

# **“Development of a Spin-Light” Polarimeter for the EIC**

**Dipangkar Dutta**  
**Mississippi State University**

---

# The Collaboration

Dipangkar Dutta (PI), James Dunne, Edward Leggett, Mitra Shabestari

*Mississippi State University, Mississippi State, MS*

Wouter Deconinck (Co-PI), Valerie Gray

*College of William and Mary, Williamsburg, VA*

Abhay Deshpande (Co-PI)

*Stony Brook University, Stony Brook, NY*

Frank Maas (Co-PI)

*Johannes Gutenberg-Universität, Mainz, Germany*

Kent Paschke (co-PI)

*University of Virginia, Charlottesville, VA*

Paul Reimer (Co-PI)

*Argonne National Lab, Argonne, IL*

**The collaboration has extensive polarimetry and detector development experience at DESY, JLab, Mainz, and RHIC.**

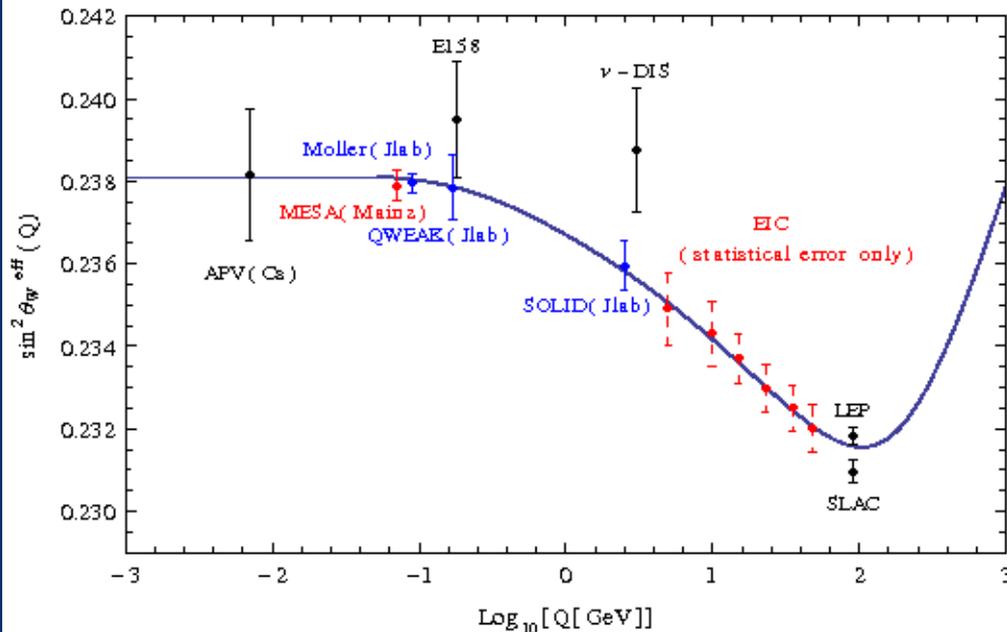
# Outline

- **Introduction**
- **Synchrotron Radiation**
  - Classical & Quantum description
- **“Spin-Light”**
- **Conceptual Design of a “Spin-Light” polarimeter**
- **Detector and Prototype Development**
- **Summary**

# Introduction

The high energy and high luminosity polarized electrons, protons and ions at the EIC promises:

- a precise 3-D mapping of the proton's internal structure
- fundamental tests of QCD (such as Bjorken sum rule)
- tests of the SM at the quantum loop level that probe “new physics”



This entire program at the EIC requires precision electron polarimetry.

We propose to develop a novel continuous non-invasive polarimeter based on the spin dependence of synchrotron radiation (SR)

Projected uncertainties of future EIC measurements of  $\sin^2 \theta_W$ .

# Synchrotron Radiation- “Electronic Light”

After its discovery, the angular and spectral distribution were worked out in classical E&M

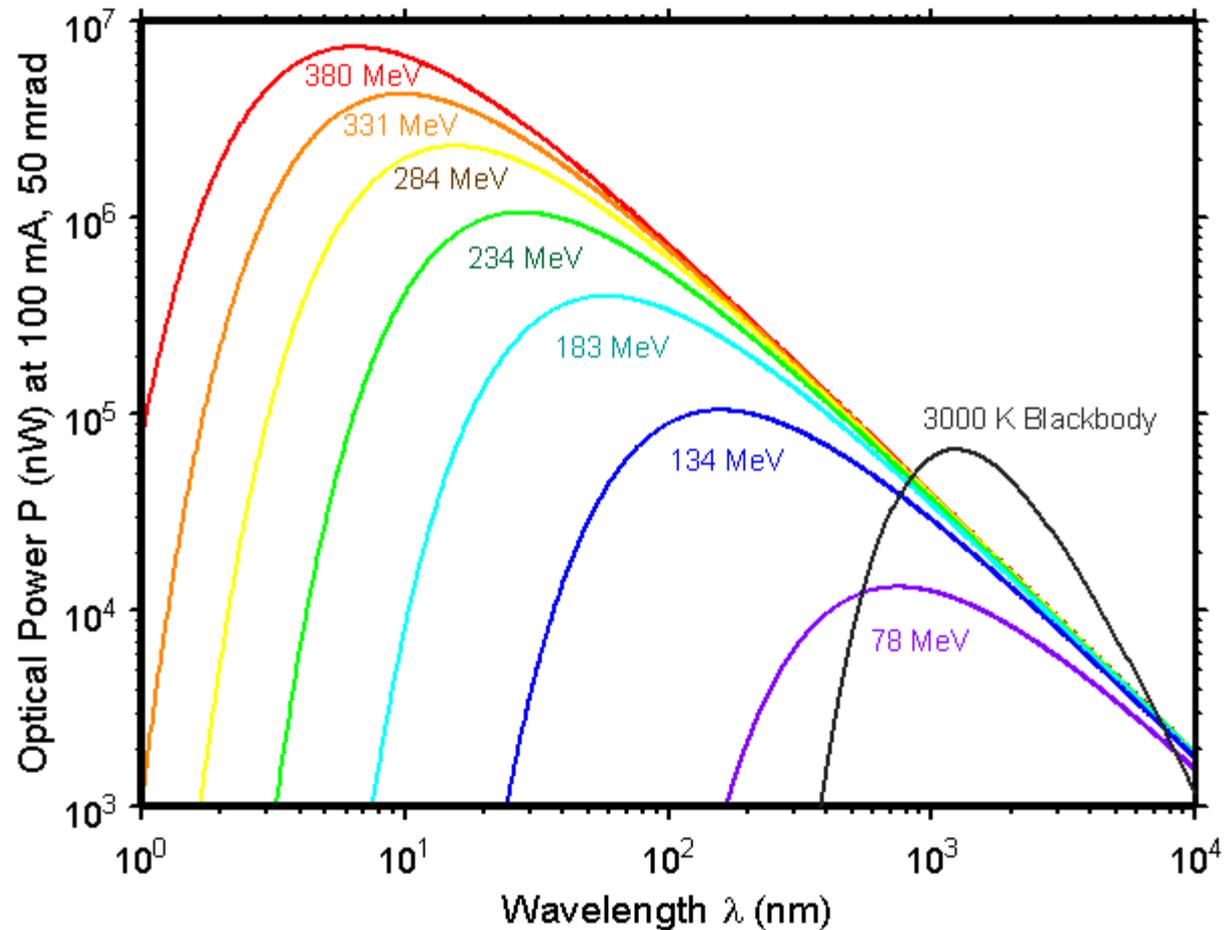
## Total Rate

$$P = \frac{dU}{dt} = \frac{2}{3} \frac{e^2 \gamma^4 c}{R^2}$$

$$\gamma = \frac{E}{m_0 c^2} \quad : \text{ Lorentz boost}$$

## Strong linear polarization

$$P_\sigma = \frac{7}{8} P ; \quad P_\pi = \frac{1}{8} P$$



# Synchrotron Radiation- “Electronic Light”

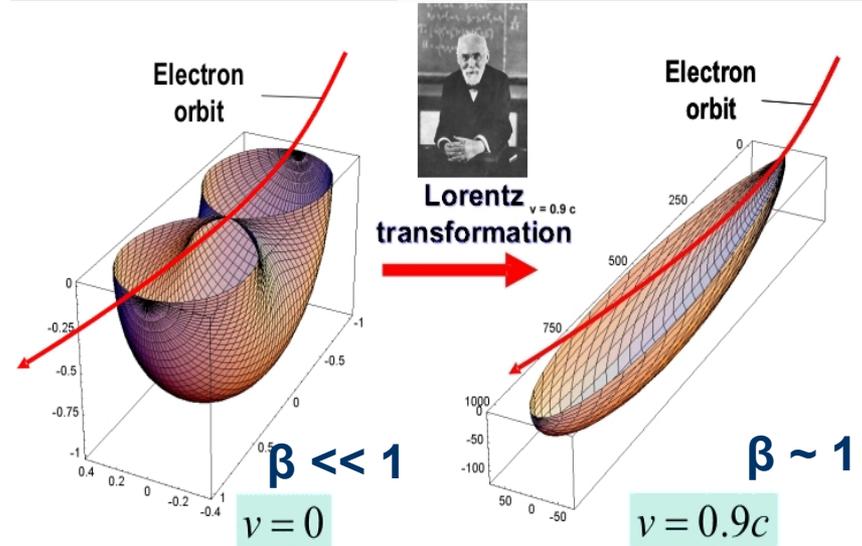
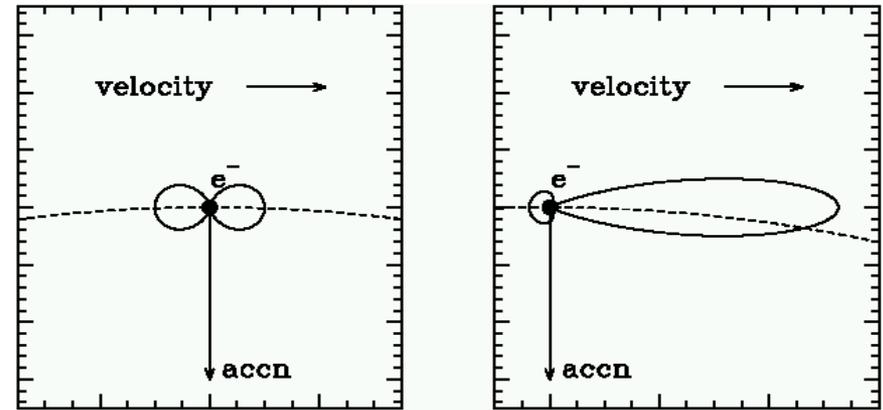
## Angular distribution

$$\frac{dU}{dt d\Omega} = \frac{e^2 \gamma^4 c}{4\pi R^2} \frac{(1 - \beta \cos\theta)^2 - (1 - \beta^2) \sin^2\theta \cos^2\phi}{(1 - \beta \cos\theta)^5}$$

For  $\gamma \gg 1$   $\theta \sim 1/\gamma$

SR emitted in a very small cone

For  $E_e = 11$  GeV,  
vert. size = 90  $\mu$ rad  
i.e. 10m from the  
source  $\sim 1$  mm height.



# Synchrotron Radiation- “Electronic Light”

## Quantum Corrections

Exact QED calculations by A.A. Sokolov and I. M.Ternov (1960s)

# Synchrotron Radiation- “Electronic Light”

## Quantum Corrections

Exact QED calculations by A.A. Sokolov and I. M. Ternov (1960s)

Classical theory (continuous SR spectrum) valid for

$$E \ll E_{\text{crit}} \sim 10^6 \text{ GeV} \text{ and } B \ll B_{\text{crit}} \sim 4 \times 10^9 \text{ T}$$

$E_{\text{crit}}$  : single SR photon carries away all of the electron's energy

# Synchrotron Radiation- “Electronic Light”

## Quantum Corrections

Exact QED calculations by A.A. Sokolov and I. M. Ternov (1960s)

Classical theory (continuous SR spectrum) valid for

$$E \ll E_{\text{crit}} \sim 10^6 \text{ GeV} \text{ and } B \ll B_{\text{crit}} \sim 4 \times 10^9 \text{ T}$$

$E_{\text{crit}}$ : single SR photon carries away all of the electron's energy

But even at lower energies

QED corrections → spin dependence of the radiated power

$$P = P^{\text{clas}} \frac{9\sqrt{3}}{16\pi} \sum_{s'} \int \frac{y dy}{(1+\xi y)^4} I_{ss'}^2 F(y); \quad y = \frac{\omega}{\omega_c}; \quad \xi = \frac{3}{2} \frac{B}{B_{\text{crit}}} \gamma; \quad \begin{array}{l} s = \text{radial quantum \#} \\ l = \text{Laguerre func.} \end{array}$$

For  $\xi \ll 1$  and electron spin  $j, j' = \pm 1$

$$P = P^{\text{clas}} \left[ \left(1 - \frac{55\sqrt{3}}{24} \xi + \frac{64}{3} \xi^2\right) - \left(\frac{1+jj'}{2}\right) \left(j\xi + \frac{5}{9} \xi^2 + \frac{245\sqrt{3}}{48} j \xi^2\right) + \left(\frac{1-jj'}{2}\right) \left(\frac{4}{3} \xi^2 + \frac{315\sqrt{3}}{432} j \xi^2\right) + \dots \right]$$

# Synchrotron Radiation- “Electronic Light”

## Quantum Corrections

Exact QED calculations by A.A. Sokolov and I. M. Ternov (1960s)

Classical theory (continuous SR spectrum) valid for

$$E \ll E_{\text{crit}} \sim 10^6 \text{ GeV} \text{ and } B \ll B_{\text{crit}} \sim 4 \times 10^9 \text{ T}$$

$E_{\text{crit}}$ : single SR photon carries away all of the electron's energy

But even at lower energies

QED corrections → spin dependence of the radiated power

$$P = P^{\text{clas}} \frac{9\sqrt{3}}{16\pi} \sum_{s'} \int \frac{y dy}{(1+\xi y)^4} I_{ss'}^2 F(y); \quad y = \frac{\omega}{\omega_c}; \quad \xi = \frac{3}{2} \frac{B}{B_{\text{crit}}} \gamma; \quad \begin{array}{l} s = \text{radial quantum \#} \\ l = \text{Laguerre func.} \end{array}$$

For  $\xi \ll 1$  and electron spin  $j, j' = \pm 1$

$$P = P^{\text{clas}} \left[ \left( 1 - \frac{55\sqrt{3}}{24} \xi + \frac{64}{3} \xi^2 \right) \left( \frac{1+jj'}{2} \right) \left( j\xi + \frac{5}{9} \xi^2 + \frac{245\sqrt{3}}{48} j\xi^2 \right) + \left( \frac{1-jj'}{2} \right) \left( \frac{4}{3} \xi^2 + \frac{315\sqrt{3}}{432} j\xi^2 \right) + \dots \right]$$

spin dependent term

spin-flip dependent term

# The Spin Dependence

To the first order in  $\xi$  the difference in SR intensity between polarized and unpolarized electrons is  $\delta = \xi j \sim 10^{-4}$  for 100  $\mu\text{A}$ , 5.0 GeV electrons

# The Spin Dependence

To the first order in  $\xi$  the difference in SR intensity between polarized and unpolarized electrons is  $\delta = \xi j \sim 10^{-4}$  for 100  $\mu\text{A}$ , 5.0 GeV electrons

Verified experimentally at the VEPP-4 storage ring in Novosibirsk  
Belomestnykh et al., NIM 227, 173 (1984)

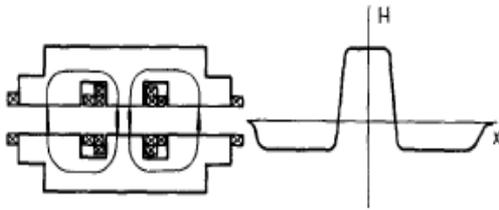


Fig. 1. The field vs the current in the 'snake'. A schematic of the 'snake' and the field distribution along its axis are shown below.

## 3 pole magnetic snake/wiggler

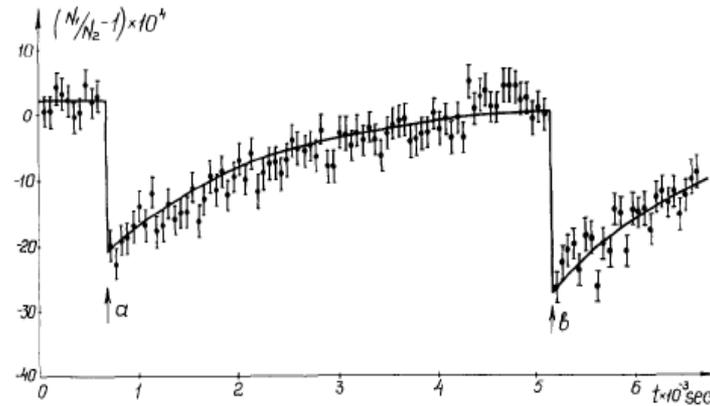


Fig. 12. The measurement results of the SR-intensity as a function of the degree of polarization of the beam. The field in the 'snake' coincides, in direction, with the storage ring guiding field. At points a and b one of the bunches ( $N_1$ ), was quickly depolarized. The measurement time at a point is 60 s. The bunch polarization time is  $\tau_p = 1740 \pm 20$  s ( $\xi = 0.726$ ).

# The Spin Dependence

To the first order in  $\xi$  the difference in SR intensity between polarized and unpolarized electrons is  $\delta = \xi j \sim 10^{-4}$  for 100  $\mu\text{A}$ , 5.0 GeV electrons

Verified experimentally at the VEPP-4 storage ring in Novosibirsk  
Belomestnykh et al., NIM 227, 173 (1984)

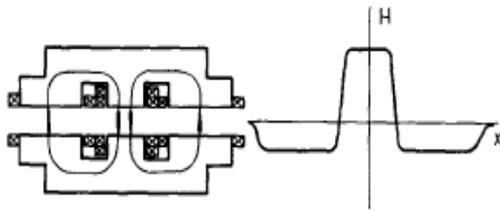


Fig. 1. The field vs the current in the 'snake'. A schematic of the 'snake' and the field distribution along its axis are shown below.

3 pole magnetic snake/wiggler

An RF field used to depolarize the electrons

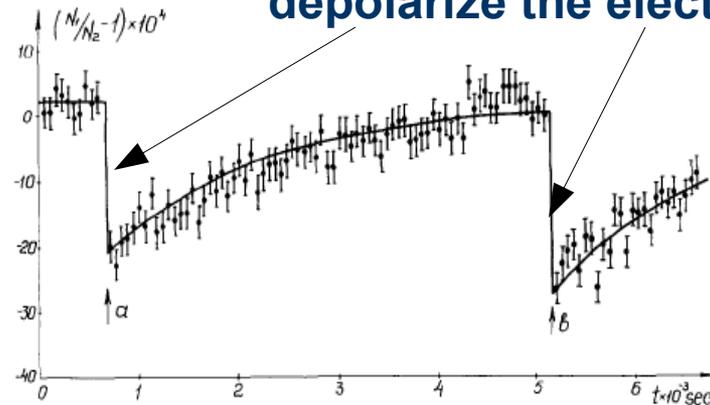


Fig. 12. The measurement results of the SR-intensity as a function of the degree of polarization of the beam. The field in the 'snake' coincides, in direction, with the storage ring guiding field. At points a and b one of the bunches ( $N_1$ ), was quickly depolarized. The measurement time at a point is 60 s. The bunch polarization time is  $\tau_p = 1740 \pm 20$  s ( $\xi = 0.726$ ).

# The Spin Dependence

To the first order in  $\xi$  the difference in SR intensity between polarized and unpolarized electrons is  $\delta = \xi j \sim 10^{-4}$  for 100  $\mu\text{A}$ , 5.0 GeV electrons

Verified experimentally at the VEPP-4 storage ring in Novosibirsk  
Belomestnykh et al., NIM 227, 173 (1984)

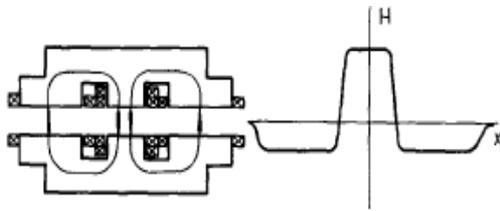


Fig. 1. The field vs the current in the 'snake'. A schematic of the 'snake' and the field distribution along its axis are shown below.

3 pole magnetic snake/wiggler

An RF field used to depolarize the electrons

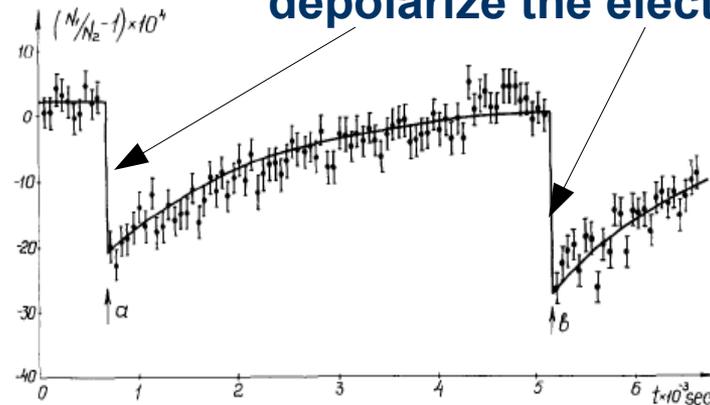


Fig. 12. The measurement results of the SR-intensity as a function of the degree of polarization of the beam. The field in the 'snake' coincides, in direction, with the storage ring guiding field. At points a and b one of the bunches ( $N_1$ ), was quickly depolarized. The measurement time at a point is 60 s. The bunch polarization time is  $\tau_p = 1740 \pm 20$  s ( $\xi = 0.726$ ).

The spin-flip term contributes only as  $\sim \xi^2$

This is responsible for the transverse self polarization of electron beams in storage rings: called the Sokolov-Ternov effect

Used to produce polarized electrons at various accelerator such as DESY

# Longitudinal “Spin Light”

For longitudinally polarized electrons

Power from  $n_e$  electrons ( ignoring spin flip and all terms  $O(\xi^2)$  )

$$P_y(\text{long}) = \frac{9 n_e c e^2}{16 \pi^3 R^2} \gamma^5 \int_0^\infty \frac{y^2 dy}{(1 + \xi y)^4} \oint d\Omega (1 + \alpha^2)^2 \left[ K_{2/3}^2(z) + \frac{\alpha^2}{1 + \alpha^2} K_{1/3}^2(z) + j \xi y \frac{\alpha}{\sqrt{1 + \alpha^2}} K_{1/3}(z) K_{2/3}(z) \right]$$

$R$  = bending radius,  $y = \frac{\omega}{\omega_c}$ ;  $\xi = \frac{3}{2} \frac{B}{B_{crit}} \gamma$ ;  $\alpha = \gamma \psi$ ;  $z = \frac{\omega}{2 \omega_c} (1 + \alpha^2)^{3/2}$ ;  $K_{1/3}, K_{2/3}$  modified Bessel function

↑  
vertical angle

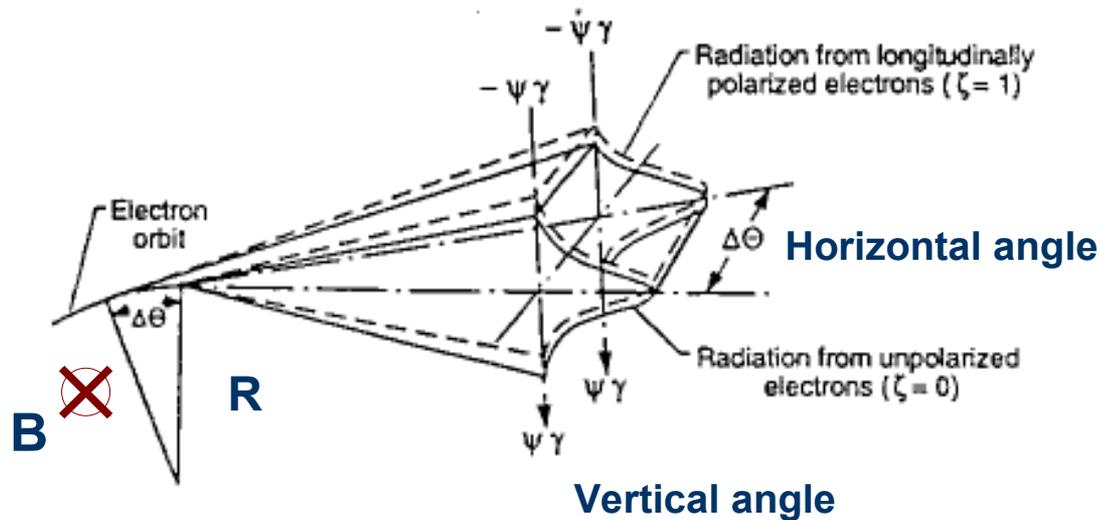


Figure 1: Geometrical definitions.

# Longitudinal “Spin Light”

For longitudinally polarized electrons

Power from  $n_e$  electrons ( ignoring spin flip and all terms  $O(\xi^2)$  )

$$P_y(\text{long}) = \frac{9 n_e c e^2}{16 \pi^3 R^2} \gamma^5 \int_0^\infty \frac{y^2 dy}{(1 + \xi y)^4} \oint d\Omega (1 + \alpha^2)^2 \left[ K_{2/3}^2(z) + \frac{\alpha^2}{1 + \alpha^2} K_{1/3}^2(z) + j \xi y \frac{\alpha}{\sqrt{1 + \alpha^2}} K_{1/3}(z) K_{2/3}(z) \right]$$

$R$  = bending radius,  $y = \frac{\omega}{\omega_c}$ ;  $\xi = \frac{3}{2} \frac{B}{B_{crit}} \gamma$ ;  $\alpha = \gamma \psi$ ;  $z = \frac{\omega}{2 \omega_c} (1 + \alpha^2)^{3/2}$ ;  $K_{1/3}, K_{2/3}$  modified Bessel function

$\uparrow$   
 vertical angle

An odd function of the vertical angle

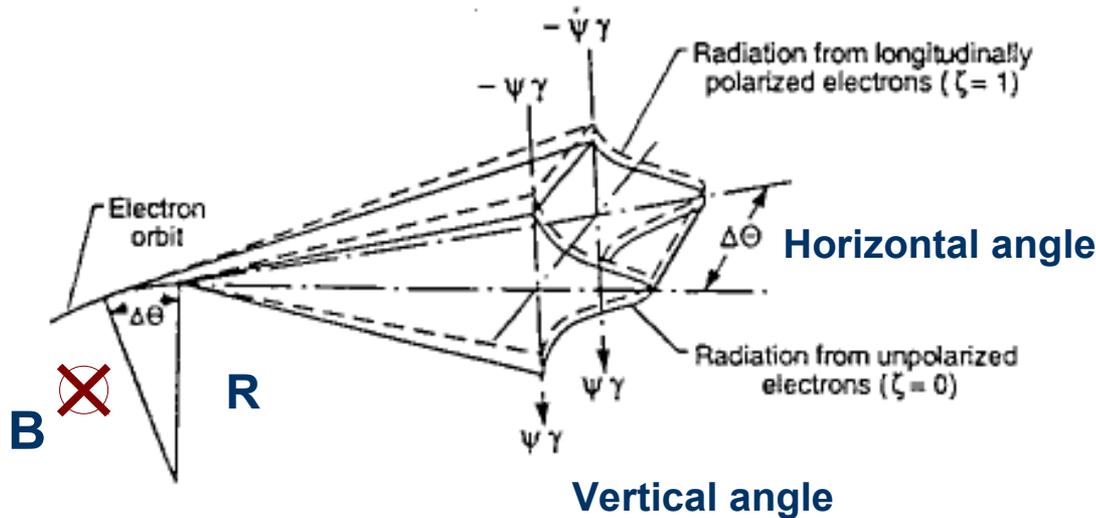


Figure 1: Geometrical definitions.

# Longitudinal “Spin Light”

For longitudinally polarized electrons

Power from  $n_e$  electrons ( ignoring spin flip and all terms  $O(\xi^2)$  )

$$P_y(\text{long}) = \frac{9 n_e c e^2}{16 \pi^3 R^2} \gamma^5 \int_0^\infty \frac{y^2 dy}{(1 + \xi y)^4} \oint d\Omega (1 + \alpha^2)^2 \left[ K_{2/3}^2(z) + \frac{\alpha^2}{1 + \alpha^2} K_{1/3}^2(z) + j \xi y \frac{\alpha}{\sqrt{1 + \alpha^2}} K_{1/3}(z) K_{2/3}(z) \right]$$

$R$  = bending radius,  $y = \frac{\omega}{\omega_c}$ ;  $\xi = \frac{3}{2} \frac{B}{B_{crit}} \gamma$ ;  $\alpha = \gamma \psi$ ;  $z = \frac{\omega}{2 \omega_c} (1 + \alpha^2)^{3/2}$ ;  $K_{1/3}, K_{2/3}$  modified Bessel function

vertical angle

An odd function of the vertical angle

Integrated over all vertical angles the total SR power is spin **independent**

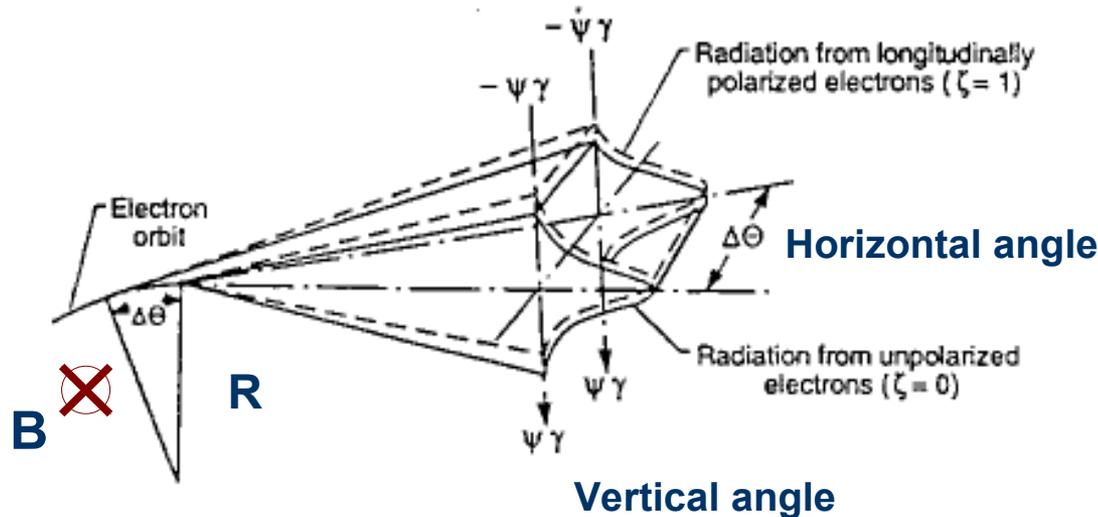
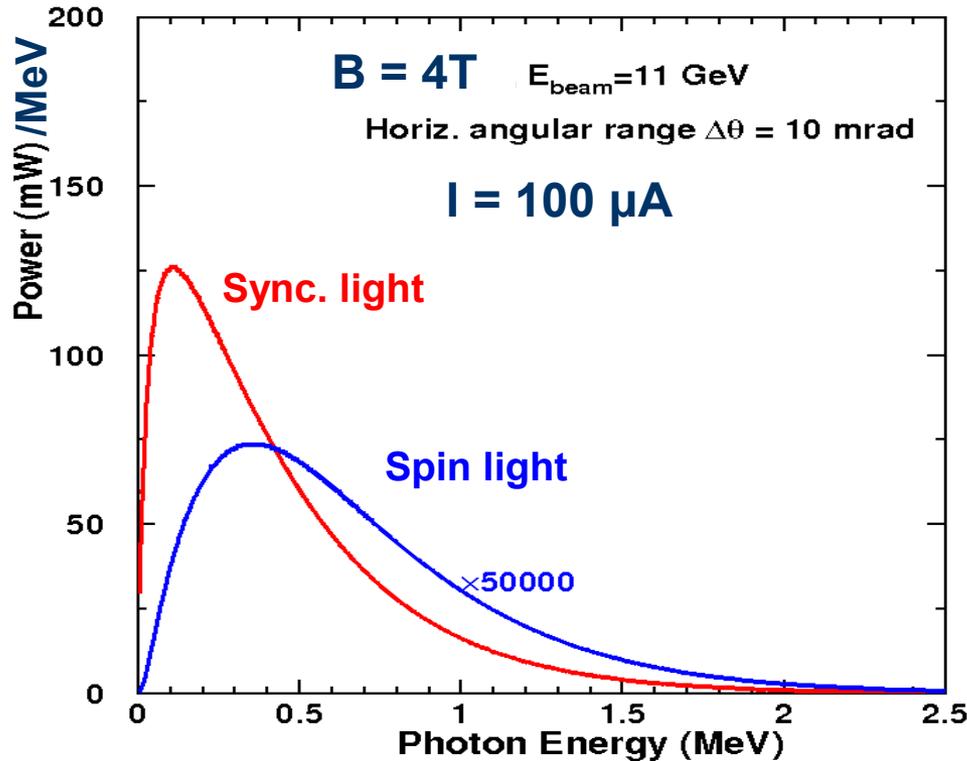


Figure 1: Geometrical definitions.

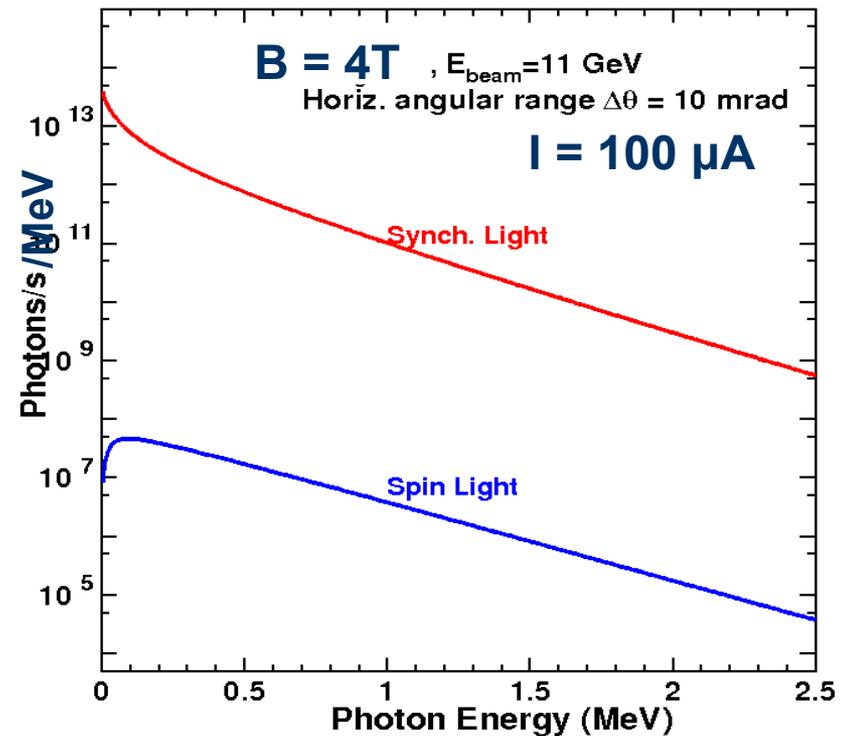
# of photons radiated above and below the orbital plane are not equal

# “Spin Light” - Some Characteristics

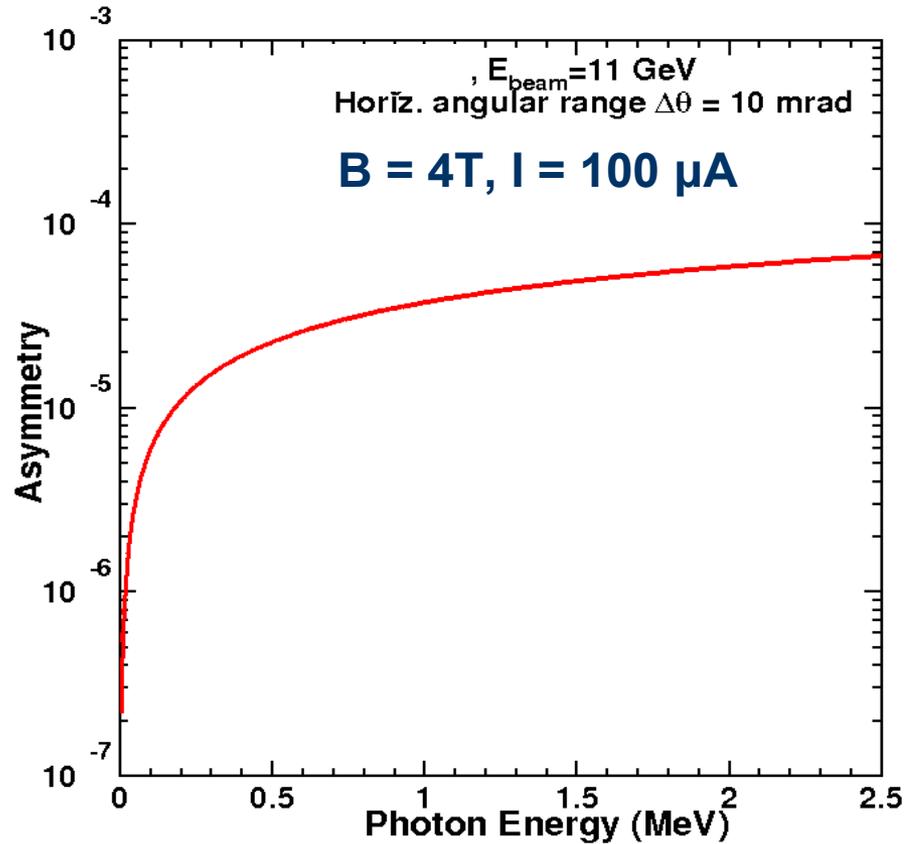


Radiated power peak at different frequencies for Sync vs Spin light

Number of photons/s  
 $\sim 10^5$  different

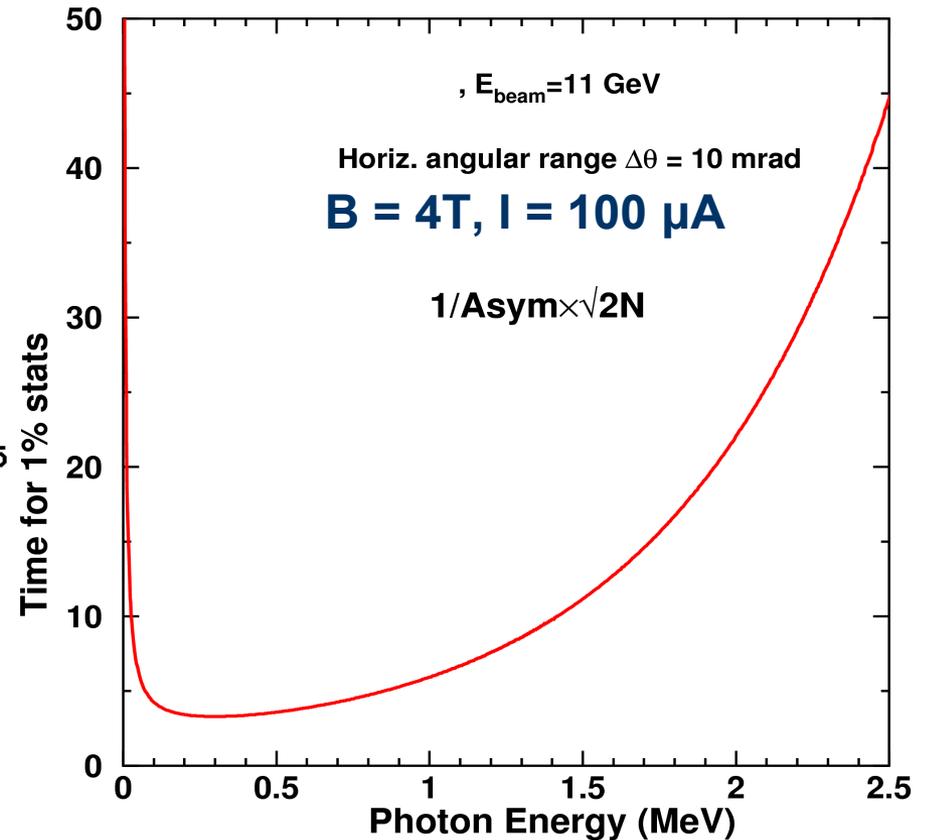


# “Spin Light” - Some Characteristics



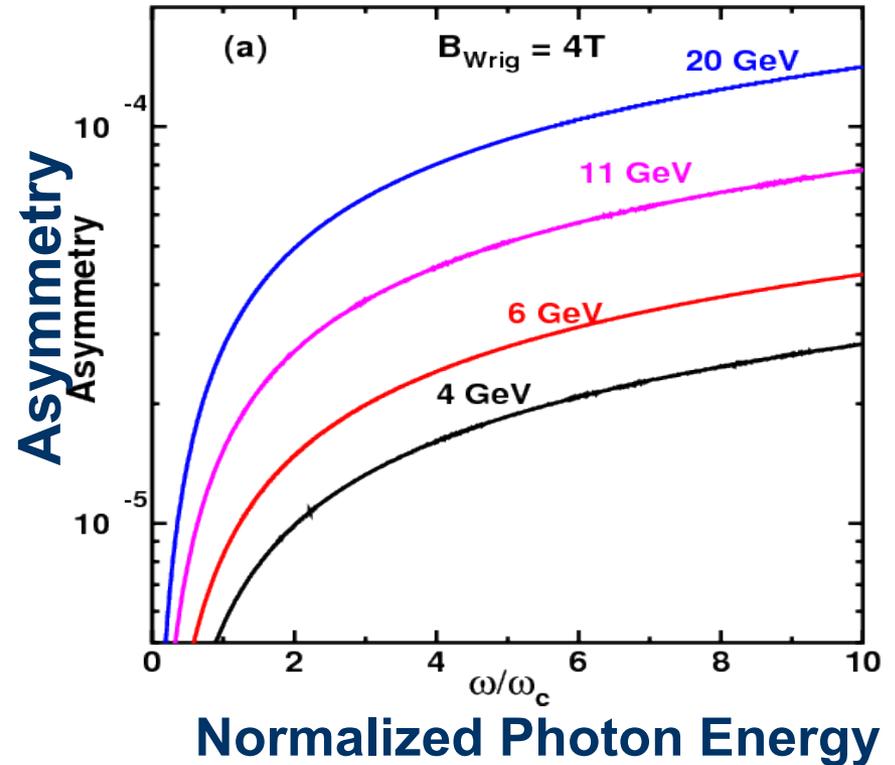
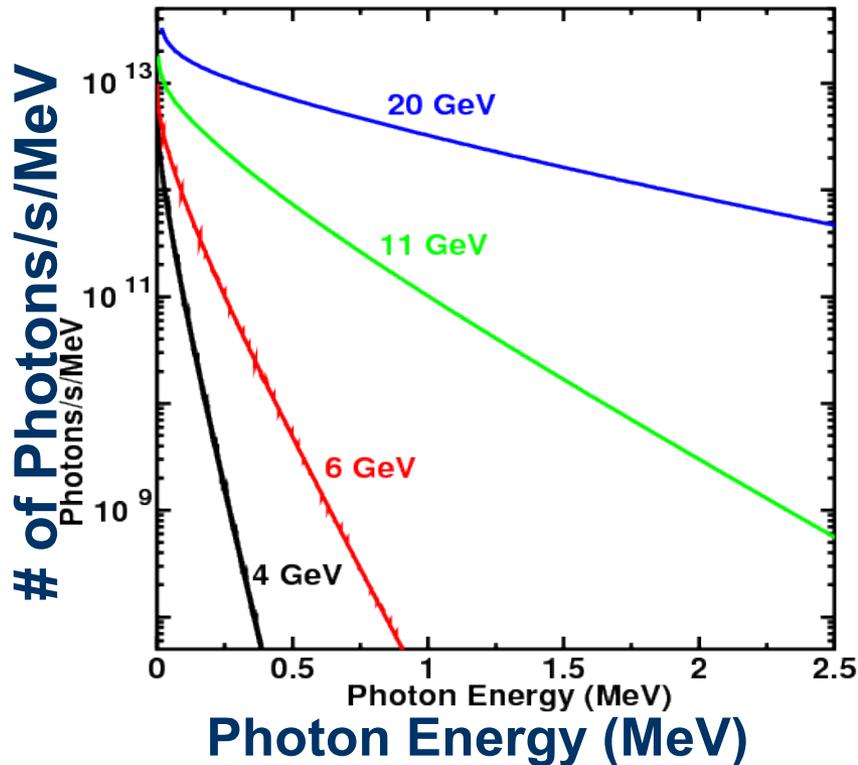
small asymmetry

But high rates imply  
1% statistics in  $\sim 10 \text{ sec}$   
Assuming ion chambers efficiency  $\sim 10\%$



# “Spin Light” - Energy Dependence

At fixed  $B = 4\text{T}$ ,  $I = 100\ \mu\text{A}$  and  $\Delta\theta = 10\ \text{mrad}$



# of photons increases sharply with energy  
but asymmetry increases slowly

# Conceptual Design

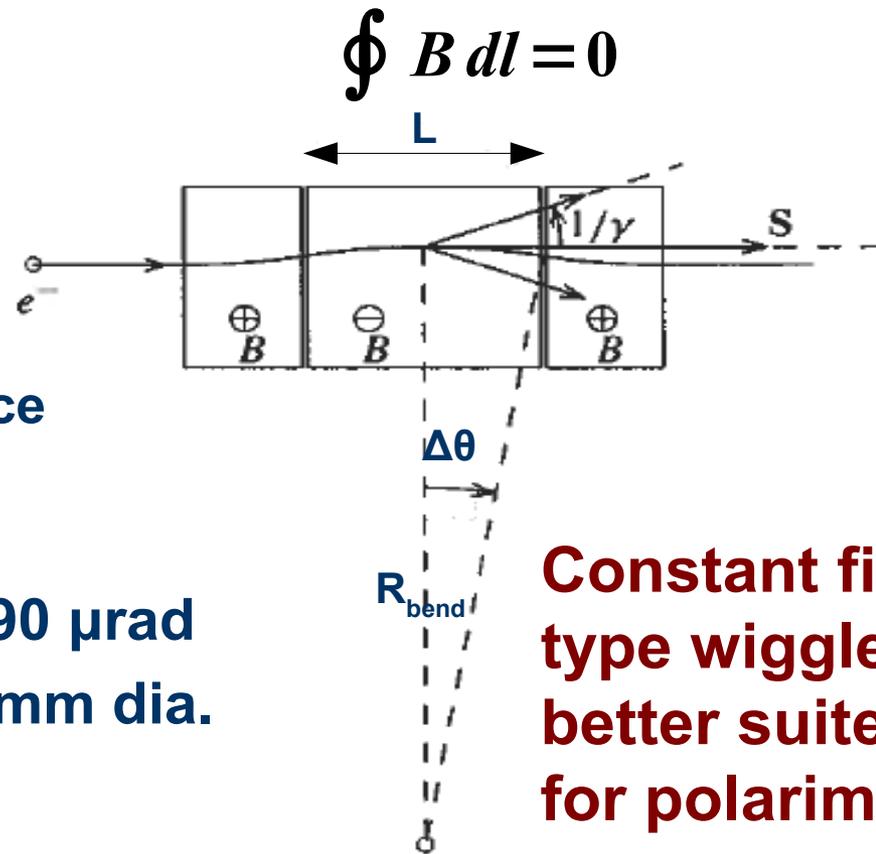
A Source of Spin Light: a 3 pole wiggler

$$R_{bend} = \frac{\gamma m_e c}{e B}$$

$$L = R_{bend} \Delta \theta$$

Horizontal angular acceptance  
 $\Delta \theta$  fixed to 10 mrad

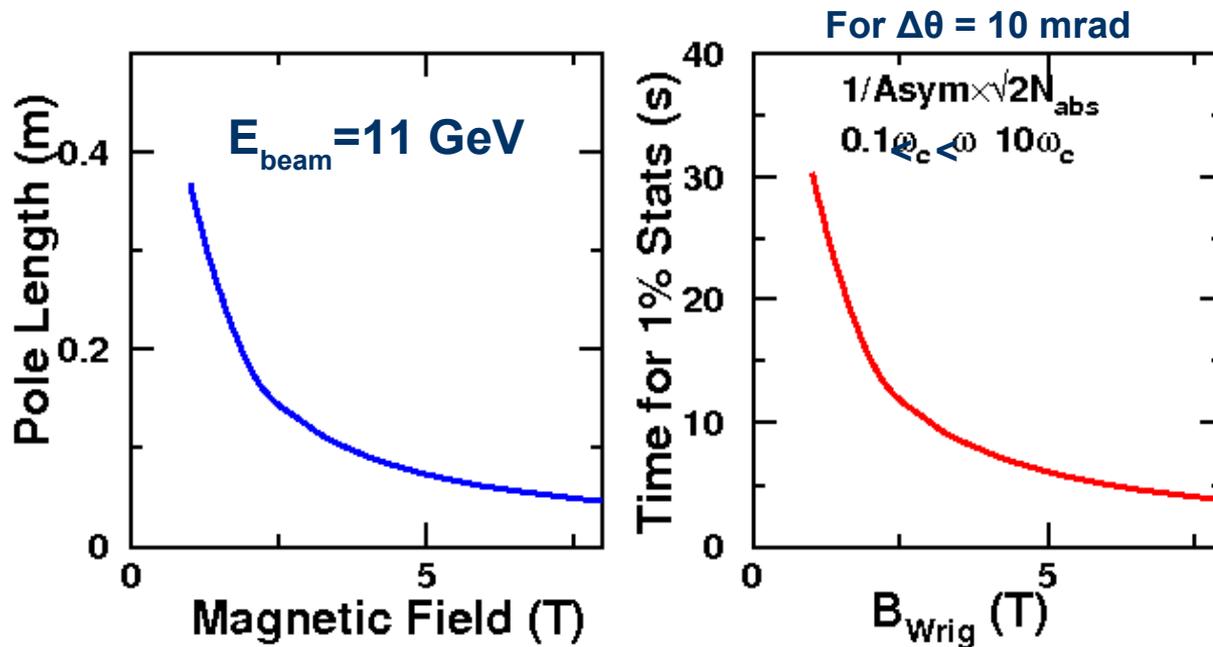
For  $E_e = 11$  GeV, spot size = 90  $\mu$ rad  
 i.e. 10m from the source  $\sim 1$  mm dia.



**Constant field  
 type wiggler  
 better suited  
 for polarimeter**

# Conceptual Design

## Choosing the $B_{\text{wiggler}}$



Pole length and figure of merit have similar dependence on the B- field.

**Choose  $B = 4\text{T}$**

# Is the Wiggler Non-invasive?

**Effects of fluctuations related to the quantum nature of SR**  
(carefully studied for the Jlab recirculating arcs for up to  $E_e = 24$  GeV)

# Is the Wiggler Non-invasive?

Effects of fluctuations related to the quantum nature of SR

(carefully studied for the Jlab recirculating arcs for up to  $E_e = 24$  GeV)

Mean # of  $\gamma$  emitted  
per  $e^-$  per radian

$$= n = 20.62 E_e$$

$$\Delta E = \sqrt{n} E_c$$

Mean energy of photon =  $E_c = 3\hbar c \gamma^3 / 2R$

For  $E_e = 11$  GeV, 10 mrad bend and  $R=10$ m:  $n \sim 2$ ,  $E_c = 199$  keV and

$$\Delta E/E = 2.5 \times 10^{-5}$$

# Is the Wiggler Non-invasive?

Effects of fluctuations related to the quantum nature of SR

(carefully studied for the Jlab recirculating arcs for up to  $E_e = 24$  GeV)

Mean # of  $\gamma$  emitted  
per  $e^-$  per radian

$$= n = 20.62 E_e$$

$$\Delta E = \sqrt{n} E_c$$

Mean energy of photon =  $E_c = 3\hbar c \gamma^3 / 2R$

For  $E_e = 11$  GeV, 10 mrad bend and  $R=10$ m:  $n \sim 2$ ,  $E_c = 199$  keV and

$$\Delta E/E = 2.5 \times 10^{-5}$$

Similarly  $\Delta \theta_e = E_\gamma \sin \Theta_\gamma / E_e = 1.5 \times 10^{-8}$  rad

Depolarization  
due to wobble

$$\Delta p = \frac{\tau_{(transit\ thru\ wiggler)}}{\tau_{(depol.)}} = \frac{L_{wrig} / \beta c}{98.66 (1900 R^2 / E^5)} \sim 1.1 \times 10^{-9} \%$$

# “Spin Light” Detector

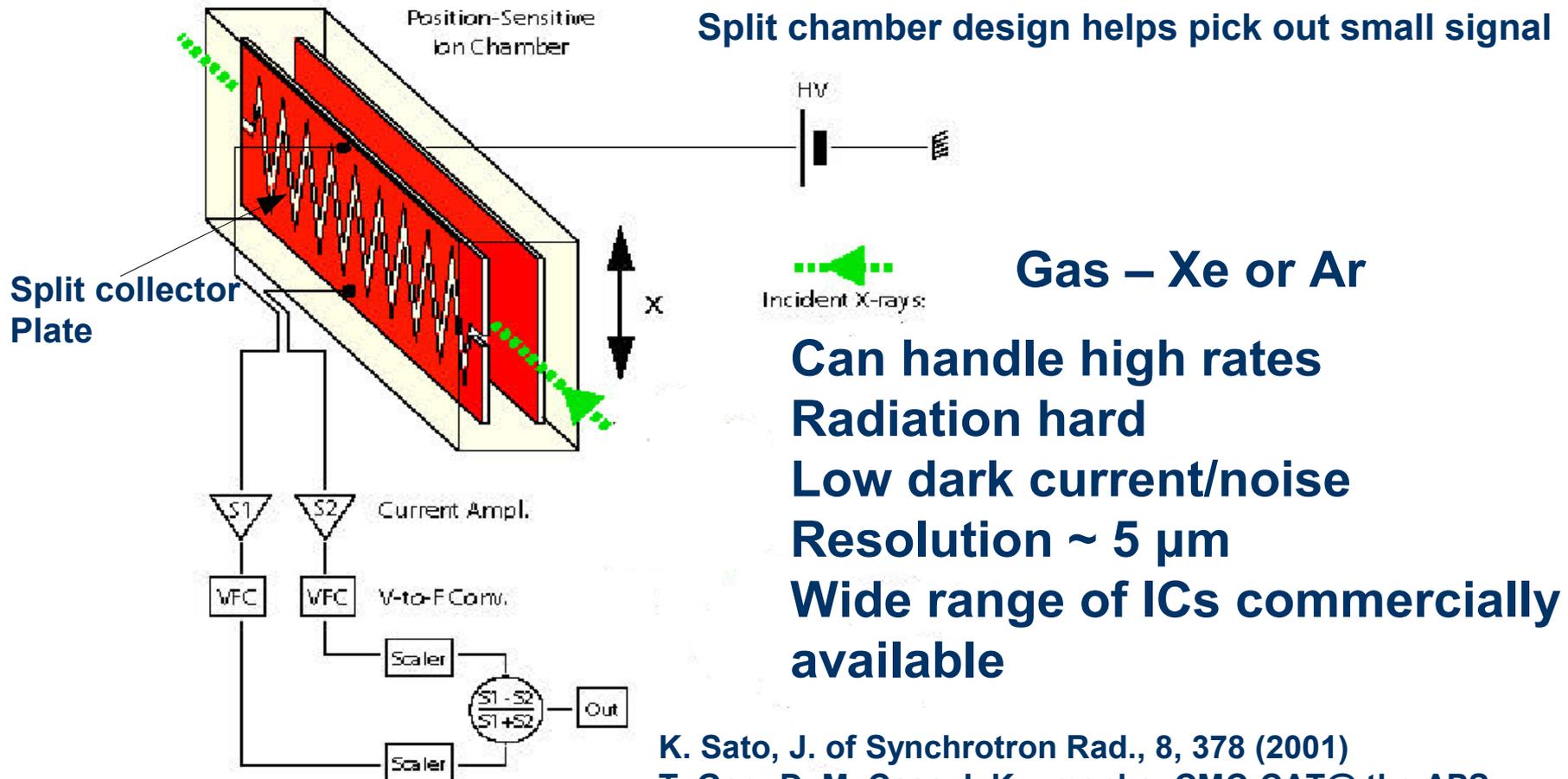
**A Detector of Synchrotron + Spin light (X-rays)**

**A transparent differential ionization chamber**

# “Spin Light” Detector

**A Detector of Synchrotron + Spin light (X-rays)**

**A transparent differential ionization chamber**



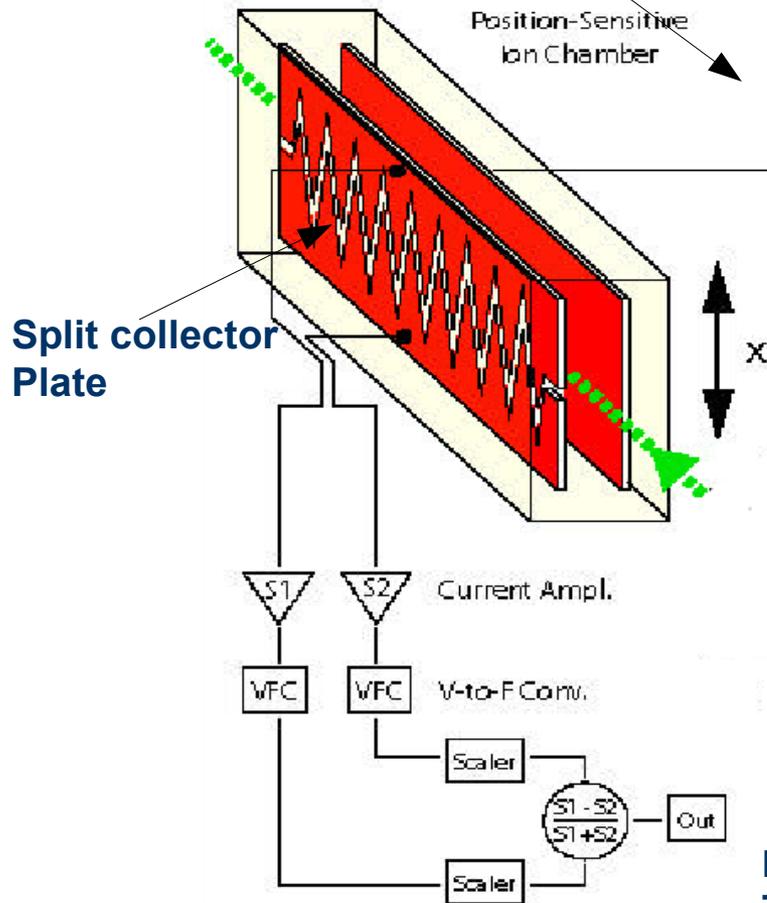
K. Sato, J. of Synchrotron Rad., 8, 378 (2001)

T. Gog, D. M. Casa, I. Kuzmenko, CMC-CAT@ the APS

# “Spin Light” Detector

**A Detector of Synchrotron + Spin light (X-rays)**

**A transparent differential ionization chamber**



**Split chamber design helps pick out small signal**  
**Visible portion can be used to center chamber**

  
Incident X-rays:

**Gas – Xe or Ar**

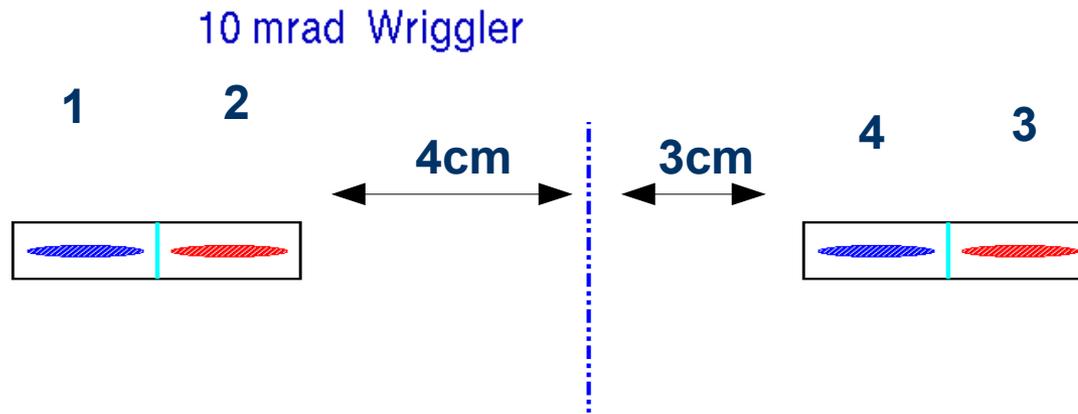
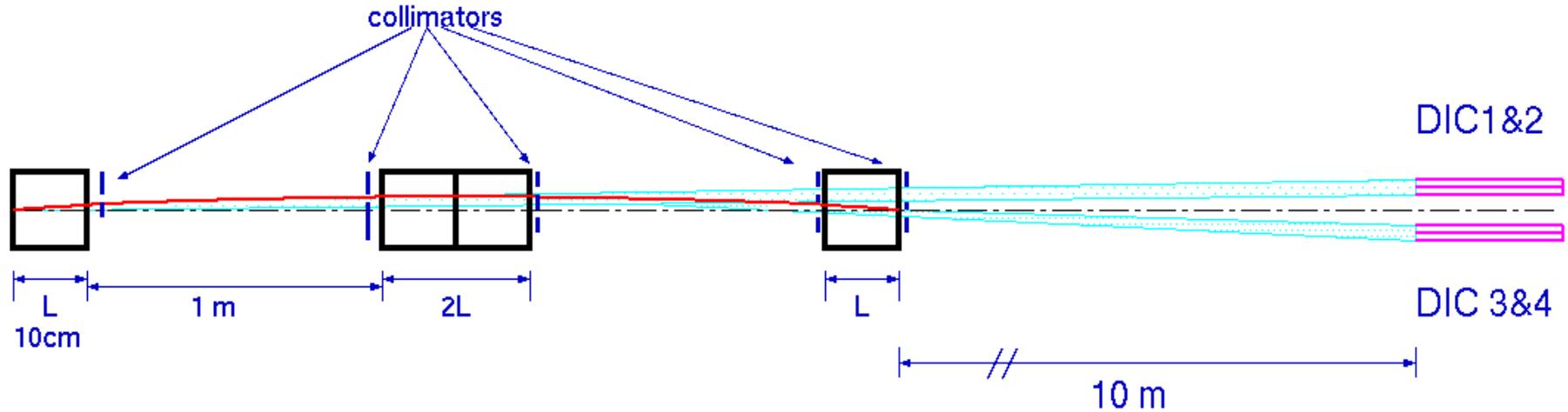
**Can handle high rates**  
**Radiation hard**  
**Low dark current/noise**  
**Resolution ~ 5  $\mu\text{m}$**   
**Wide range of ICs commercially available**

K. Sato, J. of Synchrotron Rad., 8, 378 (2001)

T. Gog, D. M. Casa, I. Kuzmenko, CMC-CAT@ the APS

# A “Spin Light” Polarimeter

## Putting it all together



**SR spot size  
10 m from the  
wigglers**

**With appropriate slits/collimators after each  
pole of the wiggler**

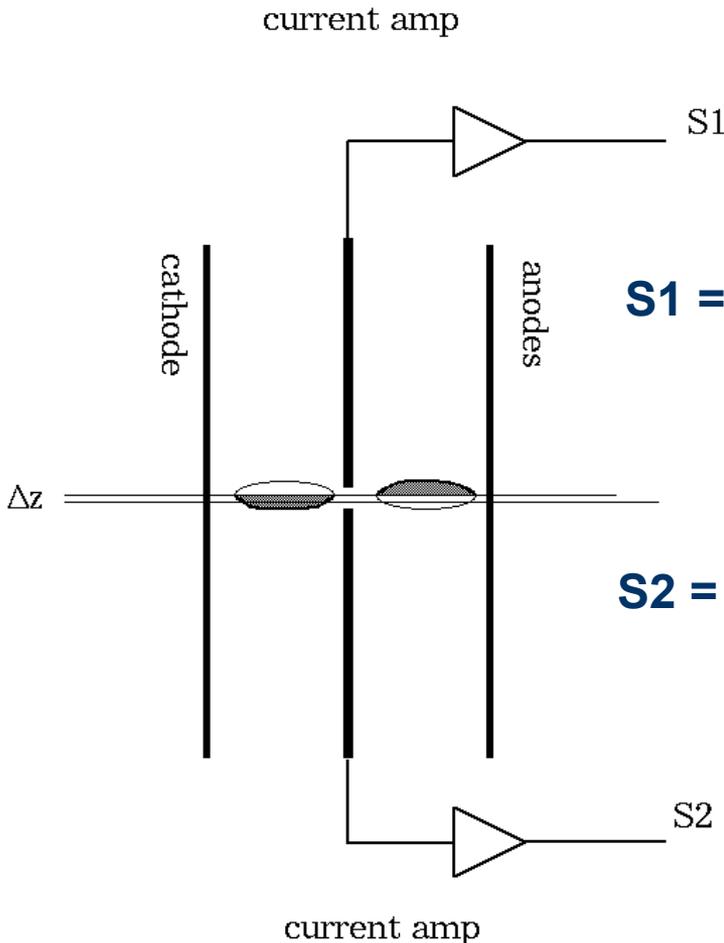
# Comparing Polarimetry Techniques

Sub-1% polarimetry requires multiple independent measurements of the beam polarization.

- **A Spin-light polarimeter is complimentary to the more popular Compton & Möller polarimeters.**
- **It will provide an independent measurement with completely different systematics**

| Compton   | Spin-Light                          | Möller                                |
|---|-------------------------------------|---------------------------------------|
| Non-invasive/continuous                                   | Non-invasive/continuous             | invasive                              |
| Analyzing power is energy dependent                       | Analyzing power is energy dependent | Analyzing power is energy independent |
| Ideal for high currents                                   | Ideal for high currents             | Used at low currents only             |
| Target is 100% polarized                                  | No target needed                    | Target is <10% polarized              |
| e & $\gamma$ detection provide 2 independent polarimeters | -                                   | -                                     |

# The Signal



## Current mode operation

$$S1 = (N_{SR}^l + \Delta N_{spin}^l + \Delta N_z^l) - (N_{SR}^r - \Delta N_{spin}^r + \Delta N_z^r)$$

Vertical beam motion cancels out

$$S2 = (N_{SR}^l - \Delta N_{spin}^l - \Delta N_z^l) - (N_{SR}^r + \Delta N_{spin}^r - \Delta N_z^r)$$

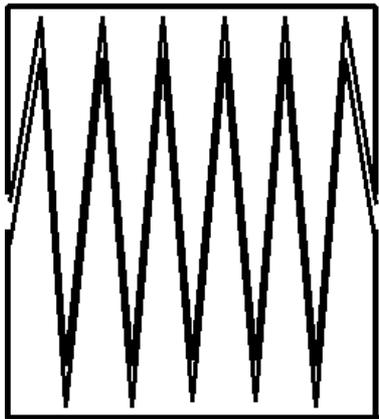
$$(S1 - S2) = 4\Delta N_{spin}$$

$$(S1 + S2) = 0$$

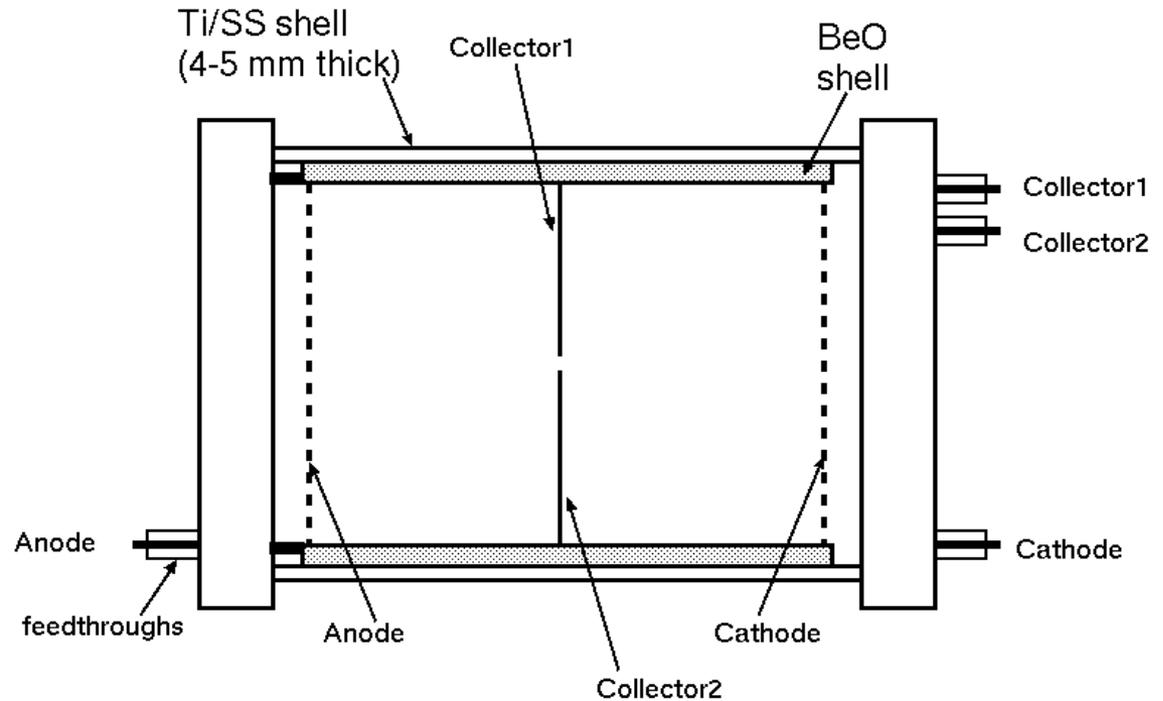
This sum for the two helicities can be used to separate out the efficiency and B-field related dilutions.

# The Detector R&D Proposal

## Stage 1: Develop a split plane differential ionization chamber



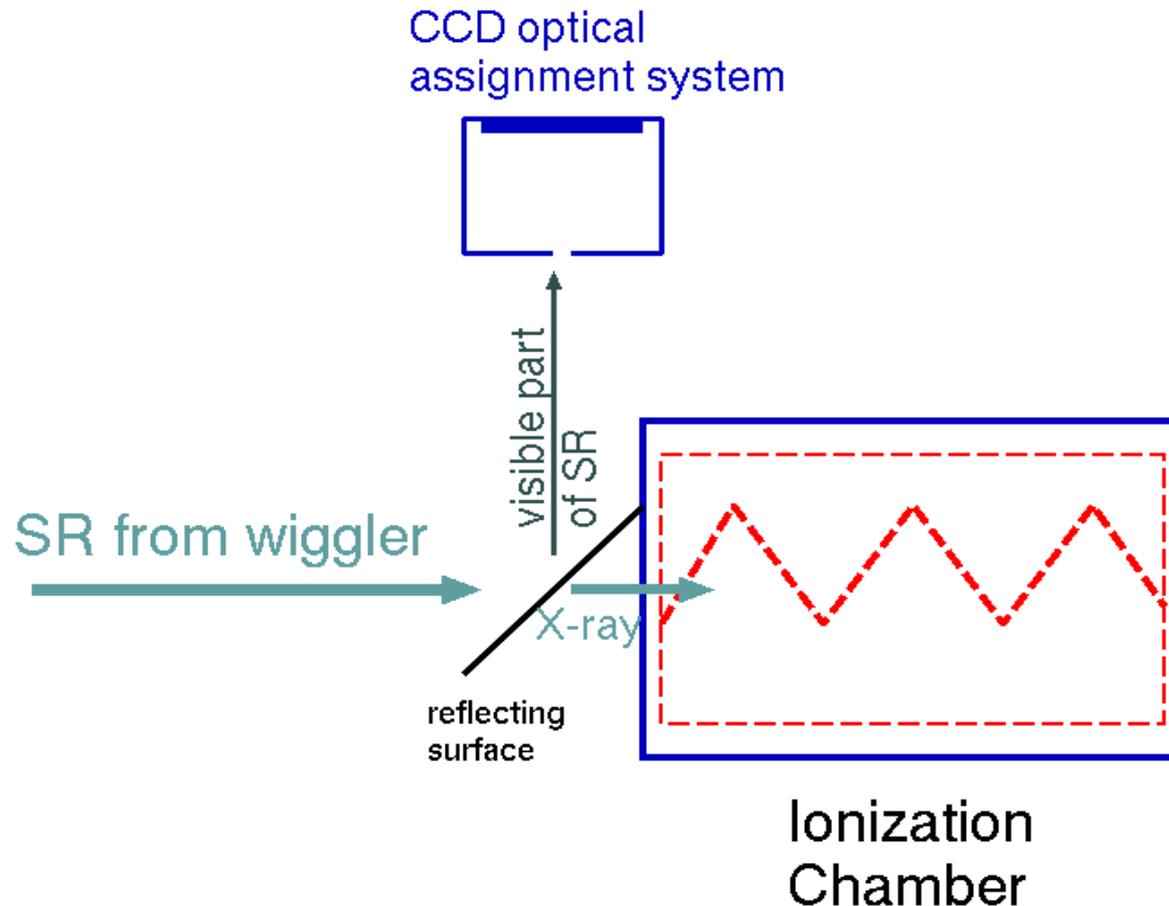
side view of  
split collector plate



Front view of DIC

# The Detector R&D Proposal

## Stage 1: Develop a CCD based ionization chamber alignment system



# The Detector R&D Proposal

## Stage 1: Test the DIC in the JLab Hall-A Compton chicane

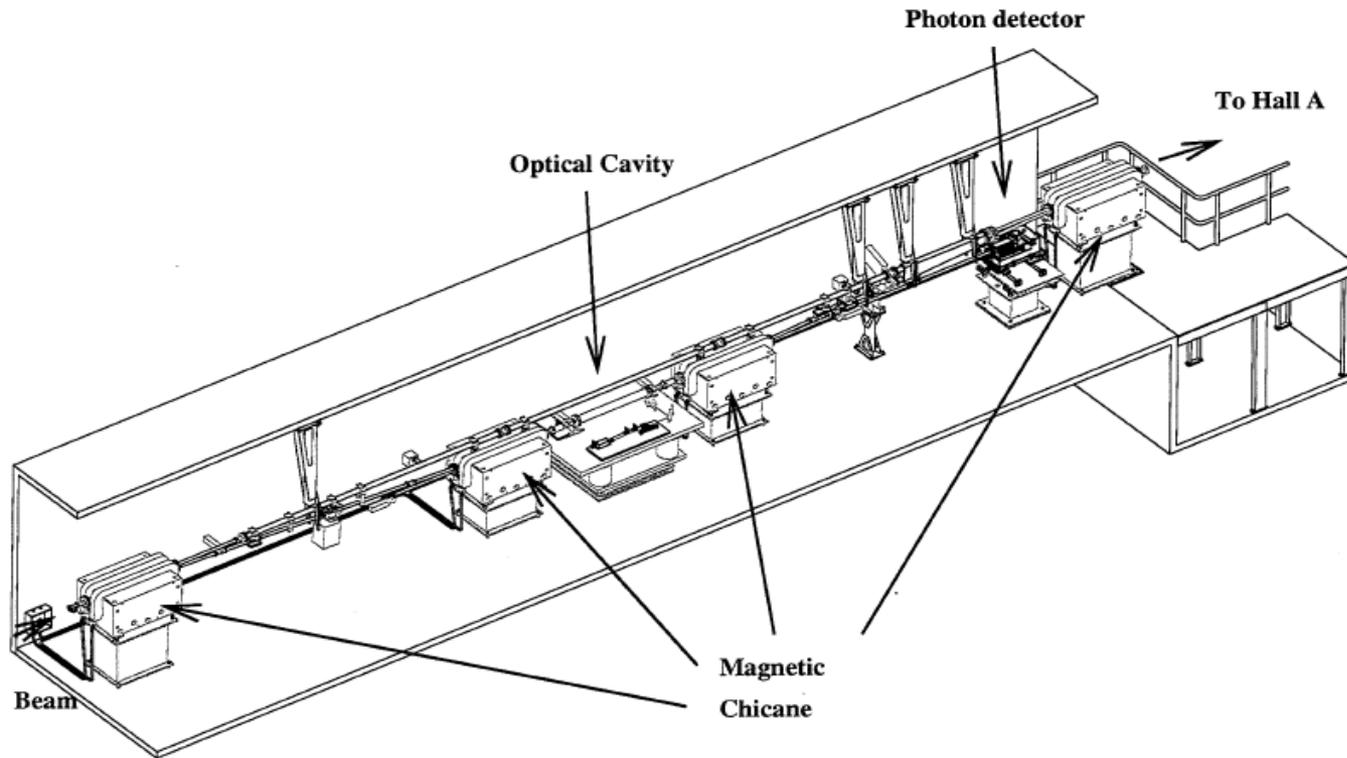


Fig. 14. The Compton Polarimeter Setup at TJNAF Hall A. Total length 15 m.

# The Detector R&D Proposal

**Stage 2: Develop a full prototype spinlight polarimeter** (a set of dual ionization chambers and a suitable wiggler magnet.)

| Activity   | Year 1 | Year 2 | Year 3 |
|--|--------|--------|--------|
| Design and build prototype DIC                     | ✓      | ✓      |        |
| Test DIC in Hall A Compton beamline                |        |        | ✓      |
| Design CCD based alignment system                  | ✓      | ✓      |        |
| Design and build set of dual DIC                   |        |        | ✓      |
| Build CCD system                                   |        |        | ✓      |
| Design wiggler magnet                              |        | ✓      |        |
| Design slits and collimators                       |        | ✓      |        |
| Identify suitable wiggler magnet at the APS        |        |        | ✓      |
| Select suitable site to test prototype polarimeter |        |        | ✓      |

# The Detector R&D Proposal

The collaboration has extensive polarimetry and detector development experience at DESY, JLab, Mainz, and RHIC.

**Students & Postdocs**      **Grad. Students:**  
**Edward Leggett** (MSU) and **Valerie Gray** (W&M) and **MSI student(s)** (Stony Brook U., during yrs. 2 and 3)  
**Post doc: M. Shabestari** (MSU, 25% FTE only)  
**Additional post-doc during 2<sup>nd</sup> / 3<sup>rd</sup> yr. desirable**

## Funding Request

| Item                   | Year 1 | Year 2 | Year 3 | Total  |
|------------------------|--------|--------|--------|--------|
| 0.5 Grad Student (MSU) | \$17k  | \$17k  | \$17k  | \$51k  |
| 0.5 Grad student (W&M) | \$17k  | \$17k  | \$ 17k | \$51k  |
| Equipment              | \$46k  | \$44k  | \$48k  | \$138k |
| Travel                 | \$10k  | \$10k  | \$10k  | \$30k  |
| Total                  | \$90k  | \$88k  | \$92k  | \$270k |

# The Detector R&D Proposal

| Equipment                       | Total cost |
|---------------------------------|------------|
| prototype DIC                   | 10000      |
| Split plane electrodes          | 5000       |
| Electronics for DIC(2 channels) |            |
| current amps                    | 8000       |
| High voltage power supplies     | 10000      |
| V-to-Fs and scalars             | 15000      |
| VME crate                       | 10000      |
| Single board computer           | 7000       |
| Gas Handling system             | 10000      |
| The Dual DICs                   |            |
| Custom dual DICs                | 12000      |
| with split collector            |            |
| additional amps                 | 8000       |
| CCD alignment system            | 25000      |
| Custom beamline vacuum          | 10000      |
| elements                        |            |
| slits and collimators           | 8000       |
| Total Equipment Cost            | 138000     |

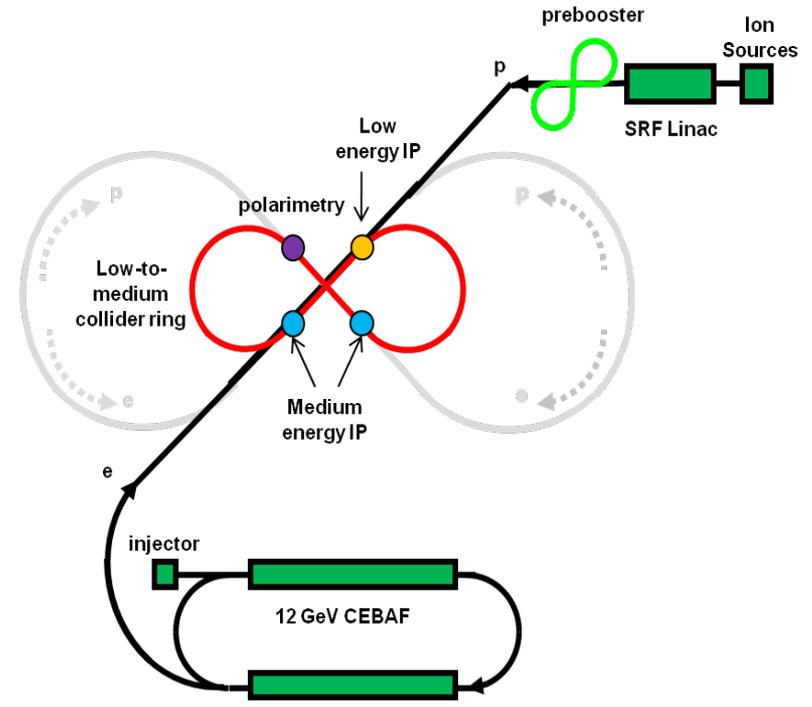
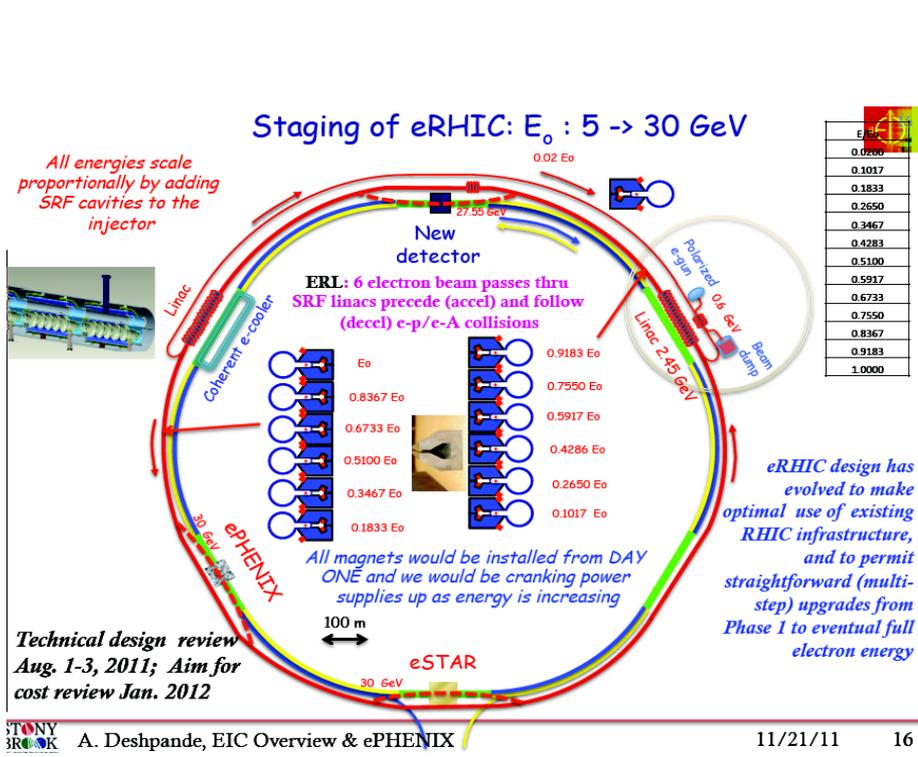
Table 1: Equipment cost breakup

# Summary

- **Spin light based polarimeter is a viable option for precision non-invasive polarimetry**
- **It is based on a well demonstrated concept (for transversely polarized electrons), the necessary technology is readily available and widely used in light sources across the world.**
- **We propose to develop a split plane ionization chamber and demonstrate proof of principle for longitudinally polarized electrons using 12 GeV beam at JLab.**
- **Begin developing a full prototype spin light polarimeter**

# Possible Locations at an EIC

A 2.5 m long straight section of e-beamline and a 10 m long free flight path (along the beamline) for SR photons to a set of DICs



## EIC@RHIC

## EIC@JLab