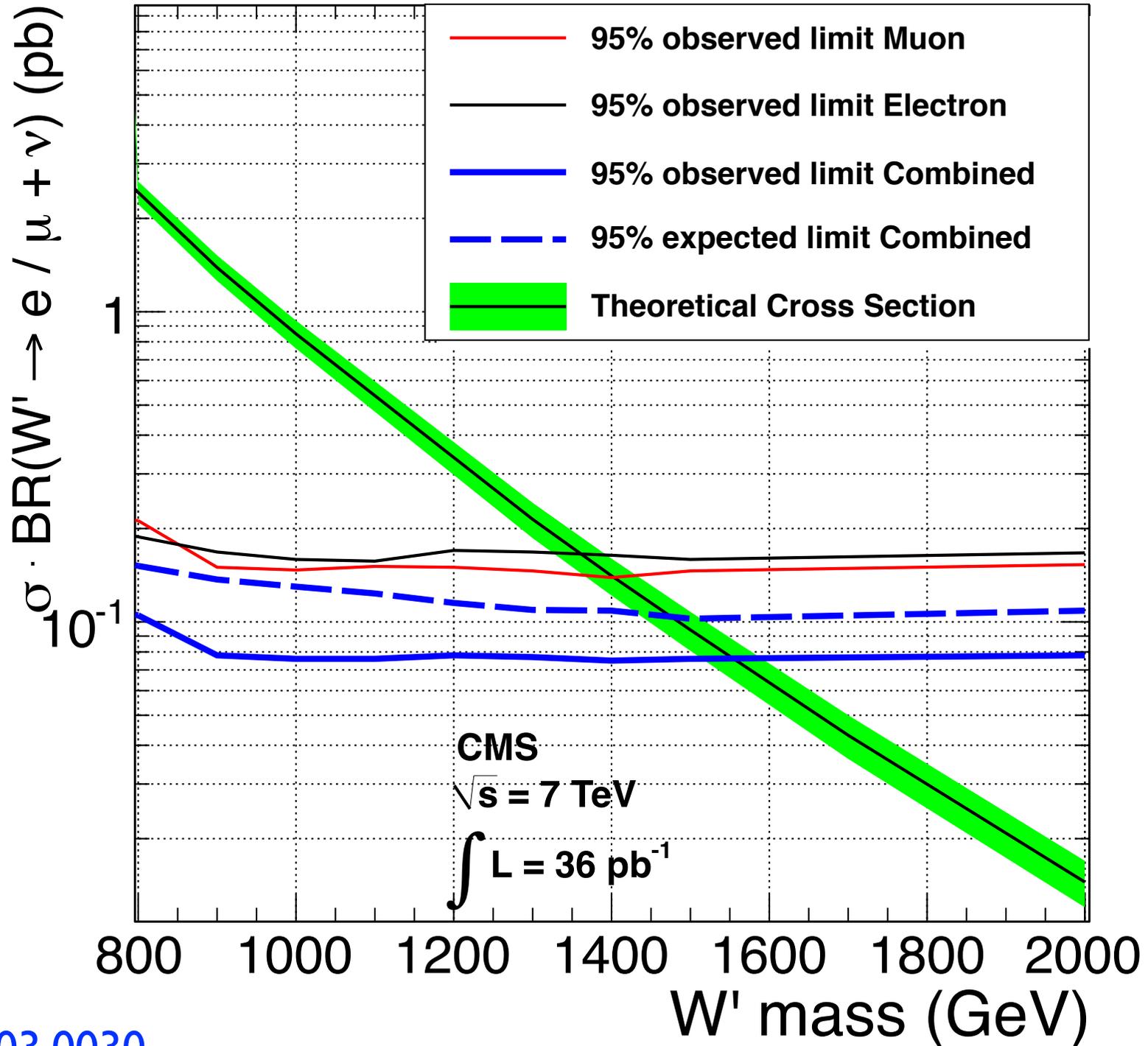


$p, p \rightarrow e, \nu$

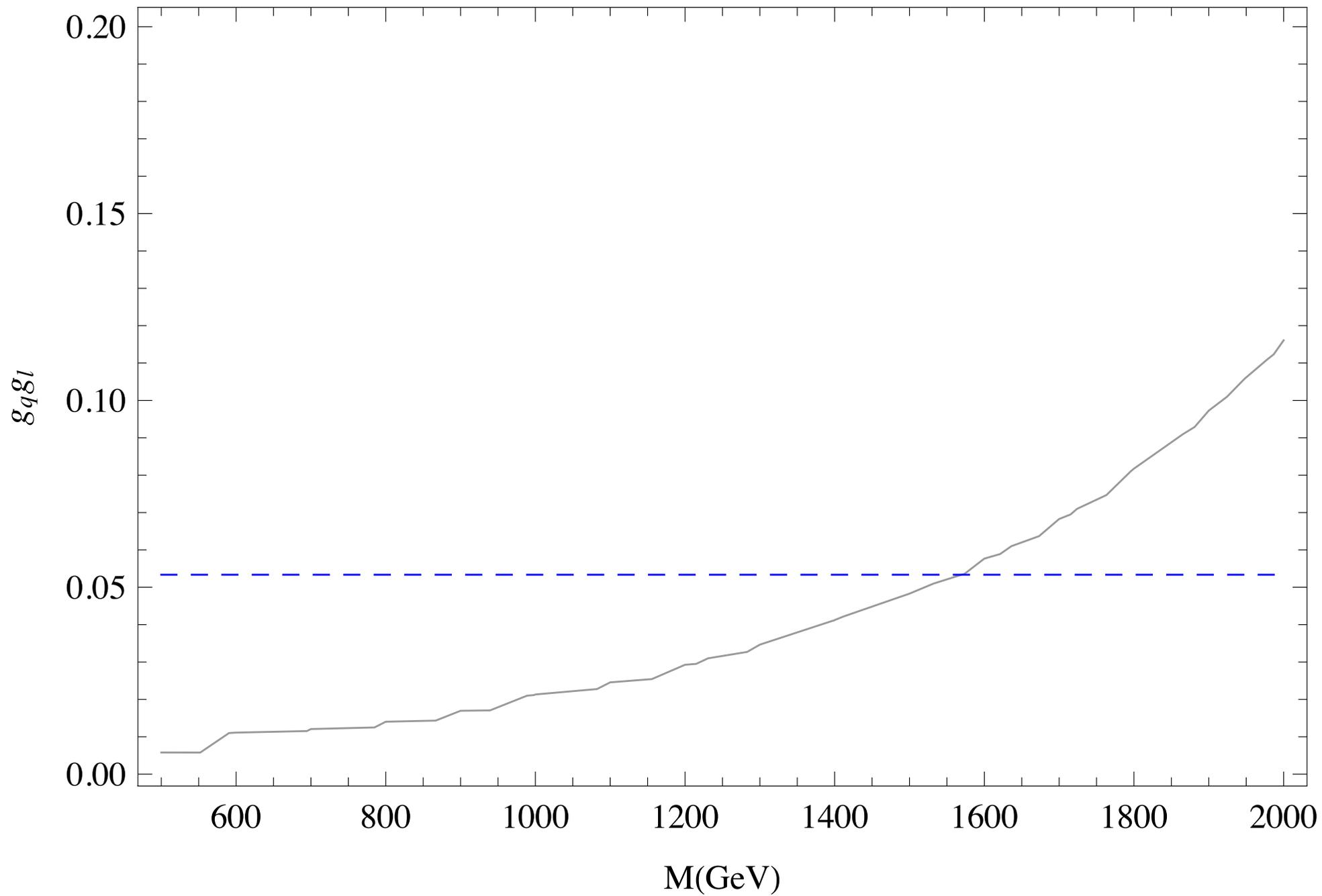


# Scalar

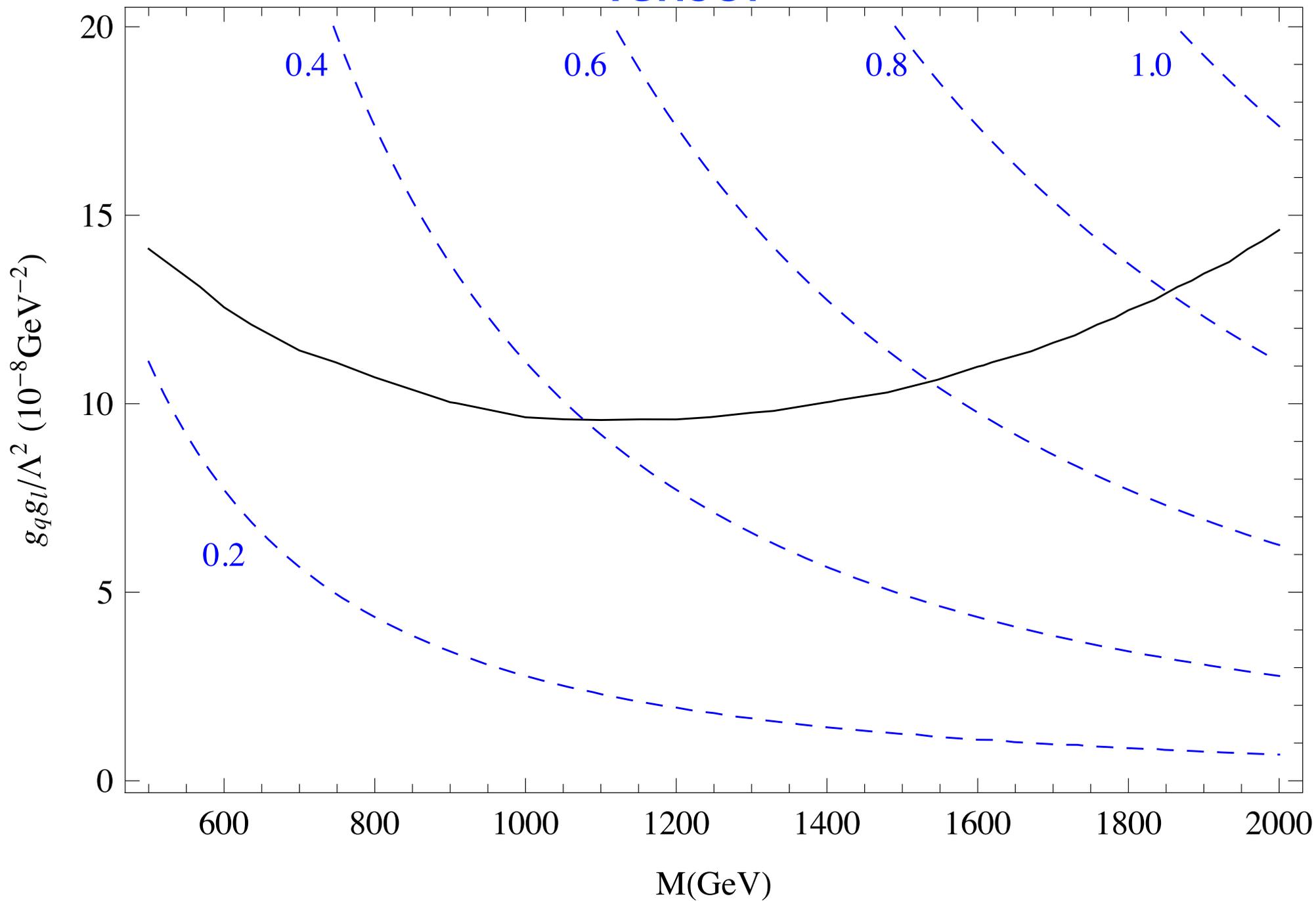




# Vector



# Tensor



What if we actually find a  $M_{\tau}$  peak?

Can we determine the spin of the resonance?

Can we determine the parity violation?

$$(\lambda, -\lambda) \rightarrow (\lambda', -\lambda')$$

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$$\mathcal{M}_W^{\lambda\lambda'} = \mathcal{C}_W^{\lambda\lambda'} \delta_{\lambda', -1} d_{1,1}^1$$

$$\mathcal{M}_V^{\lambda\lambda'} = \mathcal{C}_V^{\lambda\lambda'} d_{1,\lambda\lambda'}^1$$

$$\mathcal{M}_T^{\lambda\lambda'} = \mathcal{C}_{T2}^{\lambda\lambda'} d_{1,\lambda\lambda'}^2 + \mathcal{C}_{T1}^{\lambda\lambda'} d_{1,\lambda\lambda'}^1$$

$$d_{0,0}^0 = 1$$

$$d_{1,\pm 1}^1 = \frac{1}{2}(1 \pm \cos \theta)$$

$$d_{1,\pm 1}^2 = \frac{1}{2}(1 \pm \cos \theta)(2 \cos \theta \mp 1)$$

$$d_{2,\pm 1}^2 = -\frac{1}{2}(1 \pm \cos \theta) \sin \theta$$

$$(\lambda, \lambda) \rightarrow (\lambda', \lambda')$$

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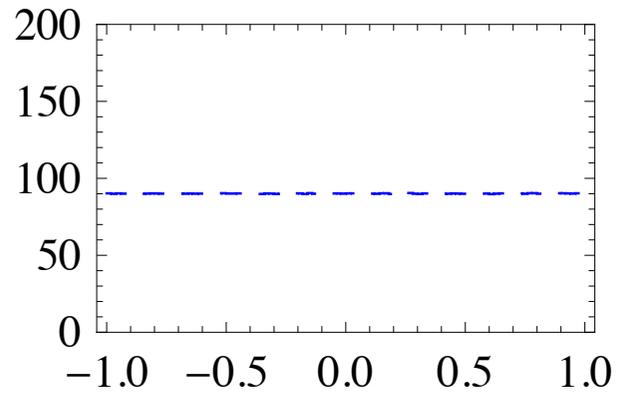
$$\mathcal{M}_S^{\lambda\lambda'} = \mathcal{C}_S^{\lambda\lambda'} d_{0,0}^0$$

What if I knew the full momentum of the neutrino?

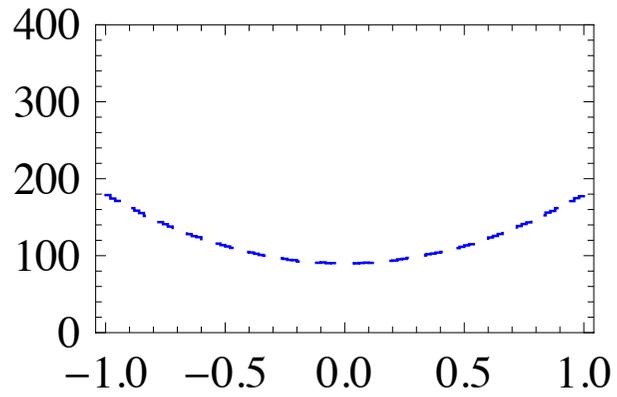
Let's concentrate on the signal for now.

But, we will include the pdfs.

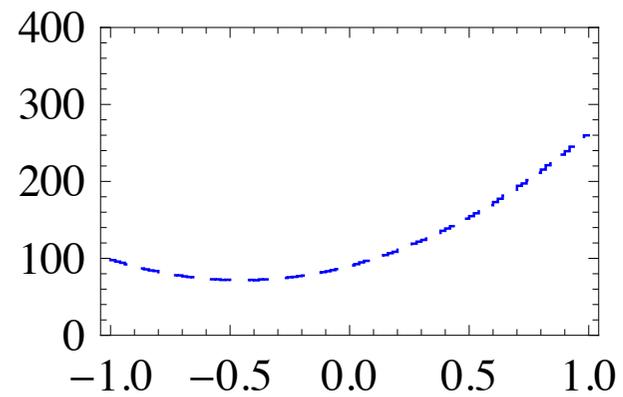
# Scalar



## Vectorial coupling

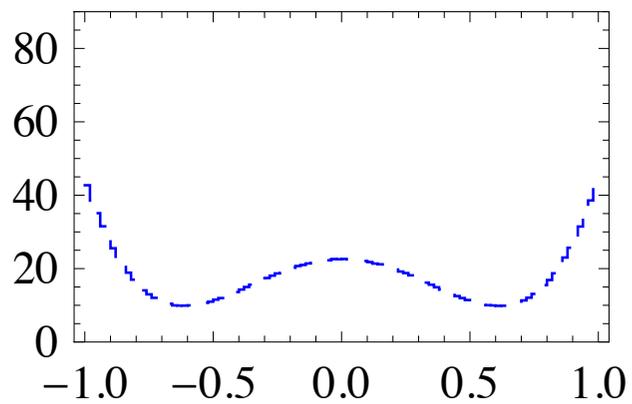


## Left-left coupling

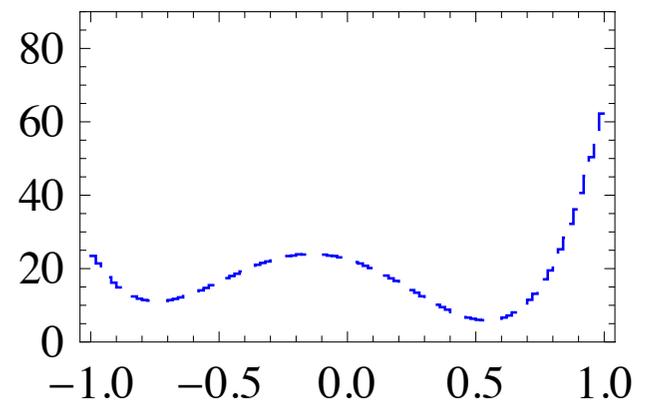


# Vector

## Vectorial coupling



## Left-left coupling



# Tensor

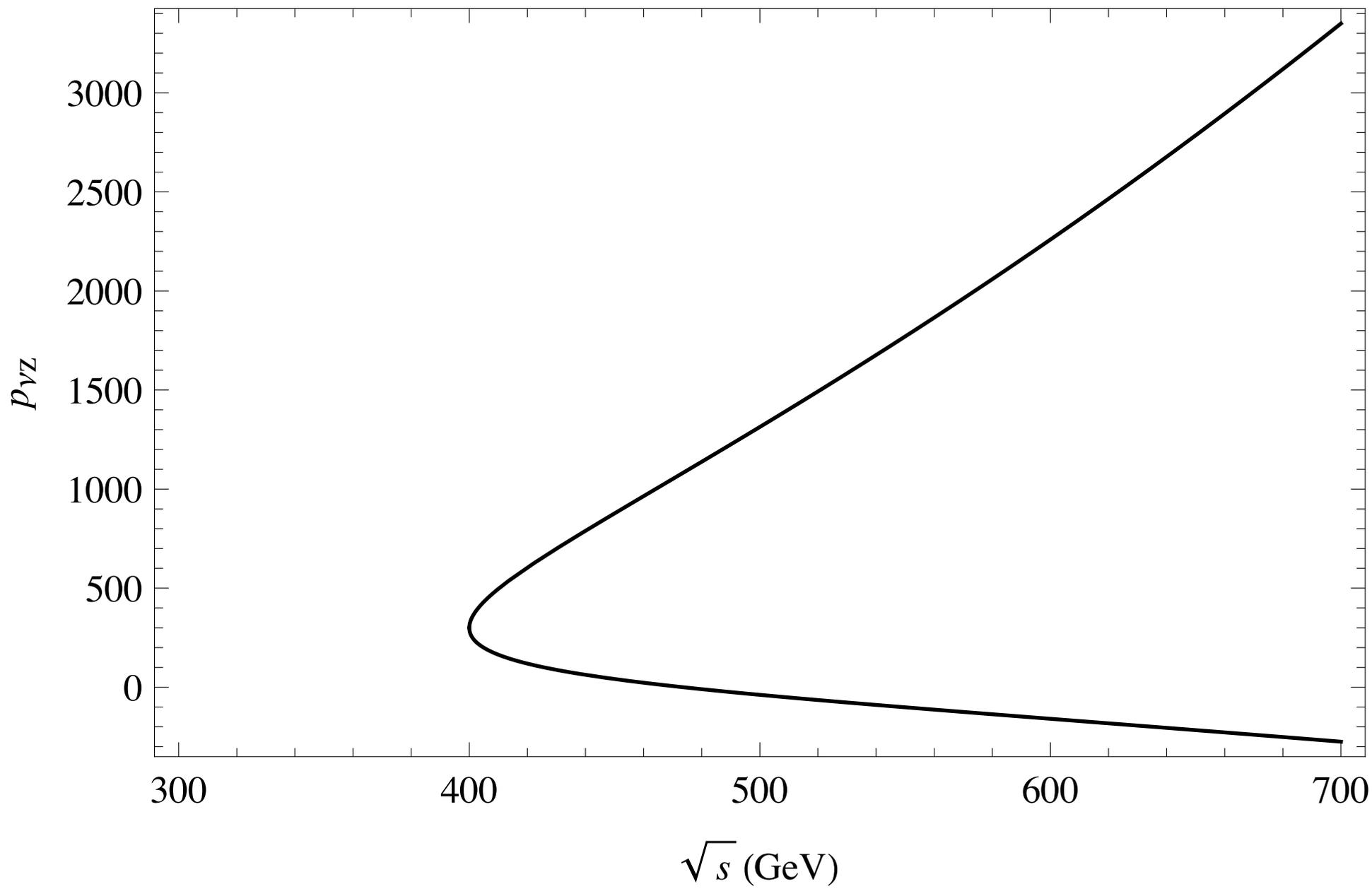
But, we don't know the full momentum of the neutrino!

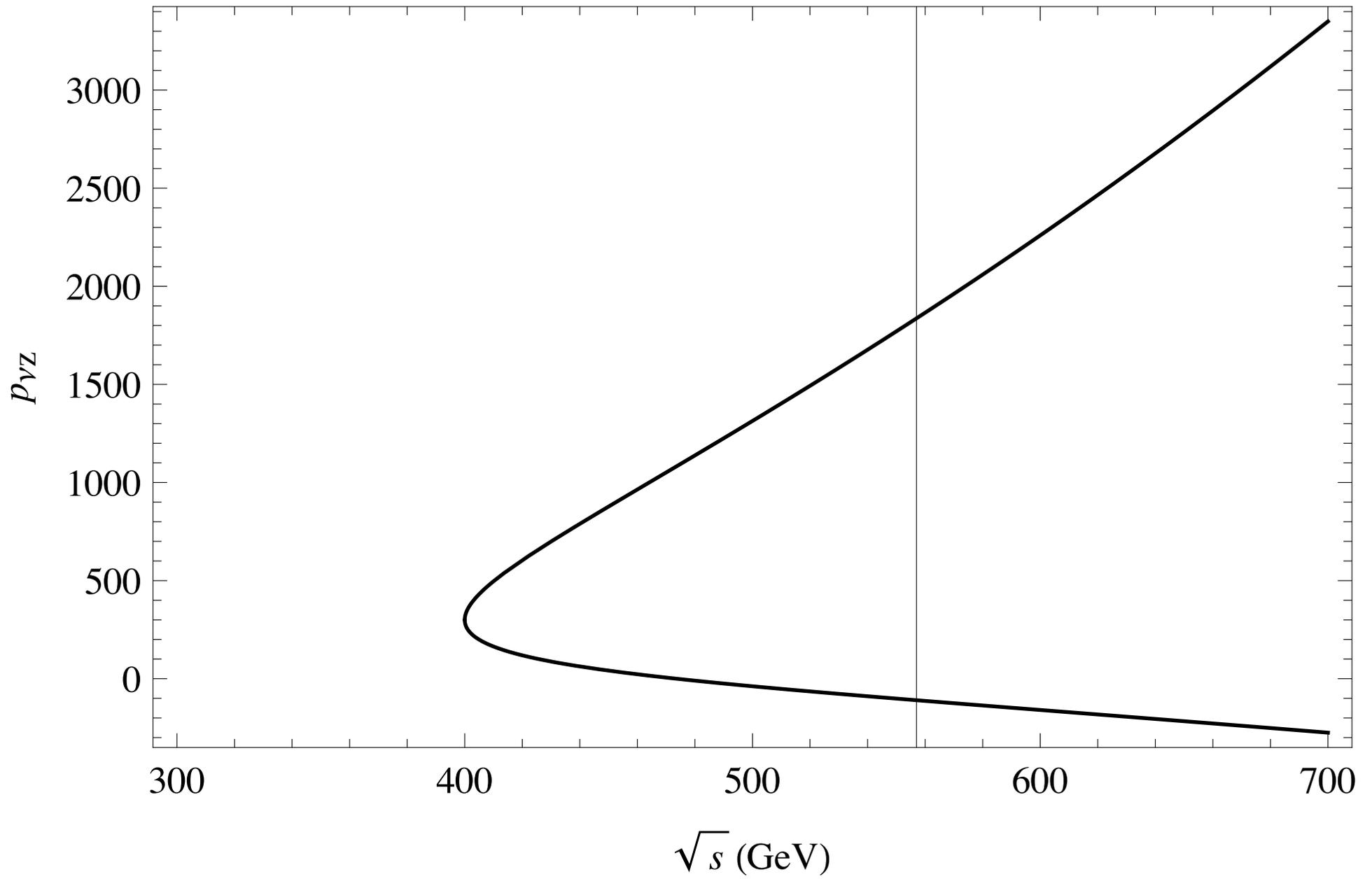
There is a well-known two-fold ambiguity in the z-component of the neutrino momentum.

$$s = (E_\ell + E_\nu)^2 - (p_{z\ell} + p_{z\nu})^2$$

$$s = (E_\ell + E_\nu)^2 - (p_{z\ell} + p_{z\nu})^2$$

$$p_{z\nu} = p_{z\ell} \left( \frac{\hat{s}}{2p_{T\ell}^2} - 1 \right) \pm \frac{E_\ell \sqrt{\hat{s}}}{p_{T\ell}} \sqrt{\frac{\hat{s}}{4p_{T\ell}^2} - 1}$$





Usually, we try to find a cut that improves our chances of guessing the right solution.

These cuts tend to reduce the signal rate and they tend to modify the angular distribution.

Can we do better?

Can we do better?

Yes!

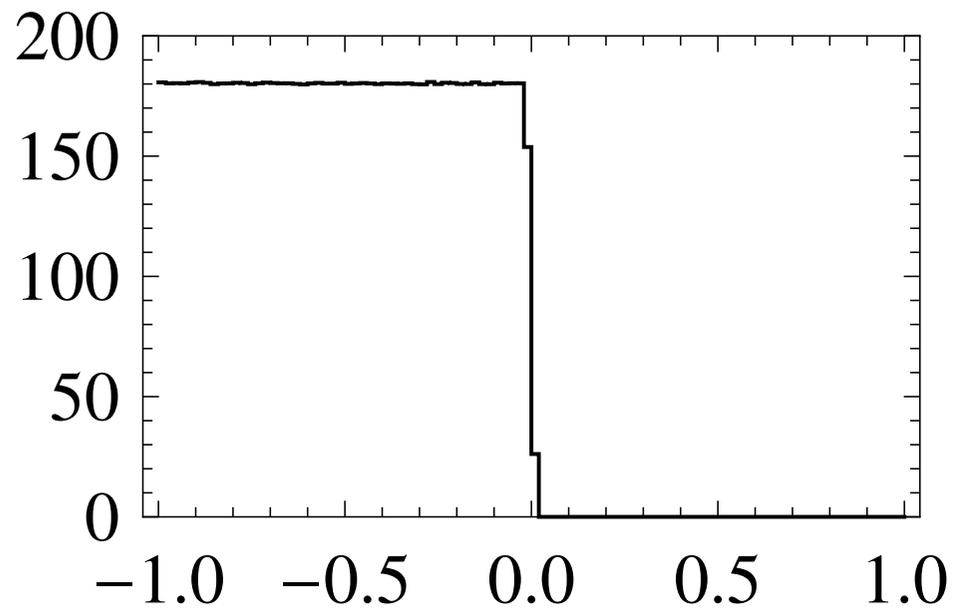
$$\cos \theta_S = -\sqrt{1 - \frac{4p_{Tl}^2}{M^2}} \operatorname{sign} \left( \frac{M}{2} - E_l \right)$$

$$\cos \theta_L = -\sqrt{1 - \frac{4p_{Tl}^2}{M^2}}$$

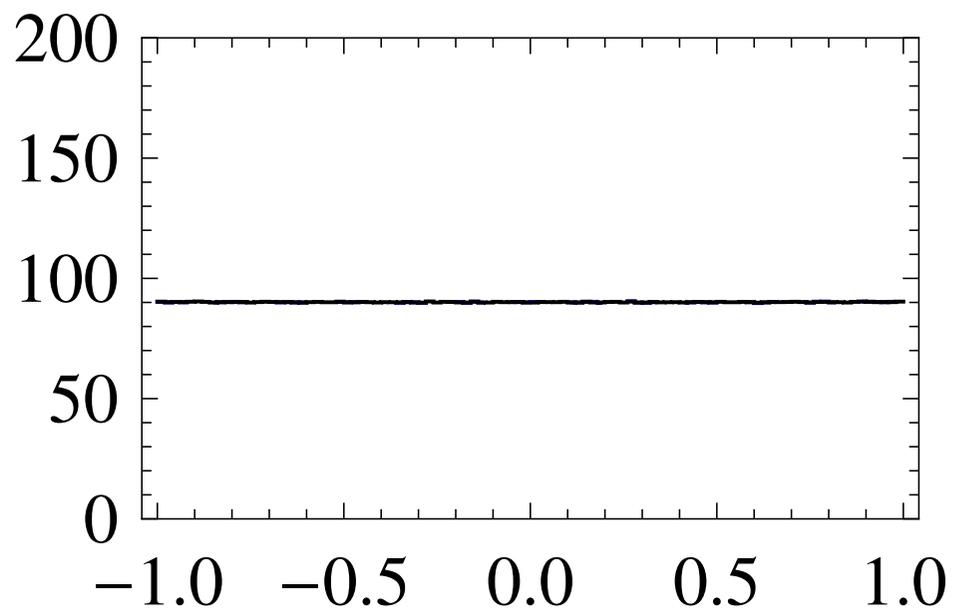
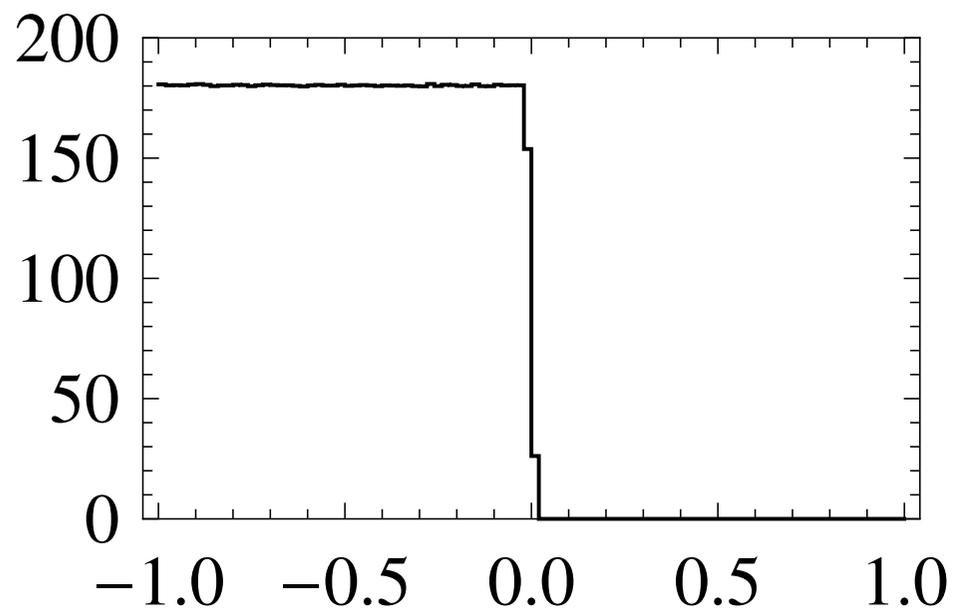
Let's look at the Large solution first:

We can reconstruct the spin.

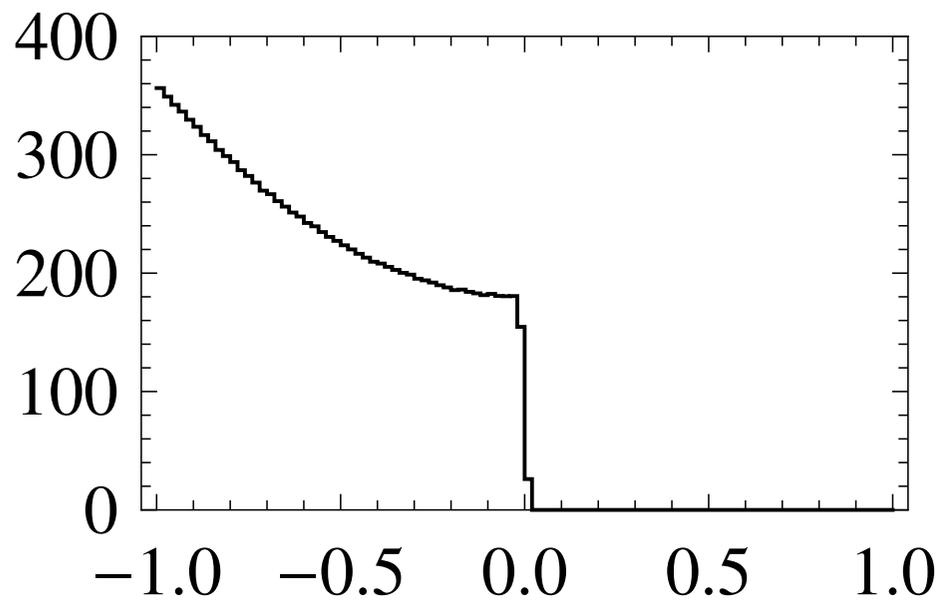
# Scalar



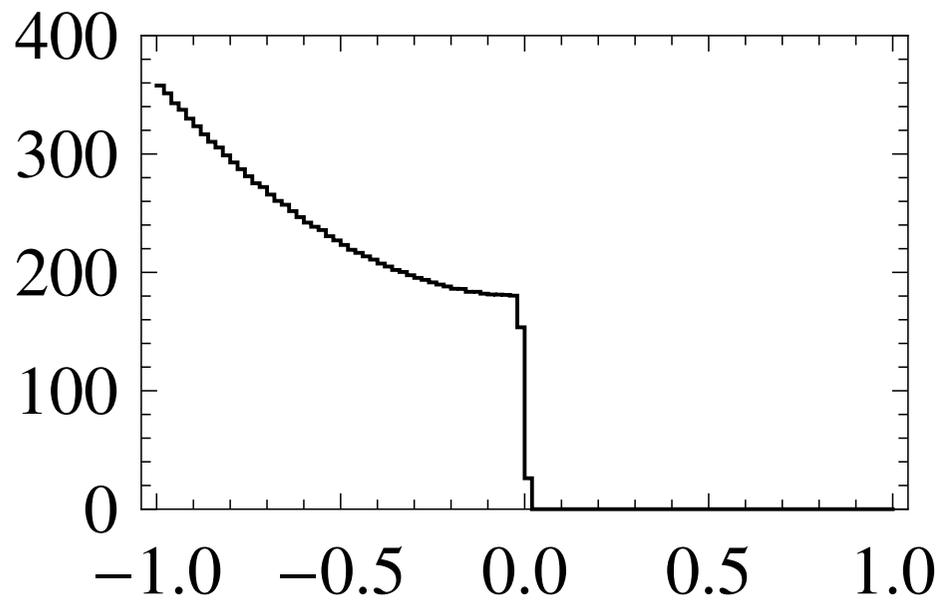
# Scalar



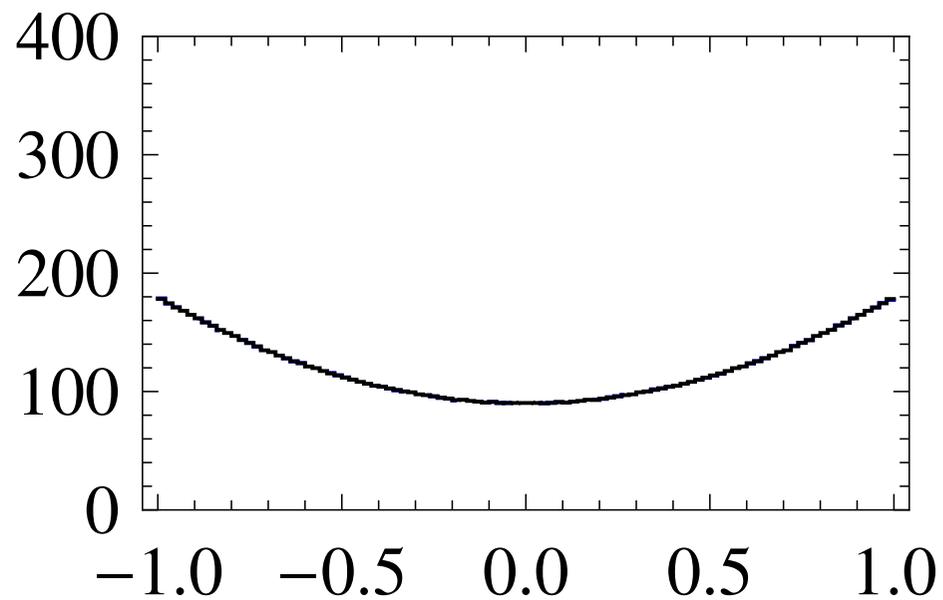
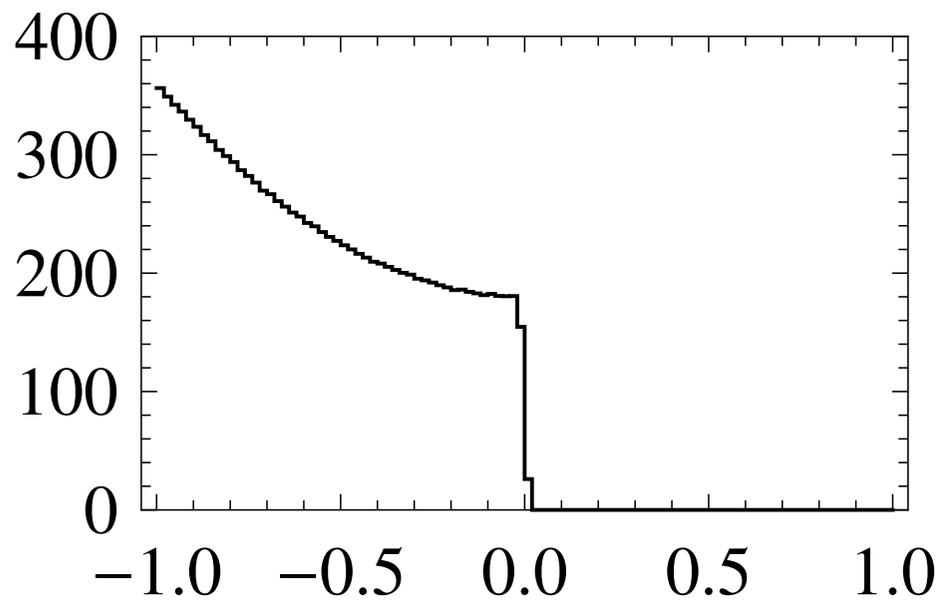
Symmetric Vector



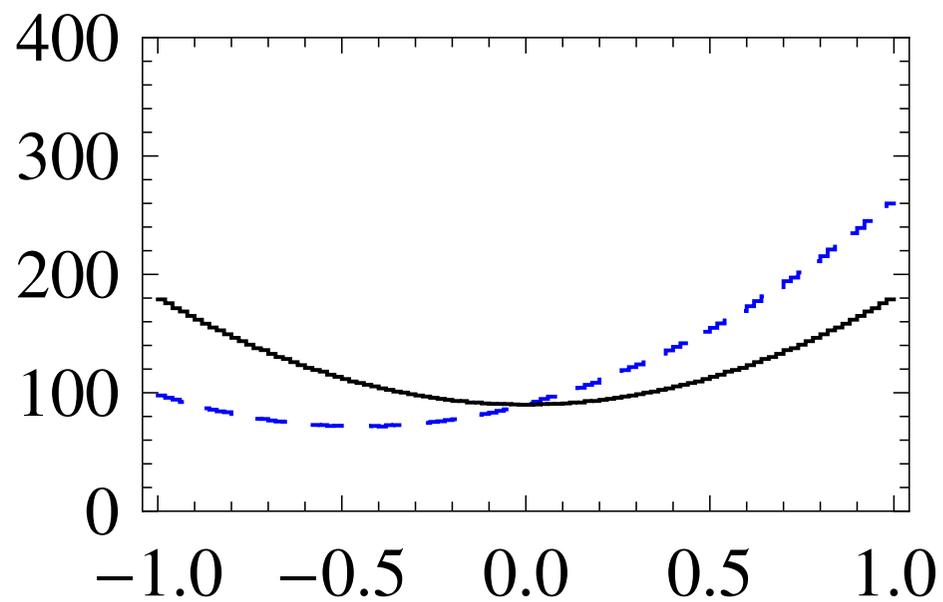
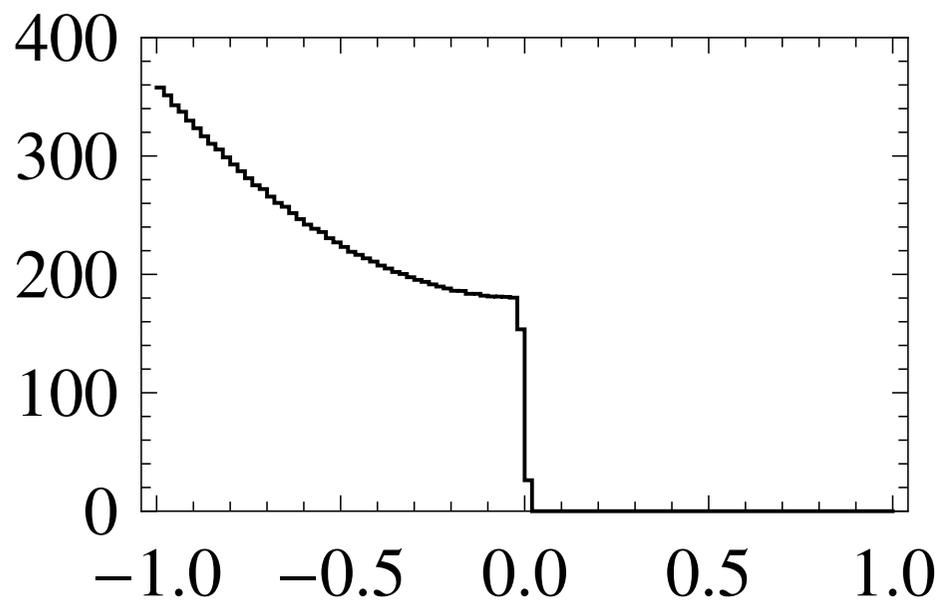
Left-Left Vector



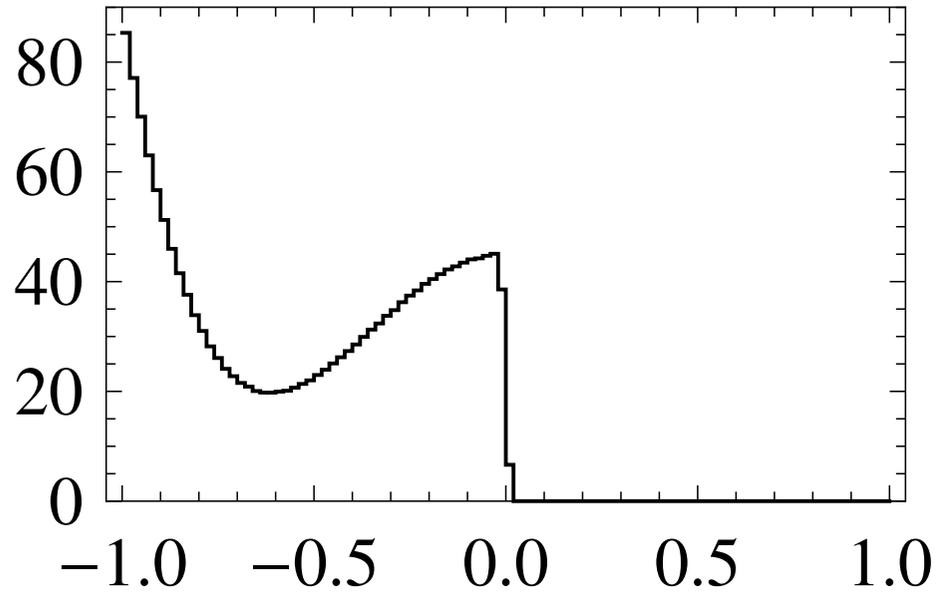
## Symmetric Vector



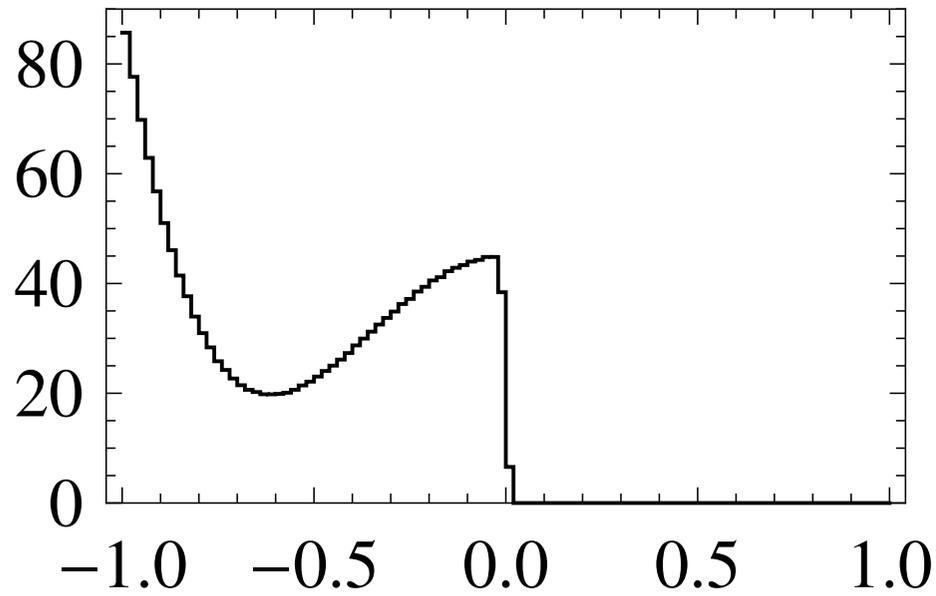
## Left-Left Vector



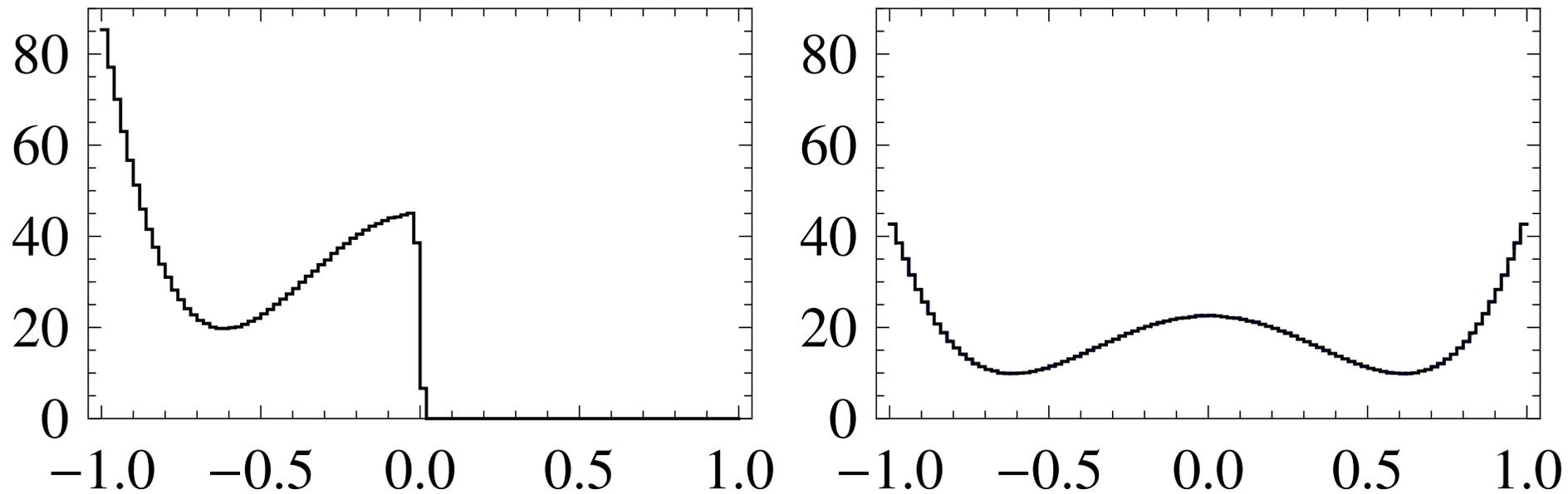
Symmetric Tensor



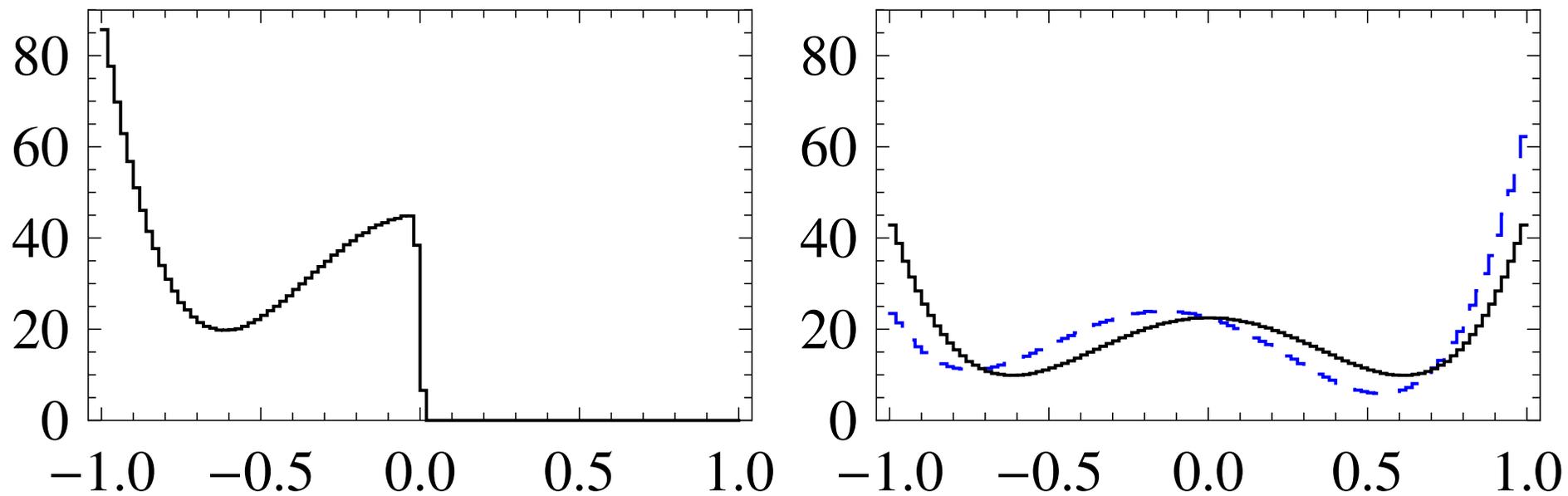
Left-Left Tensor



## Symmetric Tensor



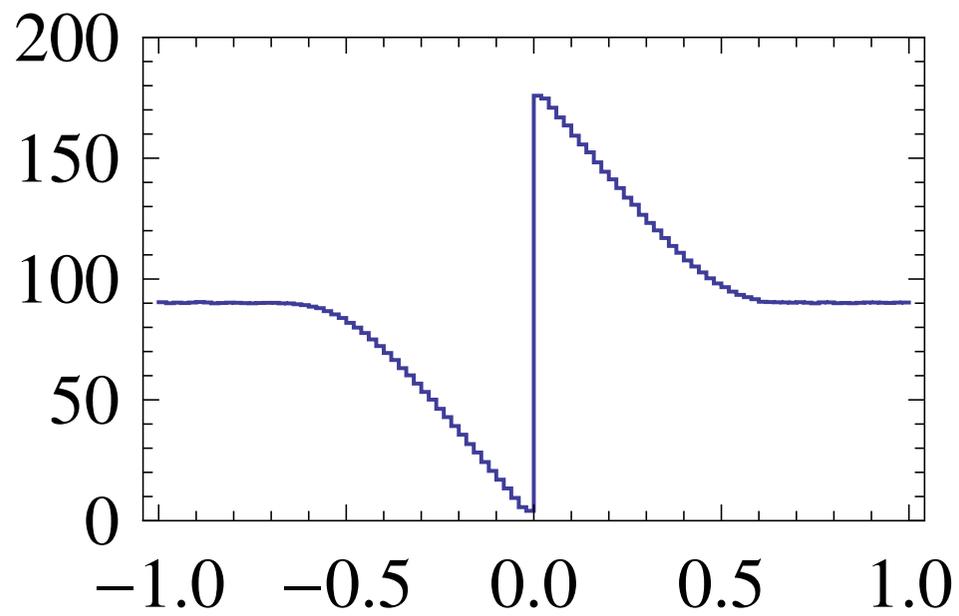
## Left-Left Tensor



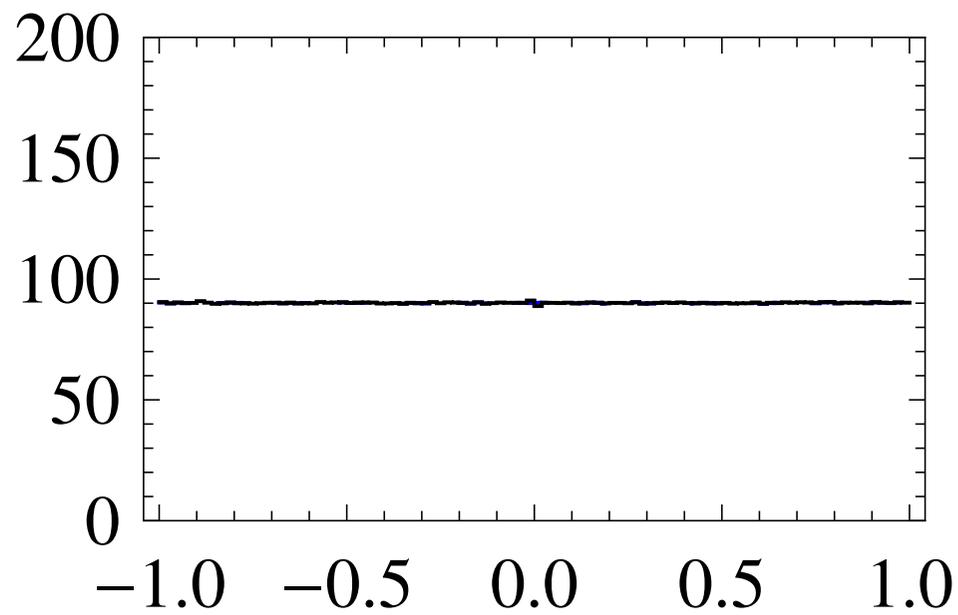
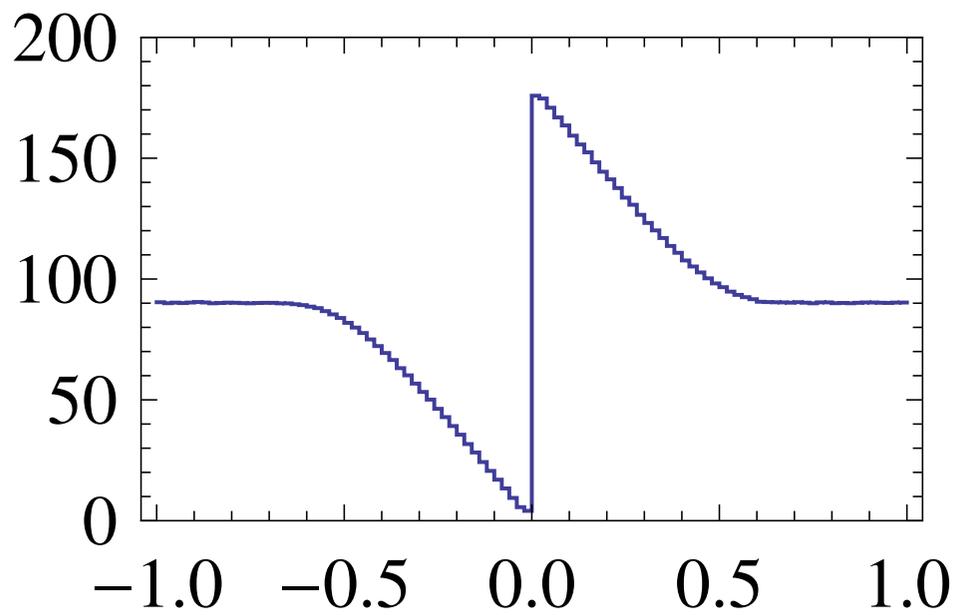
Let's look at the small solution next:

We can also reconstruct the parity violation.  
(Look in our paper for the details.)

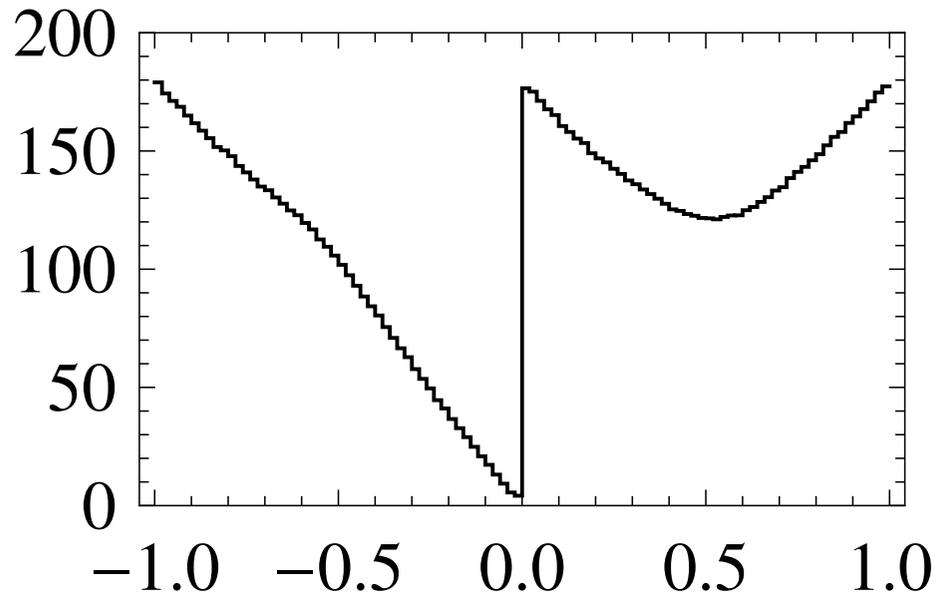
# Scalar



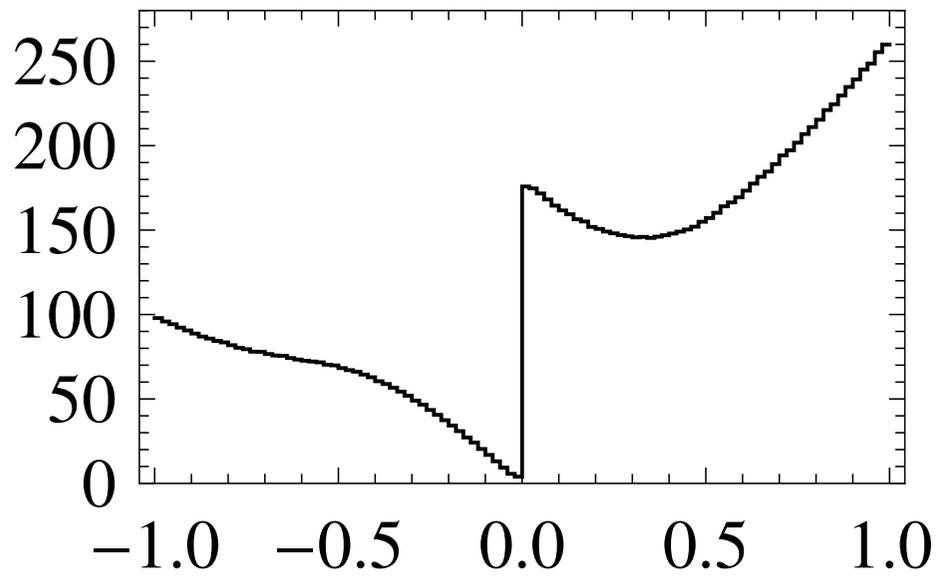
# Scalar



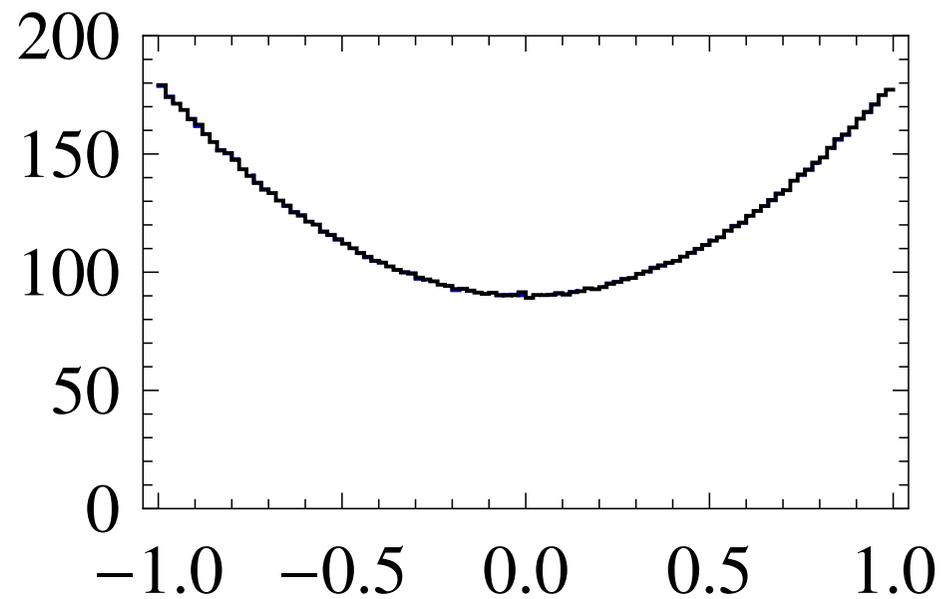
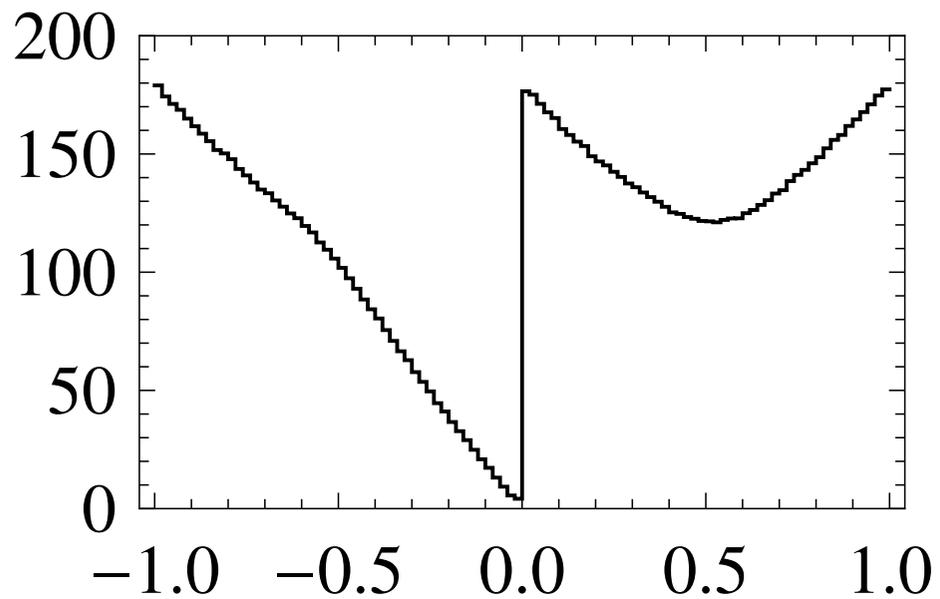
Symmetric Vector



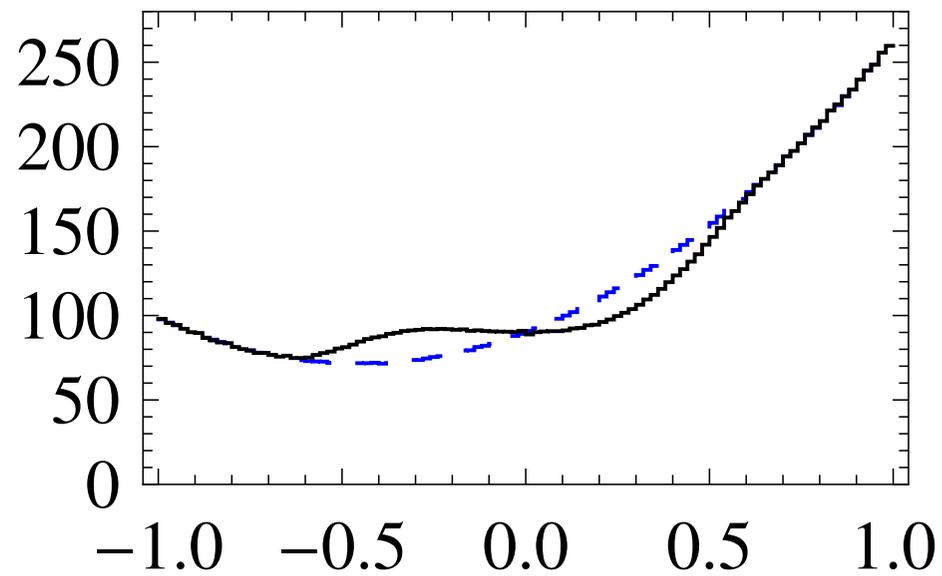
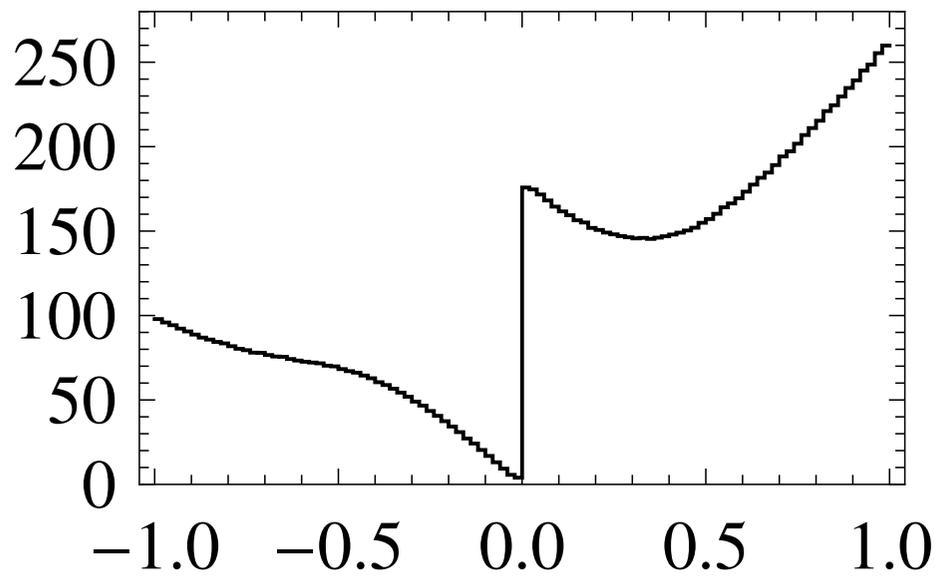
Left-Left Vector



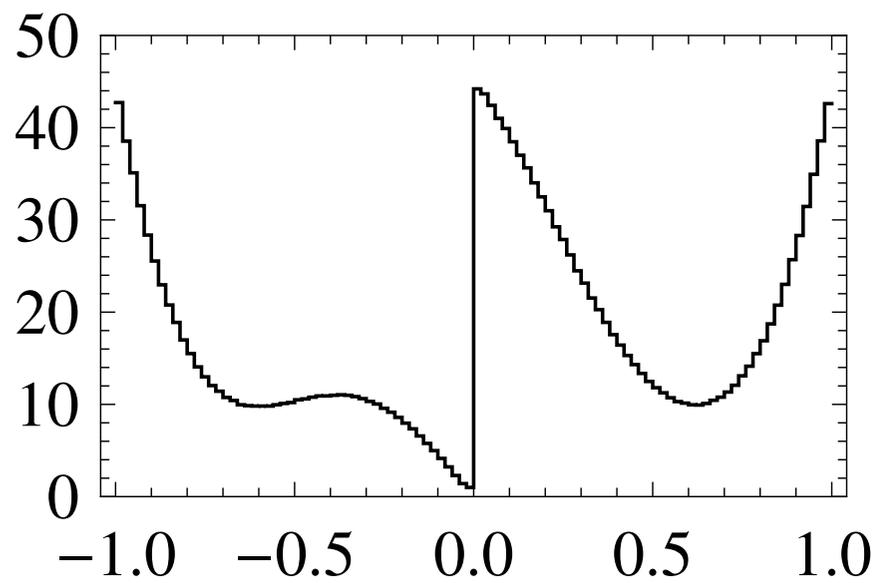
## Symmetric Vector



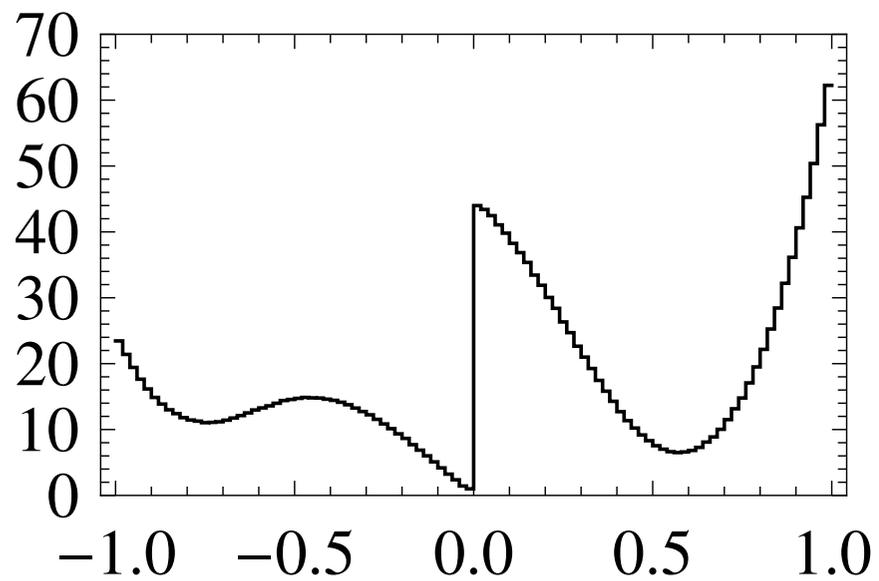
## Left-Left Vector



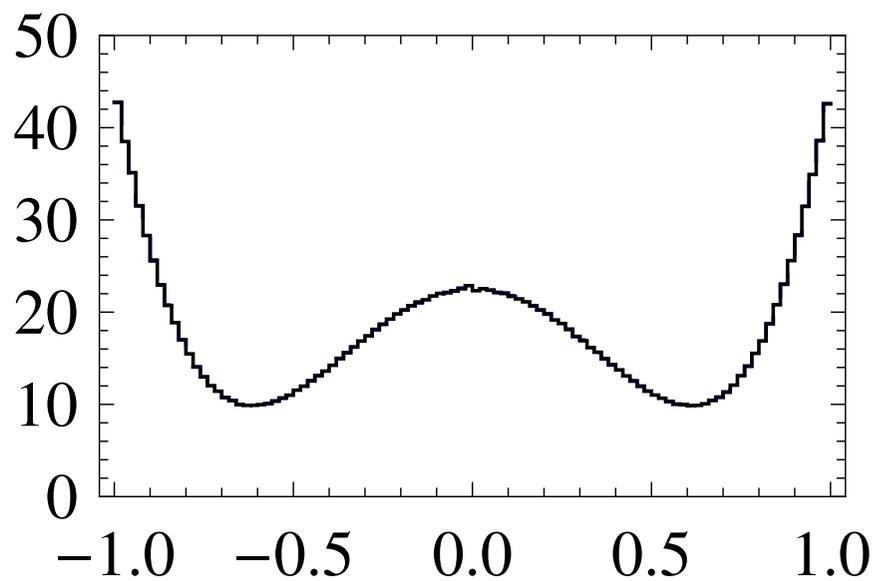
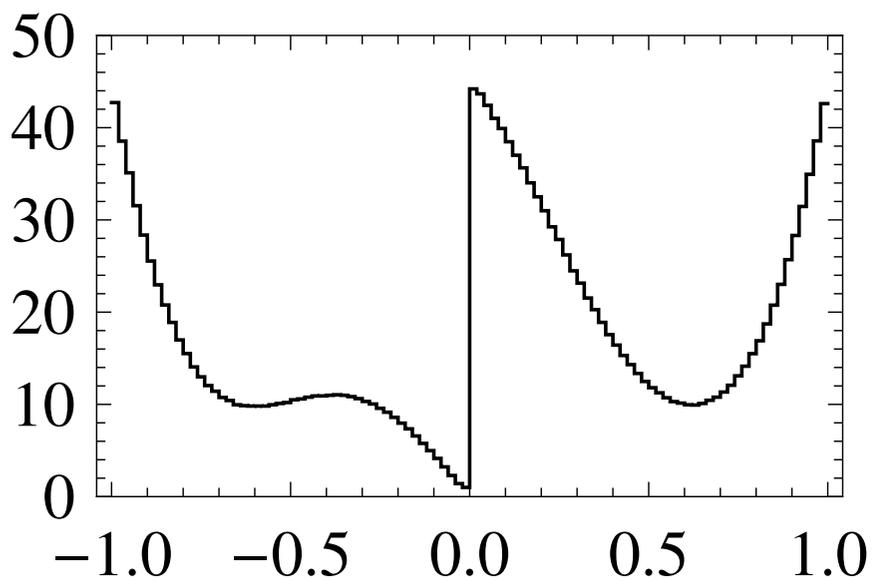
# Symmetric Tensor



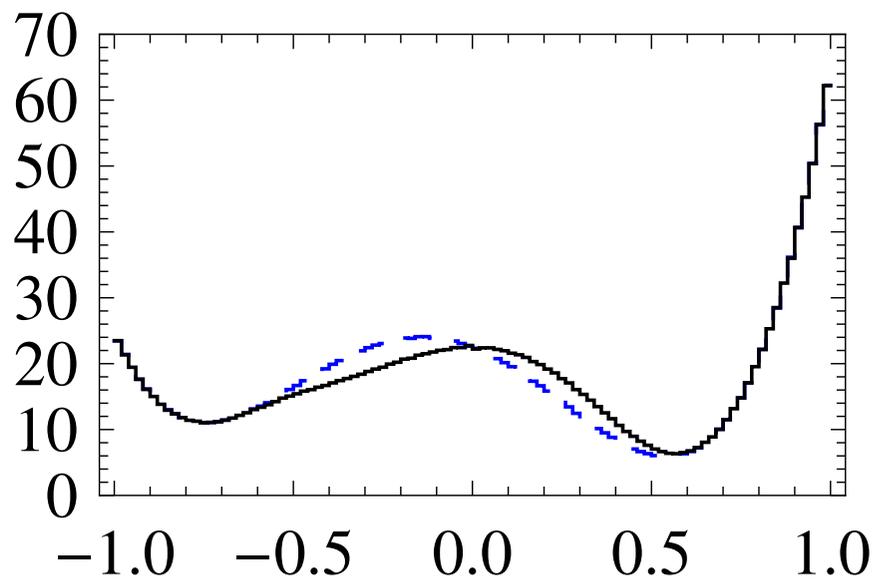
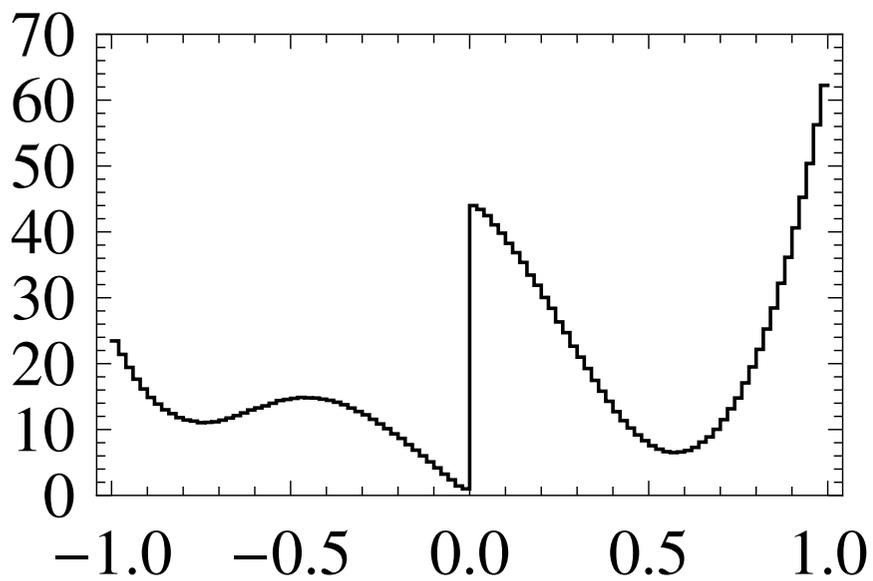
# Left-Left Tensor



## Symmetric Tensor

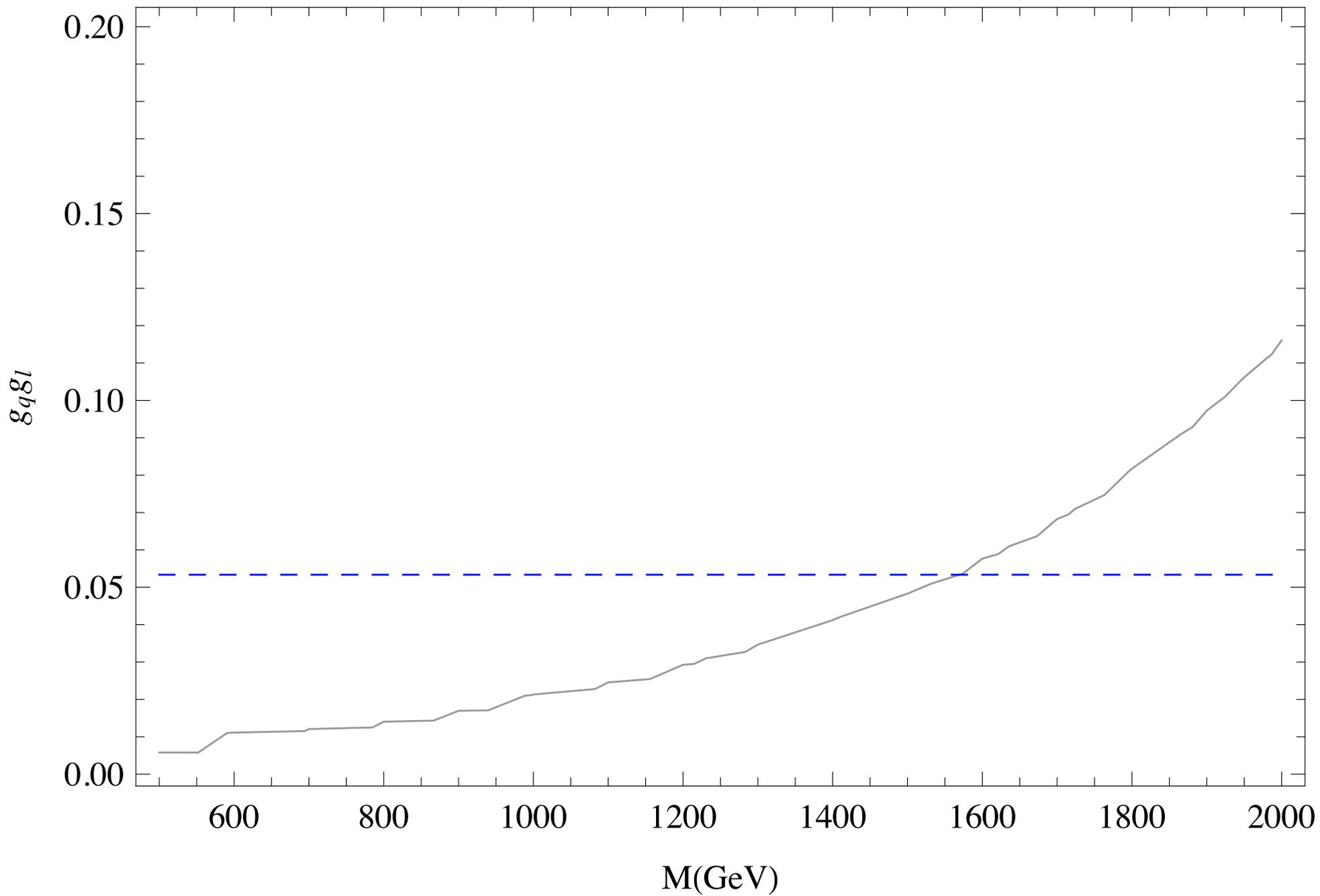


## Left-Left Tensor

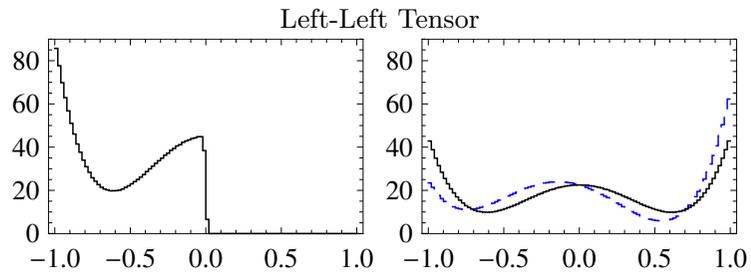
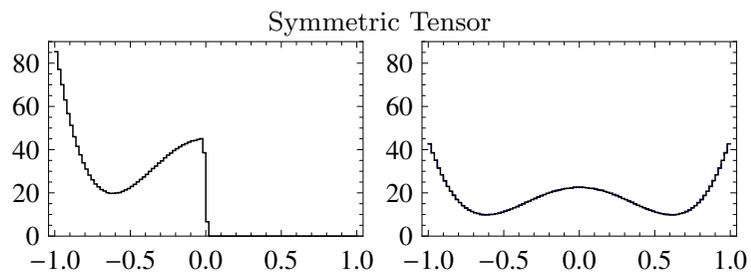
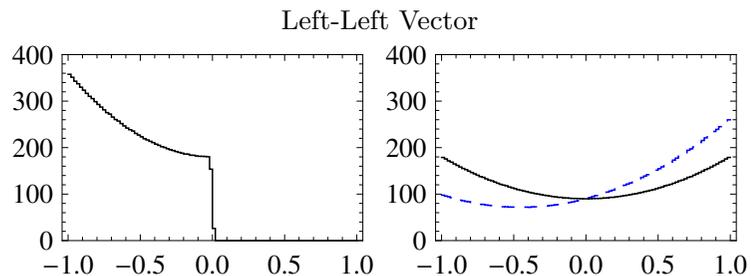
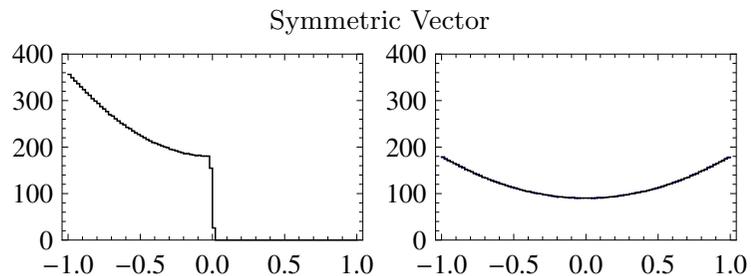
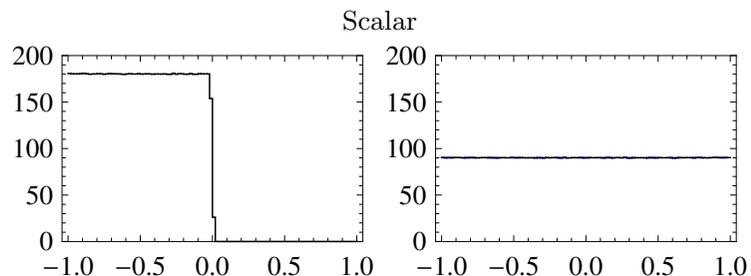


**To summarize:**

$pp \rightarrow e\nu$



# Large solution



# Small solution

